

# Pulsar emission of positrons and electrons, and its connections with gamma rays

Fiorenza Donato

Torino University

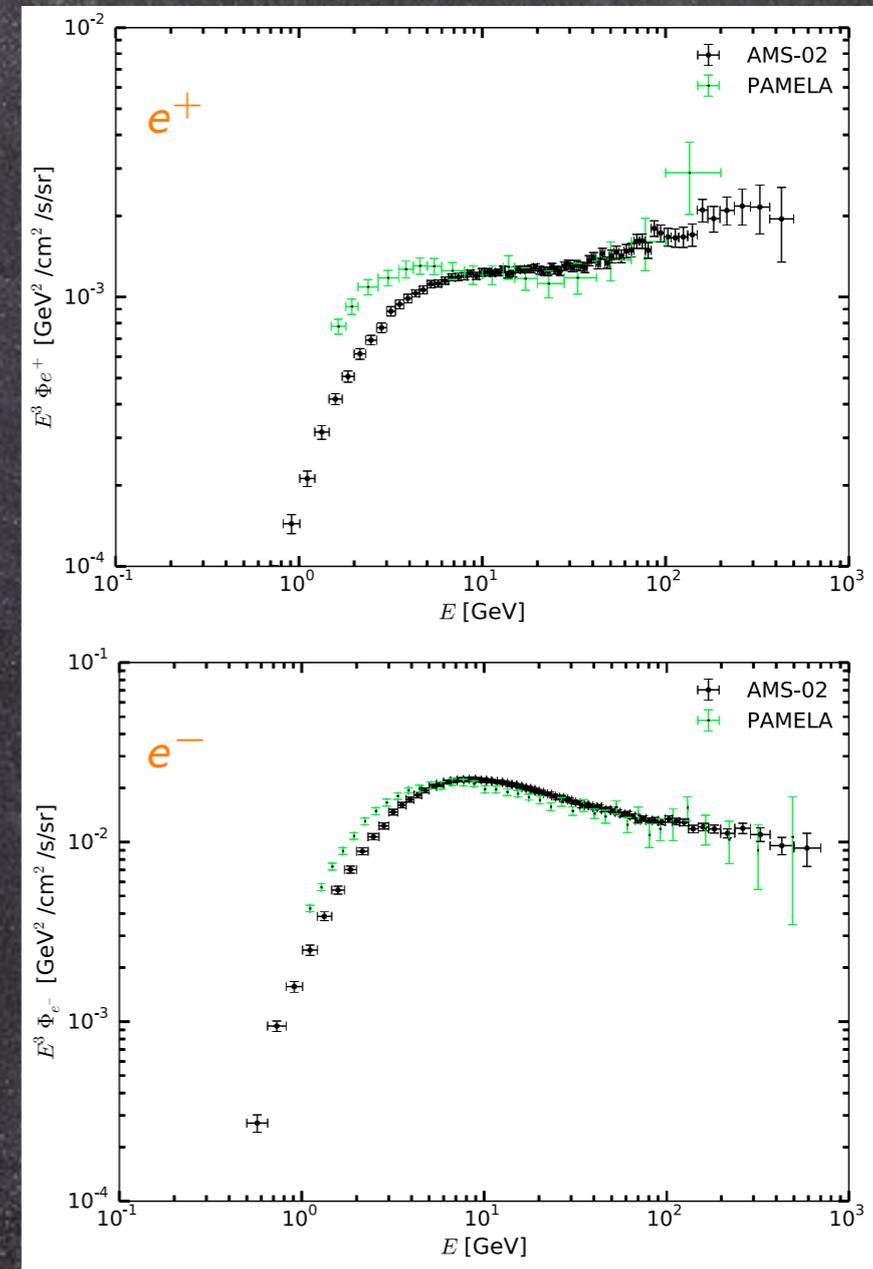
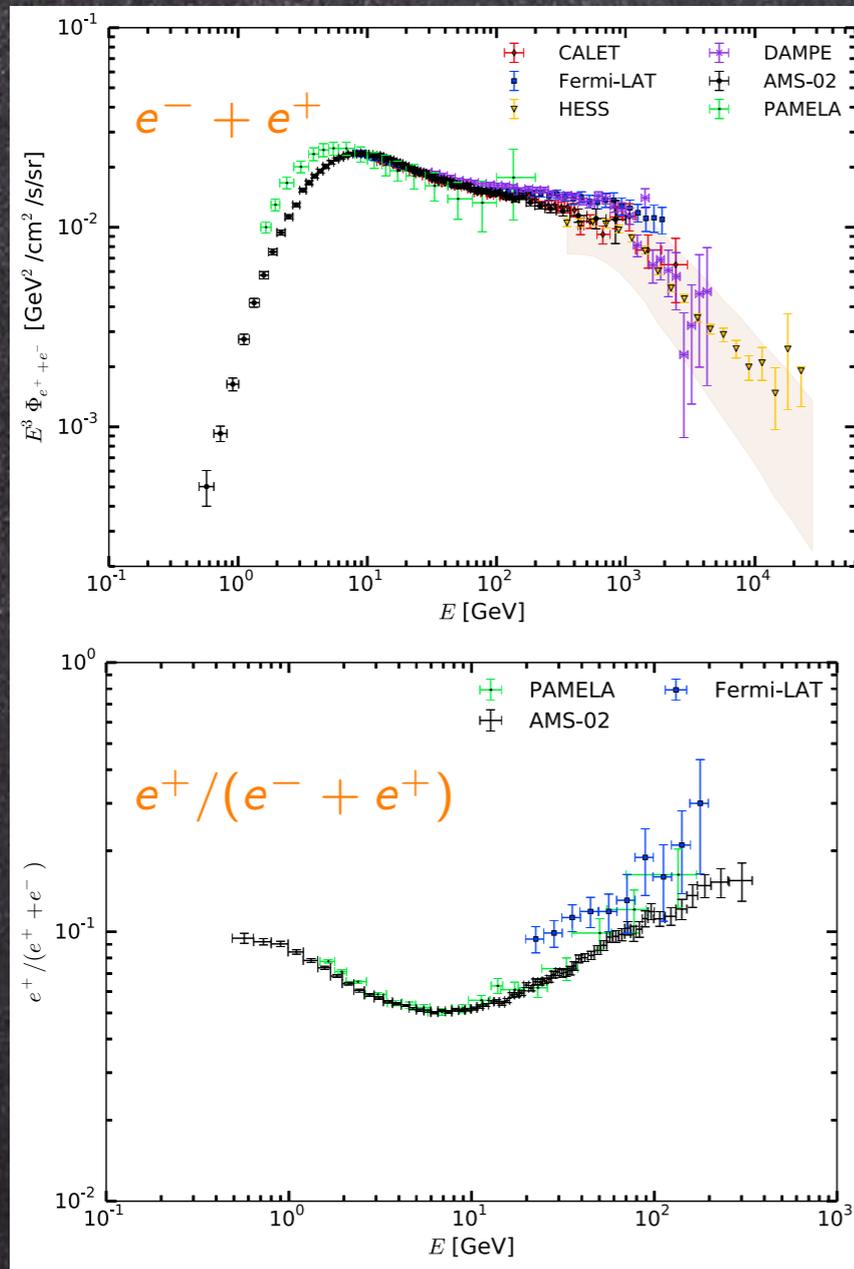
Virtual seminar on Multi-messenger astronomy

June 3rd, 2020



works done in collaboration with Mattia Di Mauro and Silvia Manconi

The journey started with the attempt  
- shared by many - to interpret the  $e^+$  data



Unprecedented statistics and energy coverage

# Sources of $e^\pm$ in the Milky Way

- Inelastic hadronic collisions (asymm.)
- Pulsar wind nebulae (PWN) (symm.)
- Supernova remnants (SNR) (only  $e^+$ )
- (Particle Dark Matter annihilation)

# From the source to a detector

**Sources:** production mechanism  
acceleration to high energies  
injection into the interstellar space

**Propagation:** diffusion on the magnetic inhomogeneities of the Galaxy  
energy losses and gains (B, E, cross sections)  
convection

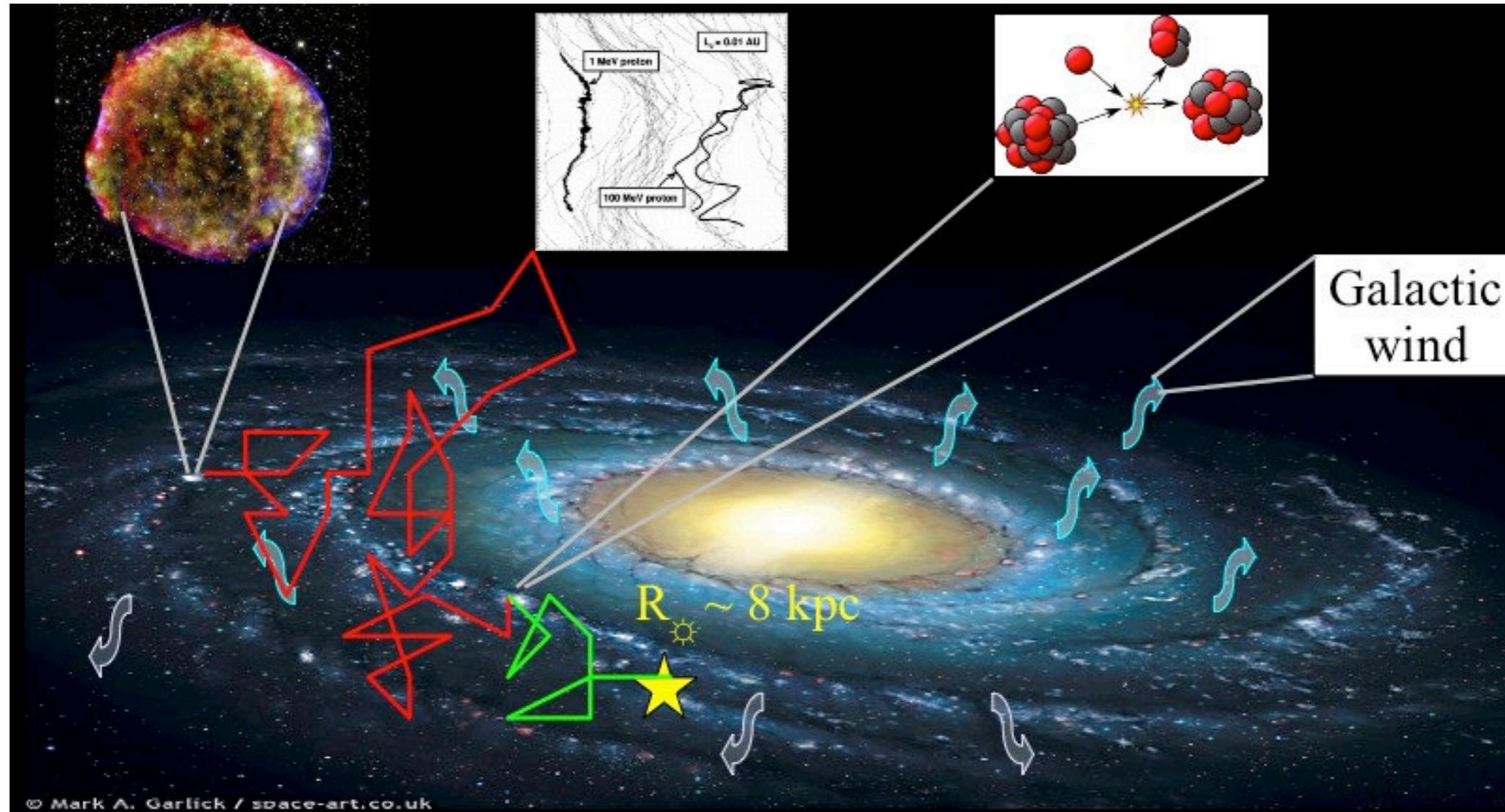
**Energy losses:** from the production site to the detector

# Propagation in the Galaxy

Sources

Diffusion

Fragmentation



from D. Maurin

The galactic disc is embedded in a diffusive halo ( $L \sim 4-10 \text{ kpc}$ ) where particles diffuse and can lose energy

# Propagation equation

$$\frac{\partial \psi}{\partial t} - \nabla \cdot \{D(E) \nabla \psi\} + \frac{\partial}{\partial E} \left\{ \frac{dE}{dt} \psi \right\} = Q(E, \mathbf{x}, t)$$

diffusion

en. losses

source spectrum

**Diffusion:**  $D(x, R)$  a priori

usually assumed isotropic in the Galaxy:  $D(R) = D_0 R^\delta$

$D_0$  and  $\delta$  fixed by secondary/primary nuclei (Kappl+15; Genolini+15 (K15))

**Energy losses:** Synchrotron on the galactic  $B \sim 3.6 \mu\text{G}$

full relativistic of Compton effect (w/ Klein-Nishijima)

on photon fields (stellar, CMB, UV, IR)

**Solution of the eq.:** semi-analytic (Maurin+ 2001, Donato+ 2004, ...), USINE codes

or fully numerical: GALPROP, DRAGON codes

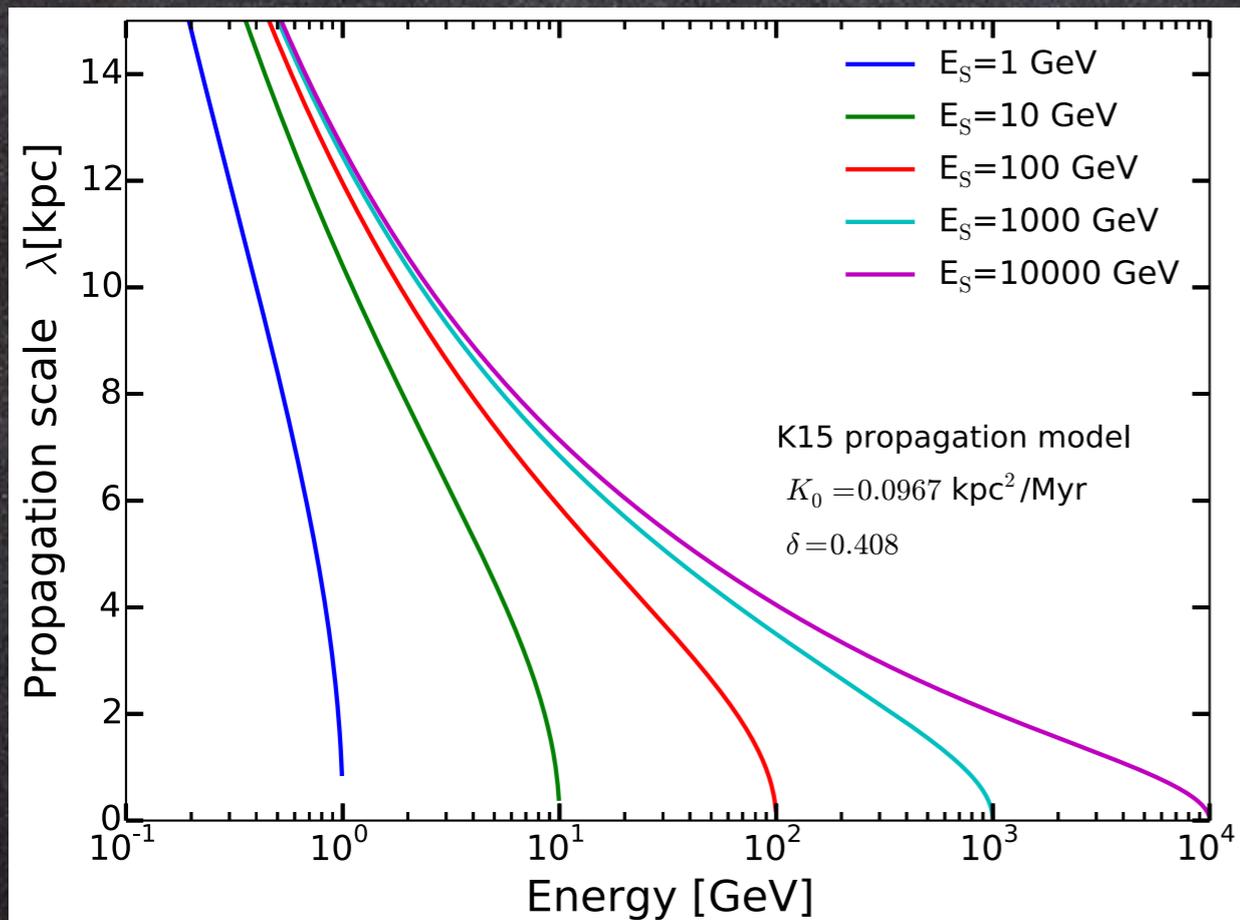
Geometry of the Galaxy: cylinder with height  $L$

# Detected $e^\pm$ are local

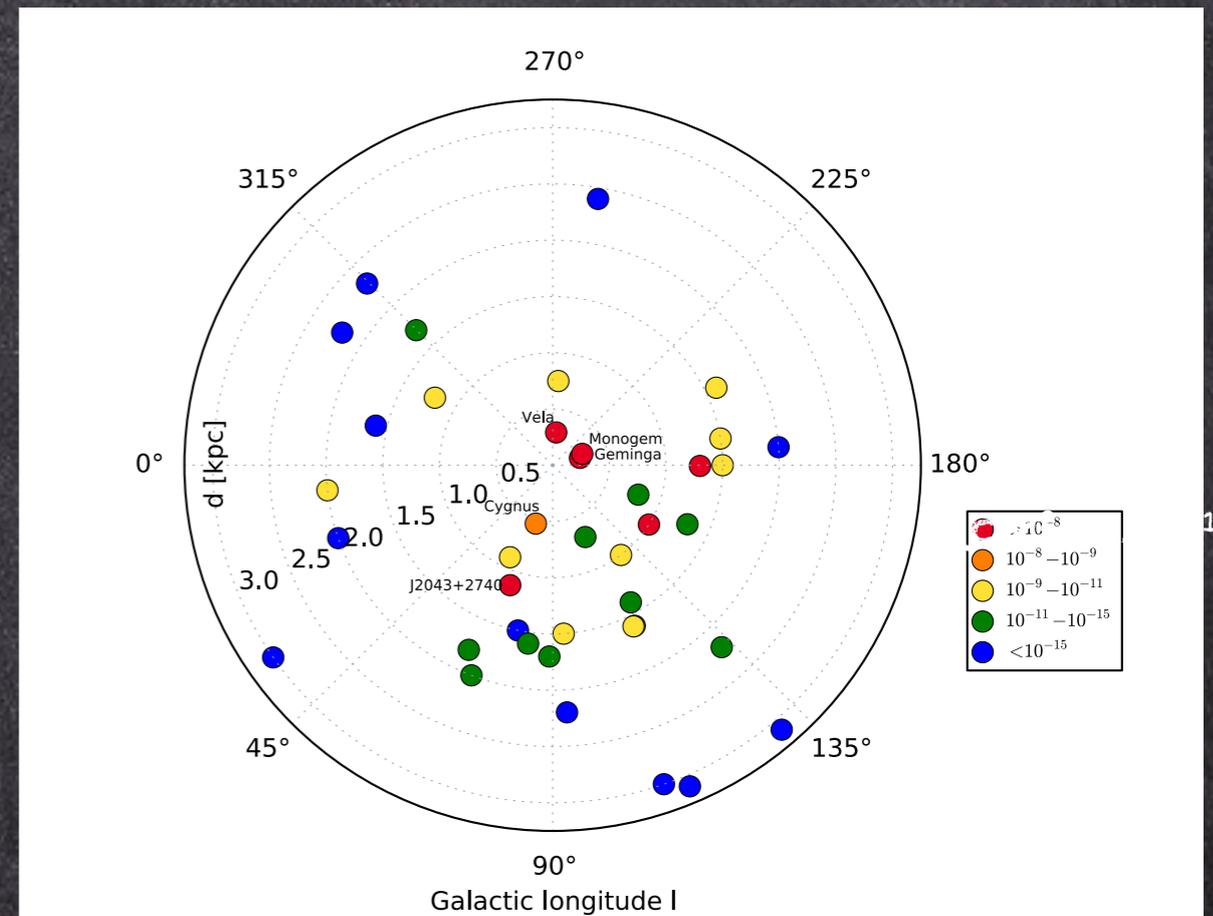
$$\lambda^2(E, E_S) = 4 \int_E^{E_S} dE' \frac{D(E')}{b_{\text{loss}}(E')}$$

Typical propagation length in the Galaxy

Manconi, Di Mauro, FD JCAP 2017



$e^-$ ,  $e^+$  have strong radiative cooling and arrive at Earth if produced within few kpc around it



Most powerful sources within 3 kpc from the Sun.  
SNRs ( $e^-$ ) and PWN ( $e^+e^-$ )

# A word worth on $e^\pm$ energy losses

Inverse Compton  $e^\pm(E_0) + \gamma(E) \rightarrow e^\pm(E_1) + \gamma(E')$

$E_0 > E_1$  a loss

$E < E'$  a gain, and a **MESSANGER**

X-rays,  $\gamma$ -rays

Synchrotron emission radiated from the interaction with a B  
 $e^\pm$  losses energy

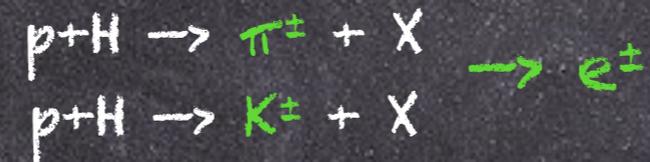
a photon is created, a **MESSENGER**

radio, X-ray

# Secondary $e^+e^-$

Delahaye+ A&A2009, 2010

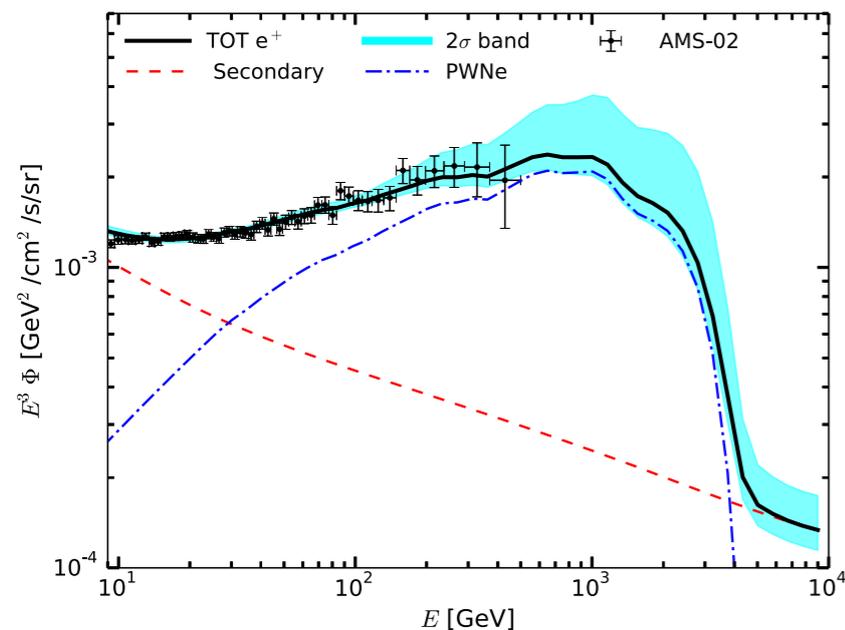
Produced by inelastic collisions of CRs (p and He mosly)  
on the interstellar GAS (H and He)



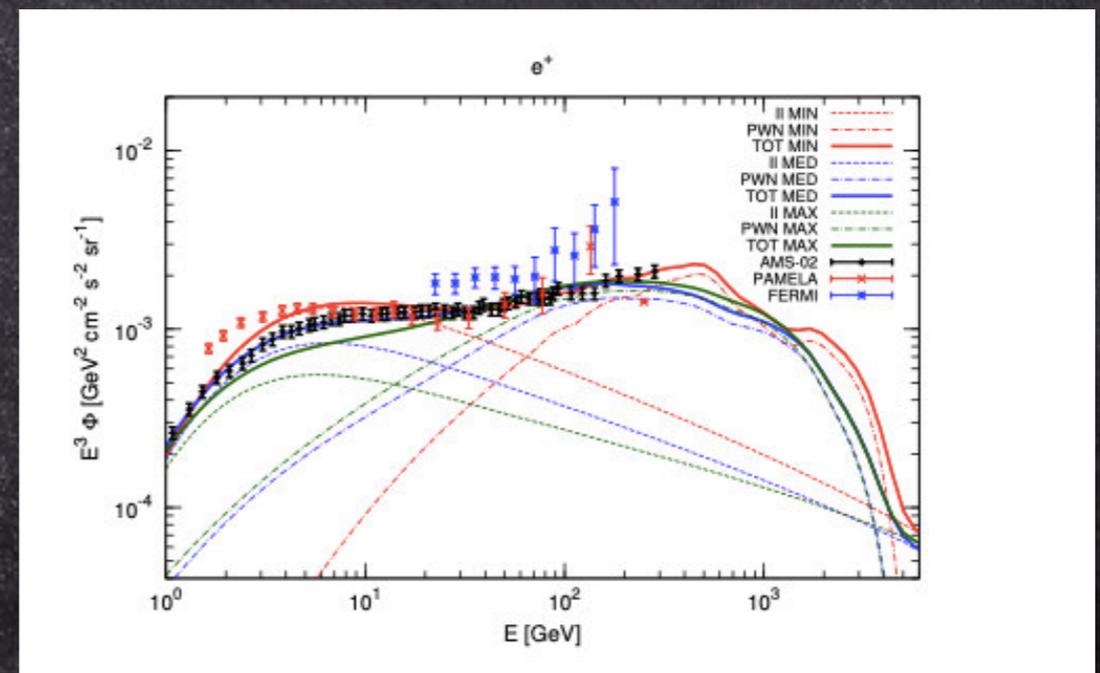
$d\sigma(p+H \rightarrow \pi^\pm + X)$  has to be measured  $\rightarrow$  UNCERTAINTY  $\sim 2$   
 $P(\pi^\pm \rightarrow e^\pm)$  is computed by QED

Need for high energy data cross sections from  
collider experiments

Di Mauro + JCAP 2019



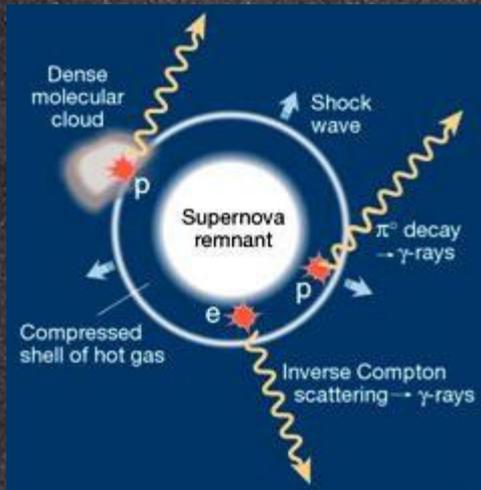
Di Mauro + JCAP 2014



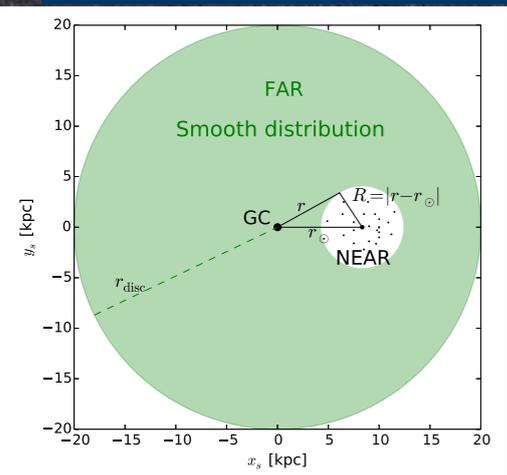
# Electrons from supernova remnants

Ellison+ ApJ 2007; Blasi 2013; Di Mauro+ JCAP 2014; ApJ 2017

SNR are considered the main sources of galactic CRs - nuclei from p to Fe, and e-



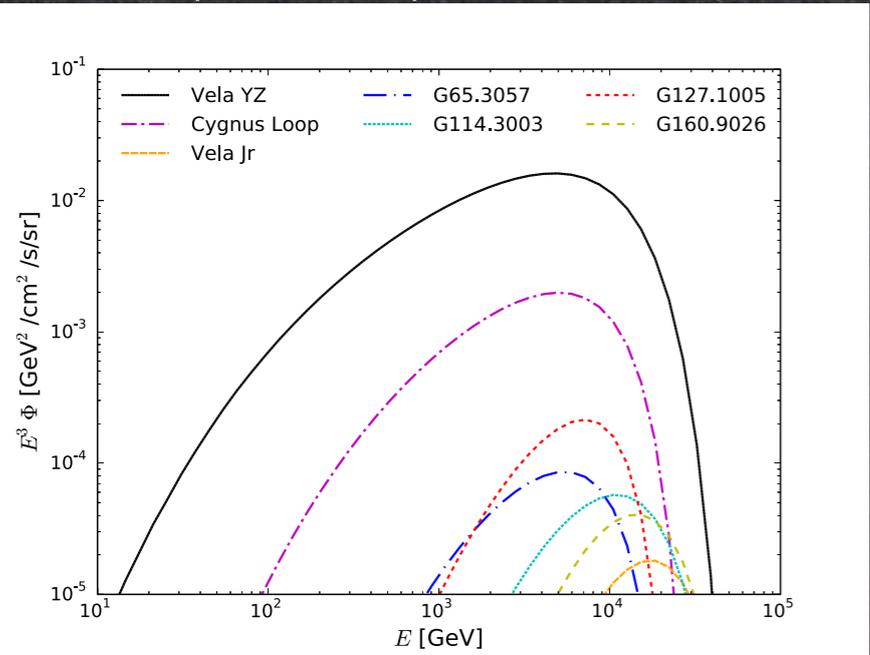
- Hadronic acceleration: evidence of  $\pi^0$  bump (Fermi-LAT+ 2010)
- Leptonic acceleration: evidence of synchrotron emission in radio and X-rays



Injection spectrum:

$$Q(E) = Q_0 \left( \frac{E}{E_0} \right)^{-\gamma} \exp \left( -\frac{E}{E_c} \right)$$

Manconi, Di Mauro, FD JCAP2017; JCAP 2019



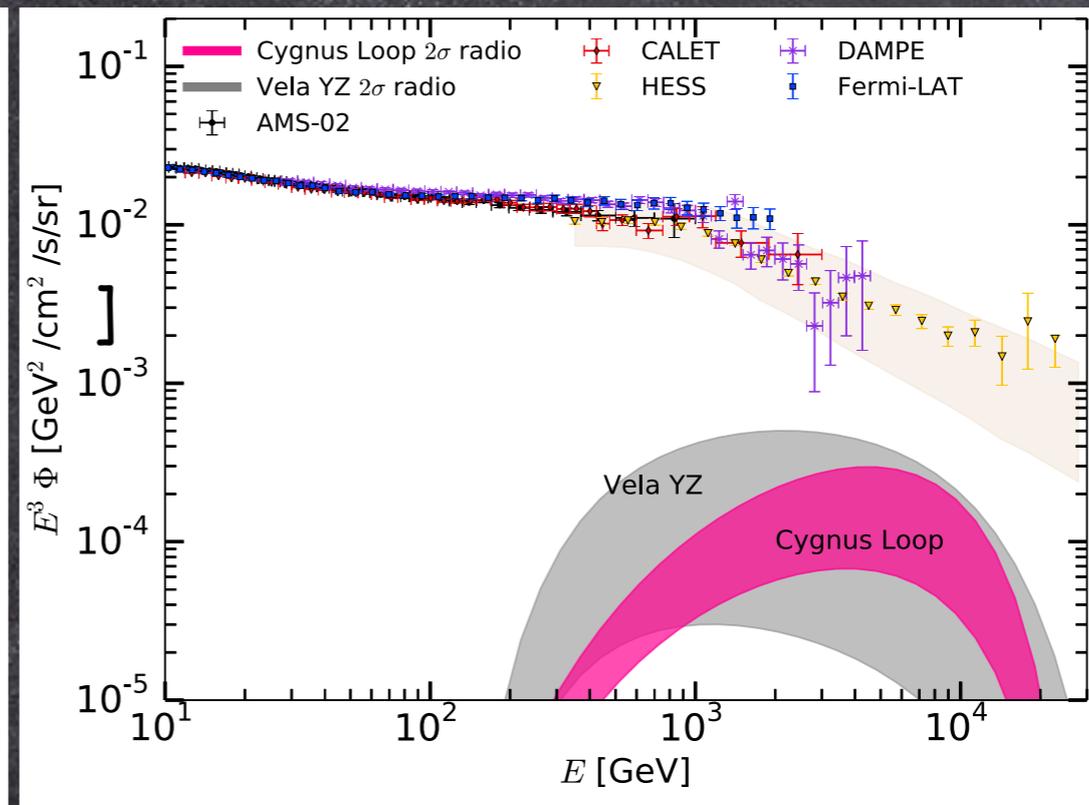
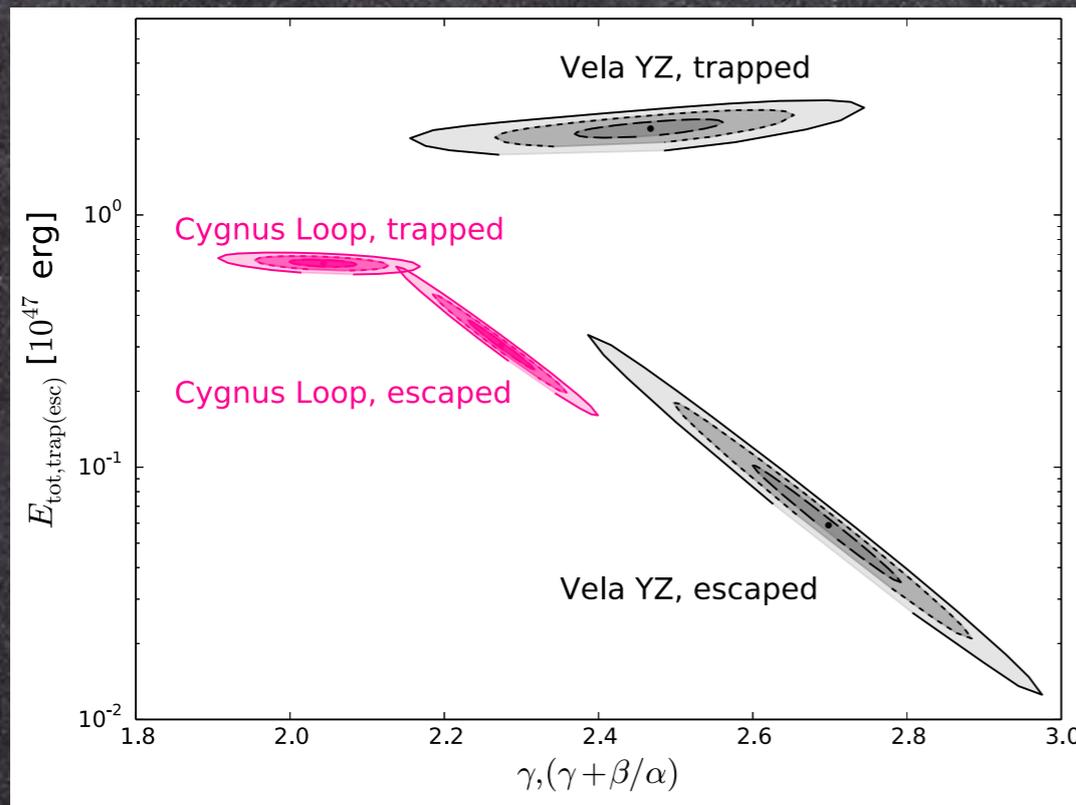
e- flux from near SNR (Vela XY and Cygnus Loop at d < 0.5 kpc)  
 Few SNR can contribute to TeV flux  
 Additional e- from a smooth SNR distribution

# SNR acceleration & radio emission

Manconi, Di Mauro, FD JCAP 2019

Hyp: Radio flux due to synchrotron emission from accelerated  $e^-$  in the SNR

$$Q_{0,\text{SNR}} = 1.2 \cdot 10^{47} \text{GeV}^{-1} (0.79)^\gamma \frac{B_r^\nu(\nu)}{\text{Jy}} \left[ \frac{d}{\text{kpc}} \right]^2 \left[ \frac{\nu}{\text{GHz}} \right]^{\frac{\gamma-1}{2}} \left[ \frac{B}{100 \mu\text{G}} \right]^{-\frac{\gamma+1}{2}}$$



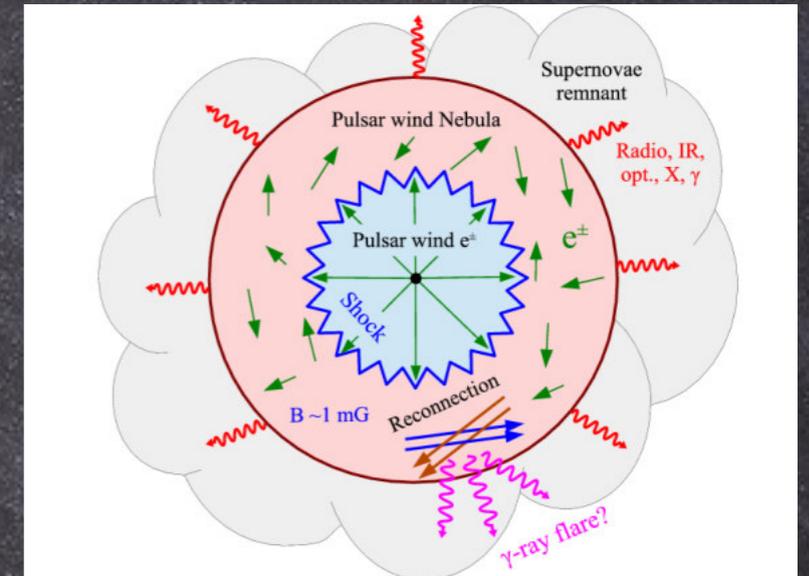
We fit all the available radio data fixing  $B_{\text{Vela}} = 36 \mu\text{G}$  and  $B_{\text{Cygnus}} = 60 \mu\text{G}$ . Vela has energetic trapped  $e^-$ , and only  $E > 88 \text{ GeV}$  have escaped (17 GeV for Cygnus). The flux of  $e^-$  as constrained by radio data contribute few % to the ( $e^+e^-$ ) data

# Pulsars (PWN) as CR $e^+e^-$ sources

Shen ApJL 1970; Amato arXiv:1312.5945; Di Mauro+ JCAP 2014

## Pulsar wind nebulae (PWNe) as engines of $e^\pm$

- High magnetic fields ( $10^9$ – $10^{12}$  G) extract wind of  $e^-$  from the pulsar surface,  $e^\pm$  pairs produced in EM cascades
- Pulsar spin-down energy ( $W_0$ ) is transferred to  $e^\pm$  pairs, accelerated to very high energy with  $Q(E) \sim E^{-\gamma}$ .
- After several kyrs  $e^\pm$  can be **released** in the ISM
- These  $e^\pm$  pairs **radiate** by IC and synch., and shine at many frequencies



$$E_{\text{tot}} = \eta W_0 = \int_0^T dt \int_{E_1}^{\infty} dE E Q(E, t)$$

The total energy  $E_{\text{tot}}$  emitted in  $e^\pm$  by a PWN is a fraction  $\eta$  (efficiency conversion) of the spin-down energy  $W_0$ . Relevant parameters:  $\gamma$  and  $\eta$

# Anisotropy in a diffusion model

Manconi, Di Mauro, FD JCAP 2017

Anisotropy should be computed by development on spherical harmonics.

For one or few sources,

we can expect only the **dipole** term to have some relevance (if any).

In diffusive propagation regime (Ginzburg & Syrovatskii 1964):

$$\Delta = \frac{3K}{c} \left| \frac{\nabla\psi}{\psi} \right|$$

For example, for a source at  $d_s$ :

$$\Delta(E)_{e^+e^-} = \frac{3K(E)}{c} \frac{2d_s}{\lambda^2(E, E_s)} \frac{\psi_{e^+e^-}^s(E)}{\psi_{e^+e^-}^{\text{tot}}(E)}$$

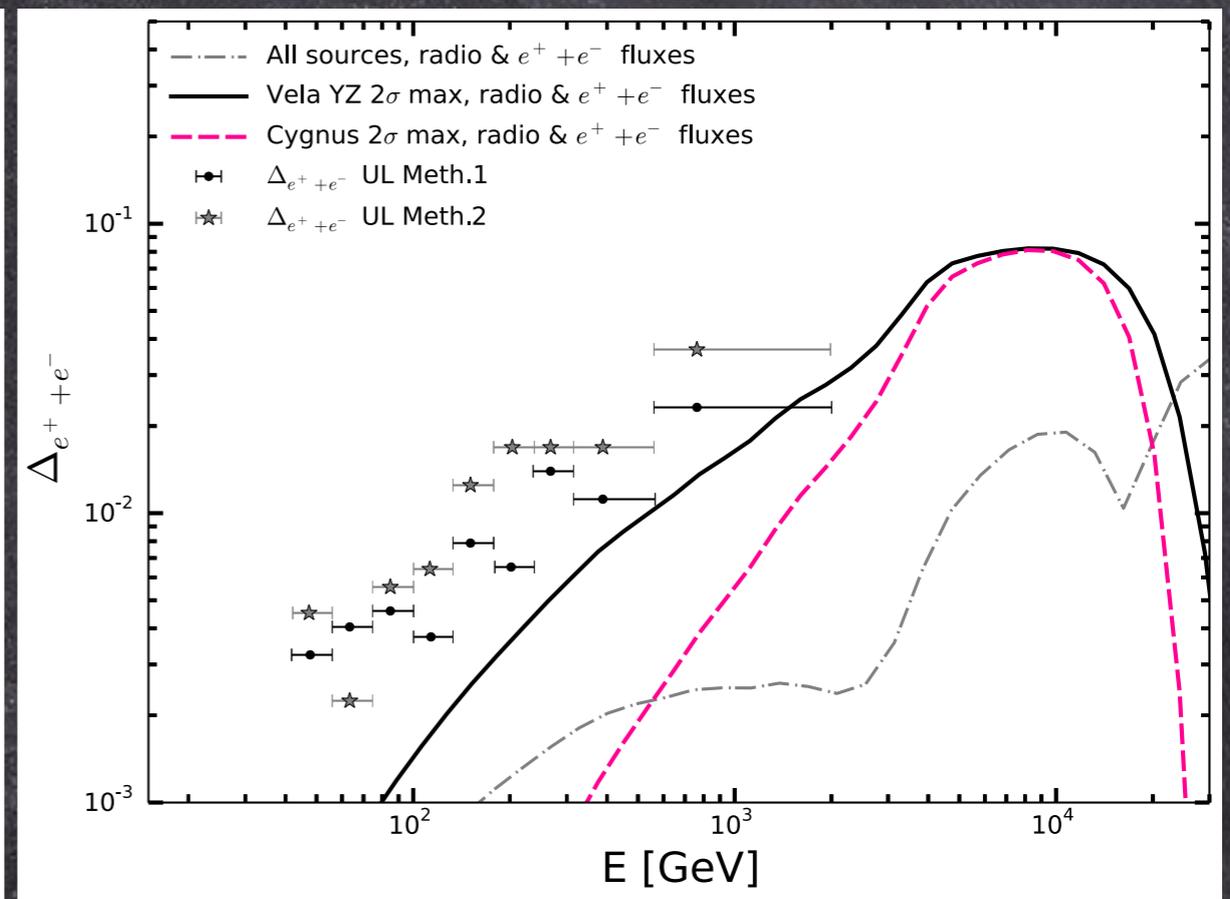
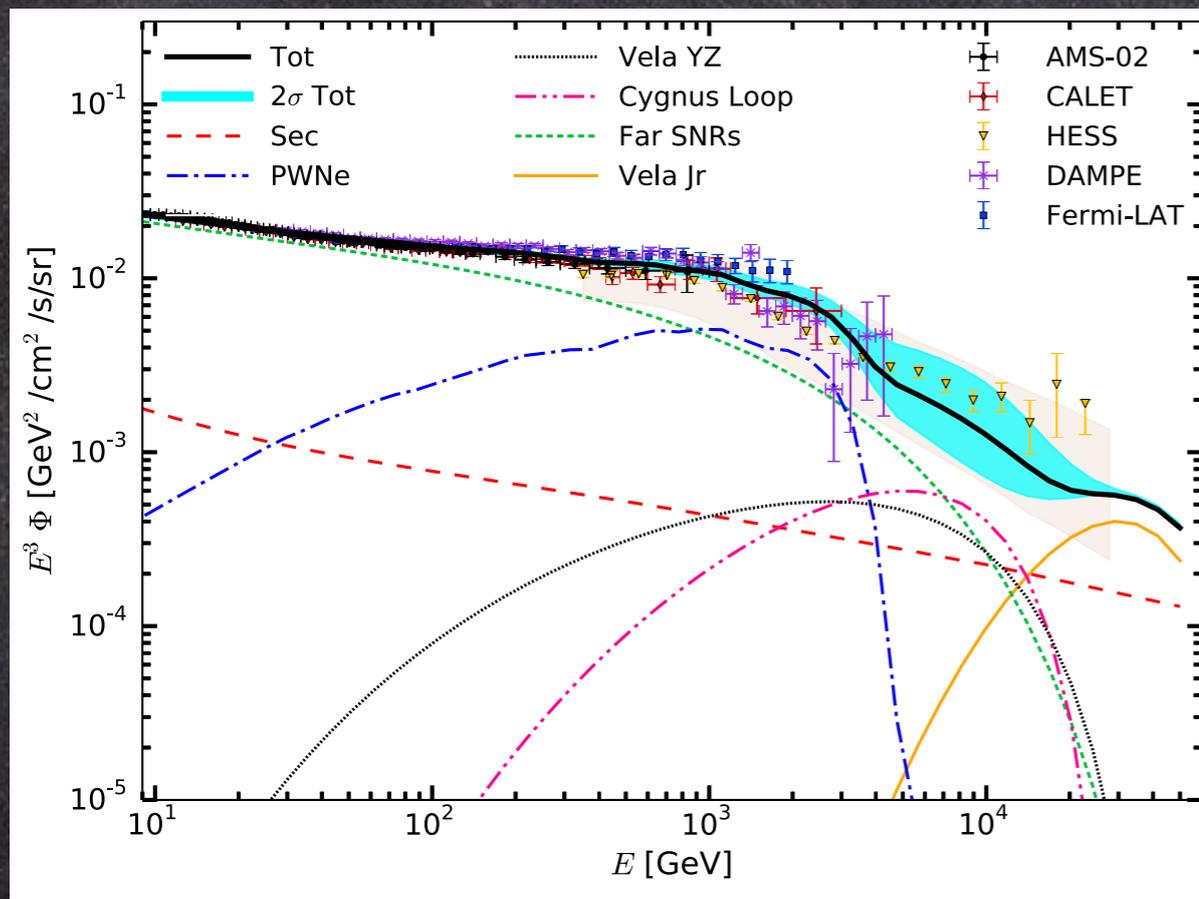
More generally, for a collection of sources (Shen & Mao, ApJL 1971):

$$\Delta(n_{\text{max}}, E) = \frac{1}{\psi^{\text{tot}}(E)} \cdot \sum_i \frac{\mathbf{r}_i \cdot \mathbf{n}_{\text{max}}}{\|\mathbf{r}_i\|} \cdot \psi_i(E) \Delta_i(E)$$

# A multi-wavelength & multi-messenger analysis

We fit the parameters selected by radio and  $e^+e^-$  flux data and check against dipole anisotropy data

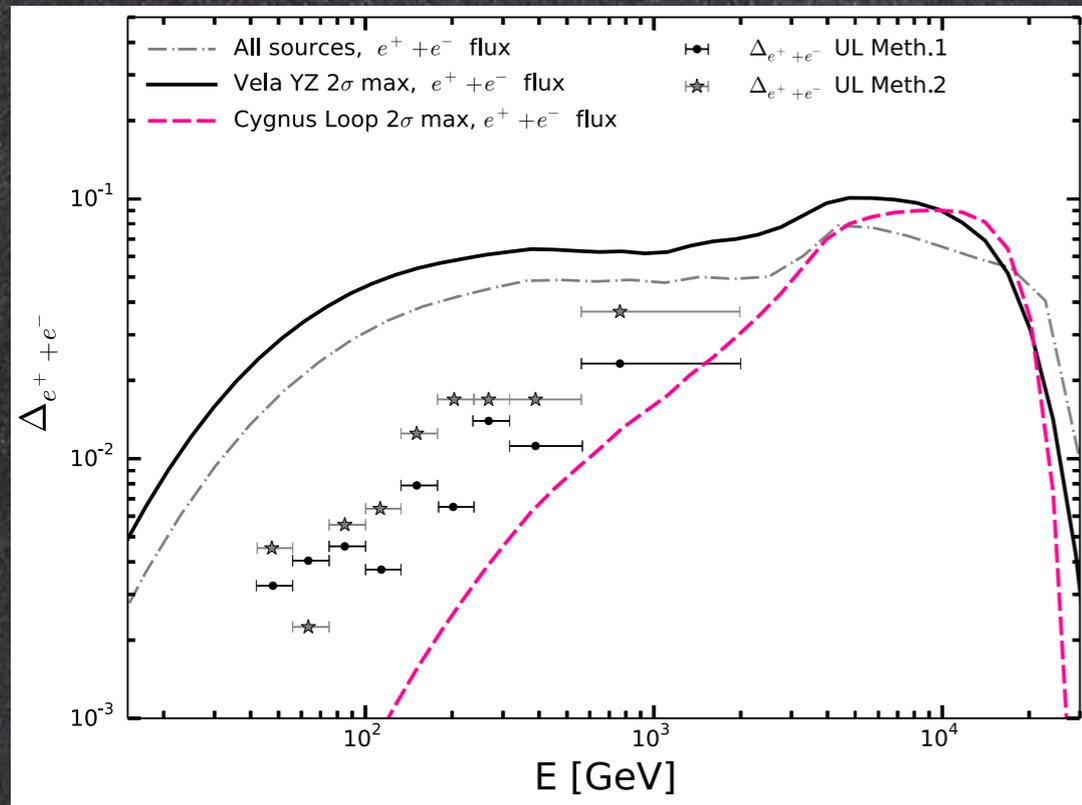
Manconi, Di Mauro, FD JCAP 2019



We find models compatible with three independent observables:  
radio flux from SNR;  $e^+e^-$  flux;  $e^+e^-$  dipole anisotropy

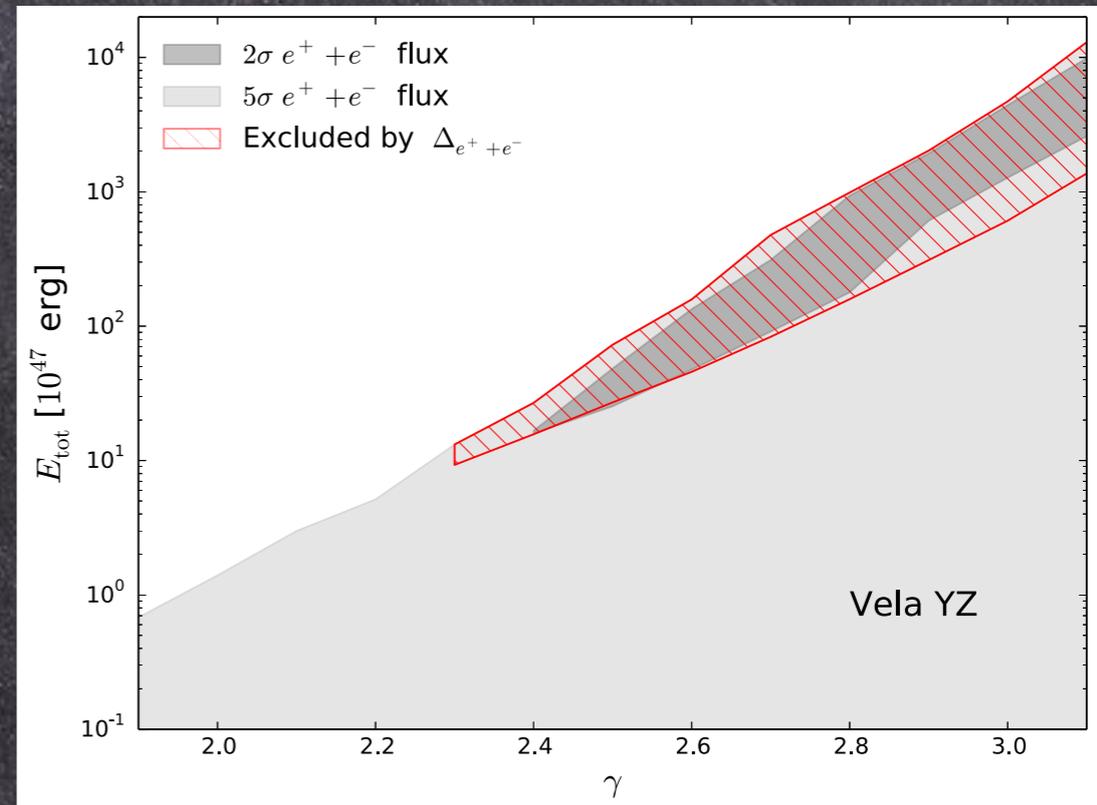
# Bounds from dipole anisotropy

Fermi-LAT data from  $e^+e^-$  dipole anisotropy are upper bounds vs E (Abdollahi+ PRL 2017)



Maximal anisotropy from  $e^+e^-$  flux  
selected configurations

Manconi, Di Mauro, FD JCAP 2019



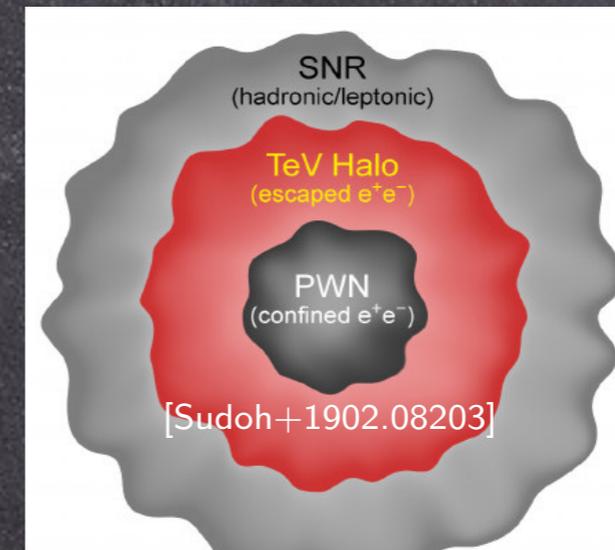
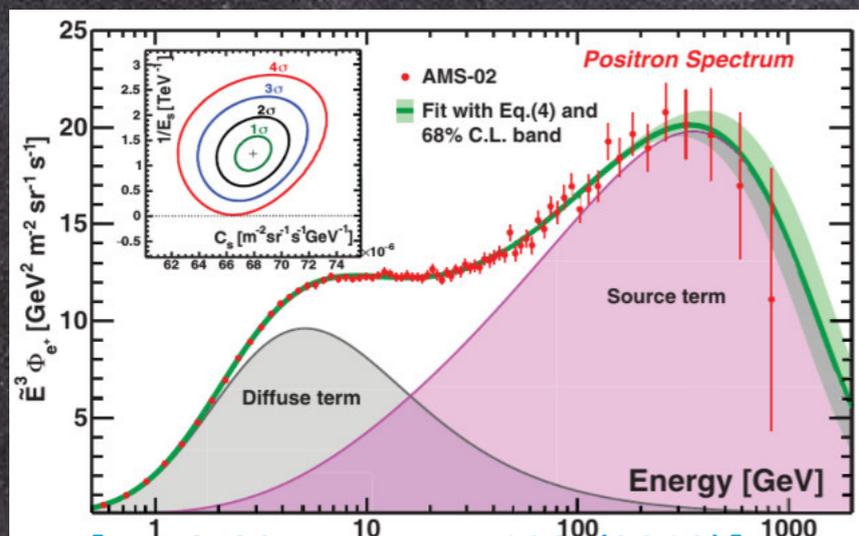
Anisotropy excludes configurations  
selected by  $e^+e^-$  flux

Dipole anisotropy in CR leptons is a valuable observable to  
study the properties of local sources - SNRs

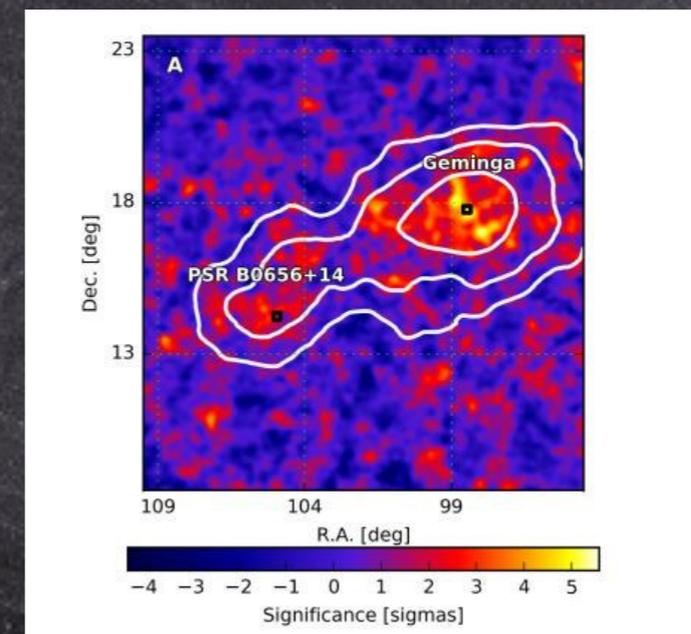
# Positron flux and $\gamma$ -ray haloes

- The positron flux shows (Pamela, AMS data) the need of primary source at high energies. **Pulsars could do the job.**

AMS Coll. PRL 2019s



- HAWC has detected a TeV gamma-ray halo around **Geminga** and **Monogem** pulsars. Interpreted as  $e^+e^-$  accelerated by the pulsar, then released in the ISM, A **low diffusion** region around the pulsars is favored by data



HAWC Collaboration, Science 358, 2017  
Milagro, Abdo et al., ApJL 2009

# What we learn from HAWC data

Di Mauro, Manconi, FD PRD 2020 Hooper+1702.08436, Fang+1803.02640  
Sudoh+1902.08203, Johannesson+1903.05509, Tang+1808.02445

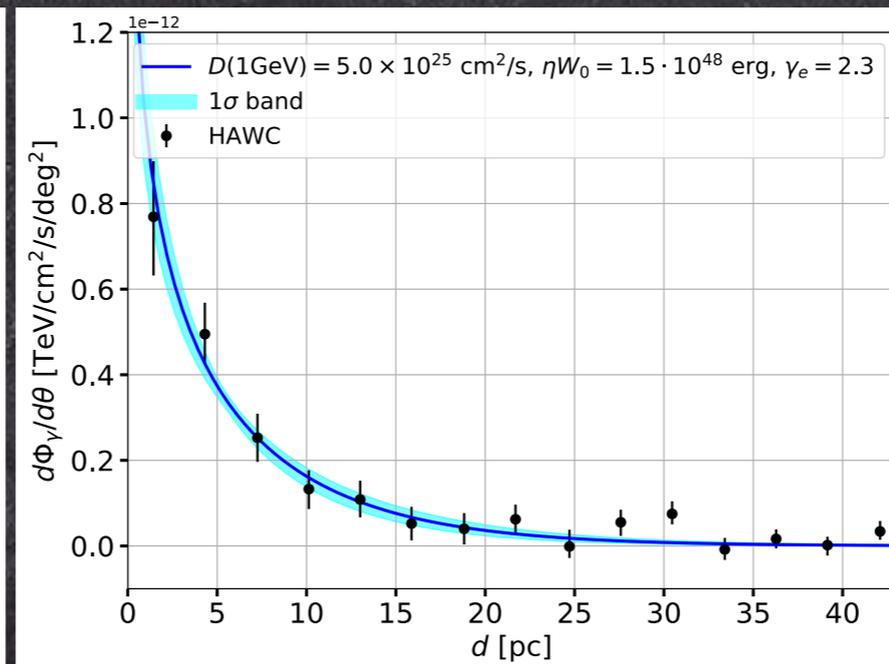
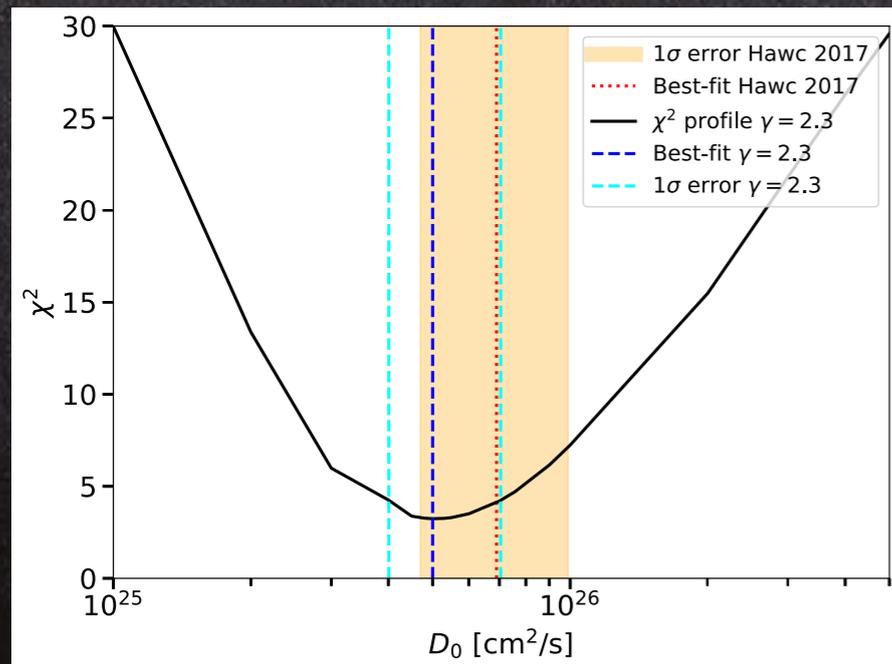
- The  $e^\pm$  injection is continuous (not burst-like)

$$Q(E, t) = L(t) \left( \frac{E}{E_0} \right)^{-\gamma} \exp \left( -\frac{E}{E_c} \right)$$

$$L(t) = \frac{L_0}{\left( 1 + \frac{t}{\tau_0} \right)^2}$$

- The spin-down luminosity converted into HAWC high energy  $e^\pm$  is  $\eta W_0 = 1.5 \cdot 10^{48}$  ( $4.2 \cdot 10^{46}$ ) erg for Geminga (Monogem)

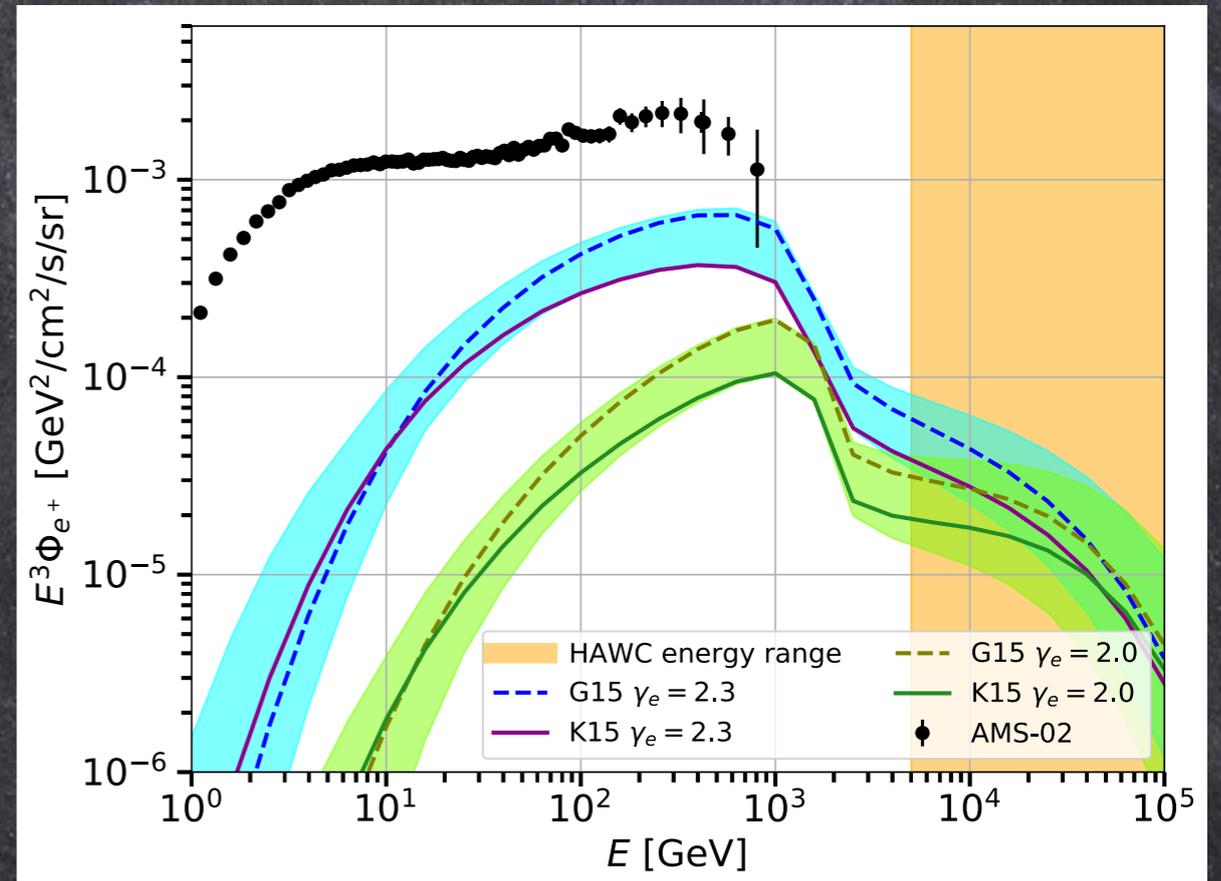
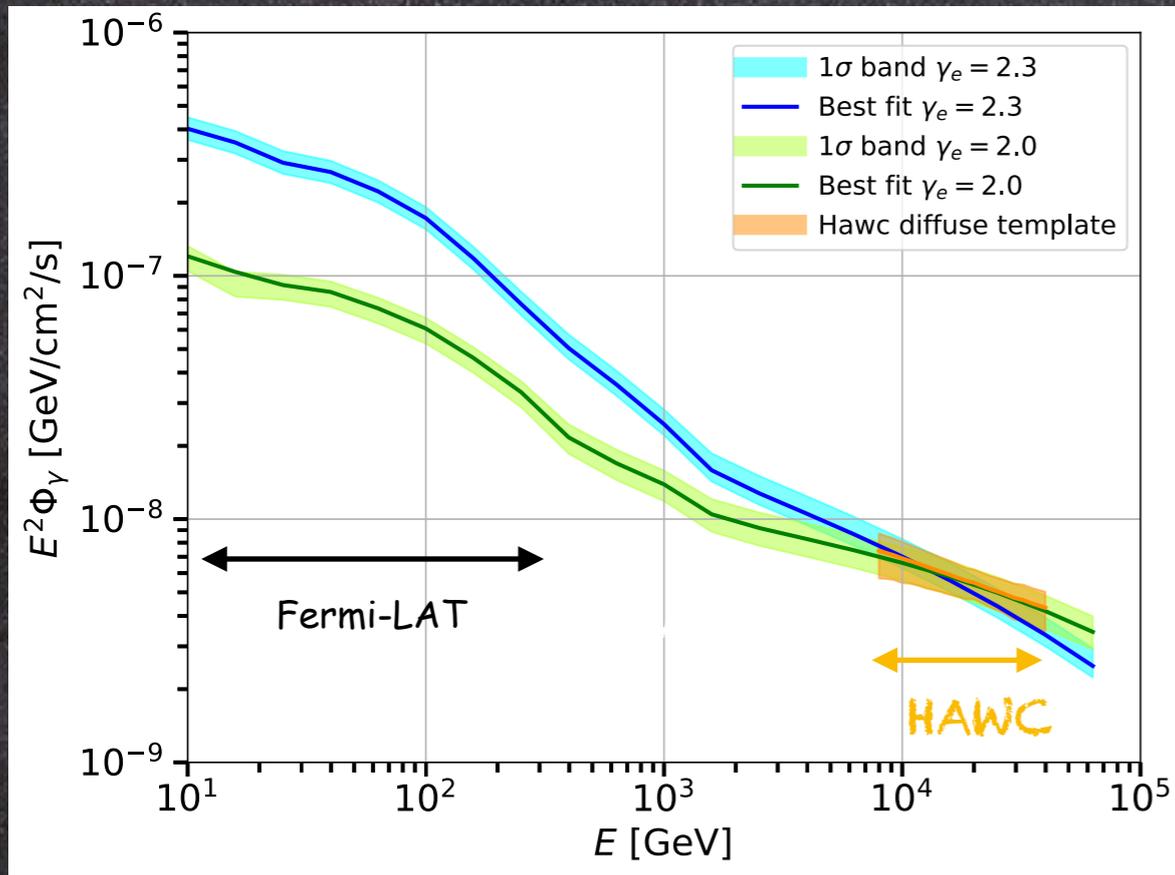
- The diffusion is inhibited around the pulsar by  $\sim 500$  times wrt the average in the ISM:  $D_0(1 \text{ GeV}) = 5 \cdot 10^{25} \text{ cm}^2/\text{s}$



Fit to Geminga  
surface brightness

# Which $e^\pm$ produce HAWC photons?

Di Mauro, Manconi, FD PRD 2020



The fit to HAWC surface brightness comes with uncertainties / degeneracies

The extrapolation down to Fermi-LAT energies gives remarkable differences

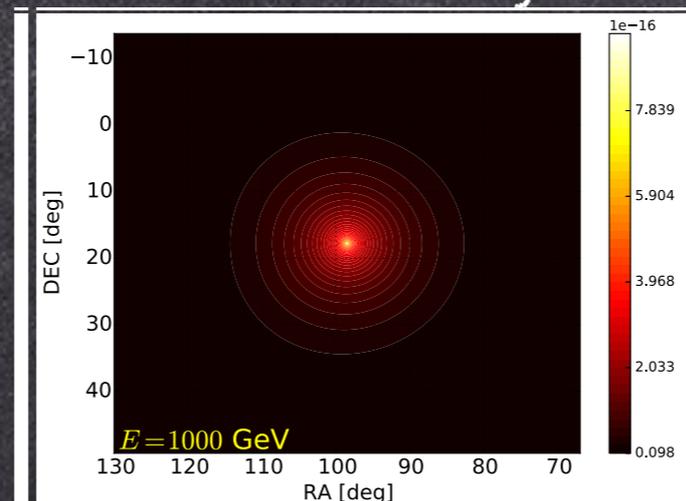
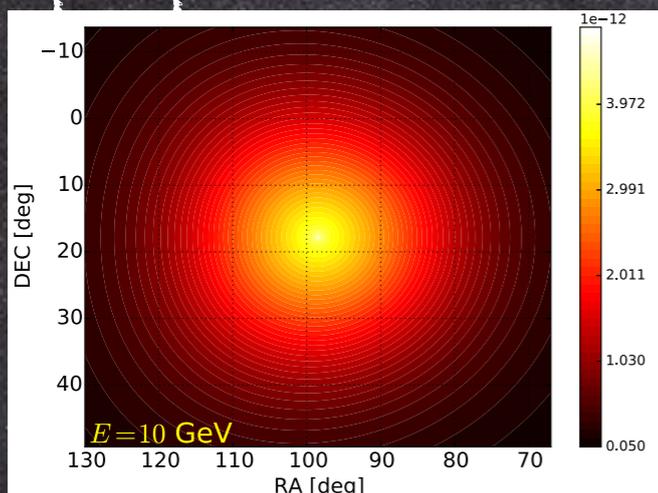
The HAWC data do not constrain the  $e^\pm$  measured by experiments (AMS02)

The  $e^\pm$  injection power spectrum is one key parameter

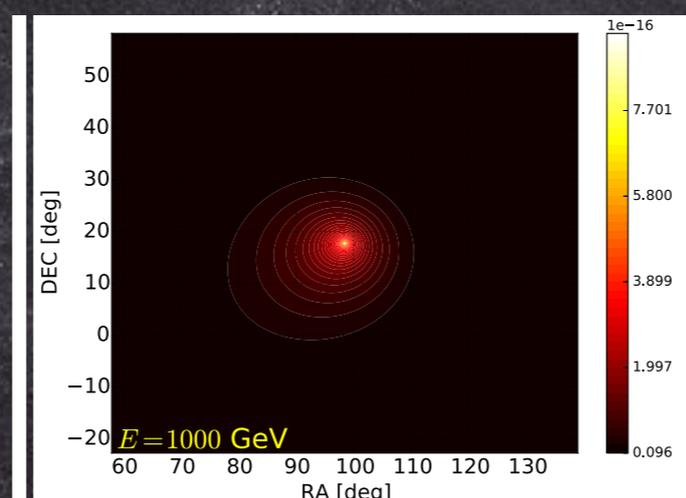
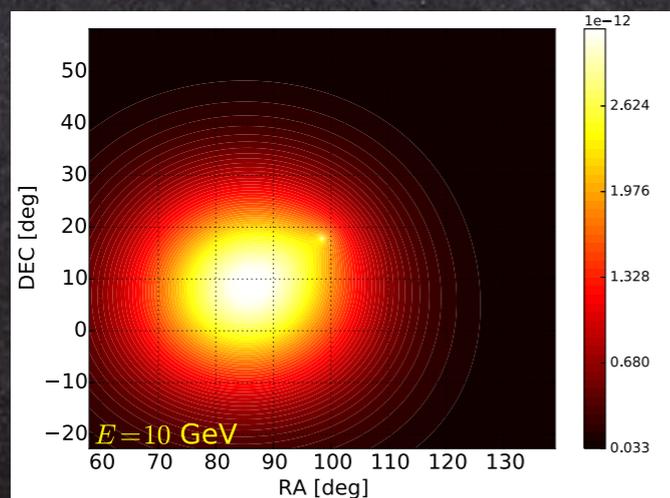
# Searching for $\gamma$ -rays around Geminga in the Fermi-LAT data

Di Mauro, Manconi, FD PRD 2020

- We implement a **Inverse Compton Scattering** template with background Interstellar radiation field (ISRF, needs a model)
- Pulsar proper motion:  $v_T = 211$  km/s (Faherty+AS2007) (70 pc travelled)



Without proper motion



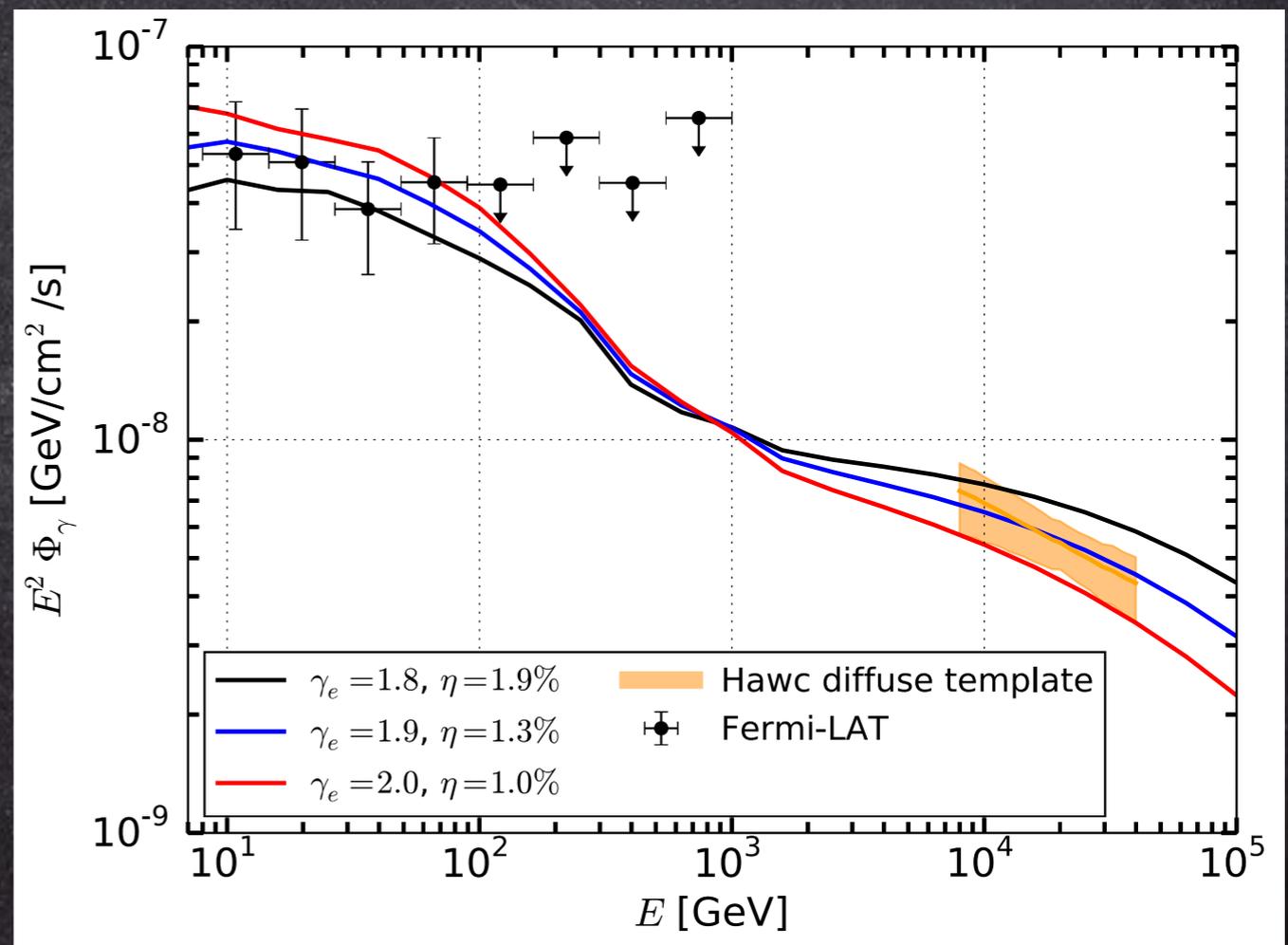
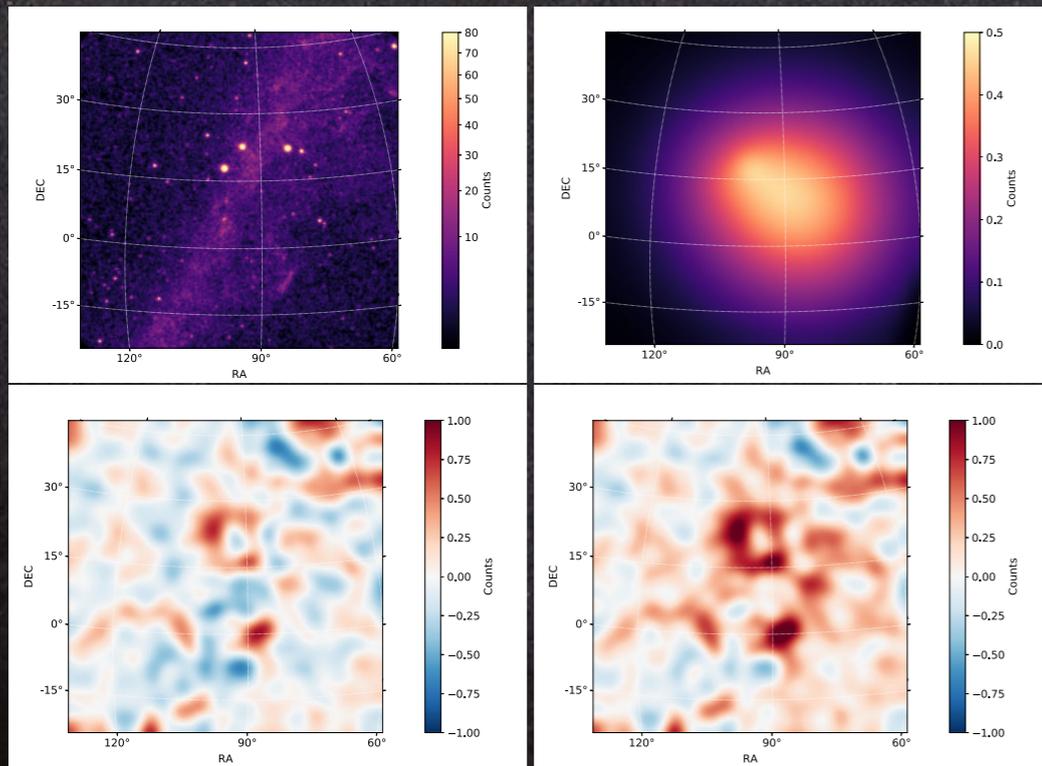
With proper motion

# Detection of a $\gamma$ -ray halo in Fermi-LAT data around Geminga

We detect a  $\gamma$ -ray halo around Geminga at  $7.8-11.8\sigma$   
(depending on background models)

- Fit improves with proper motion included.
- Diffusion  $D(1\text{GeV}) = 1.6-3.5 \cdot 10^{26} \text{ cm}^2/\text{s}$  (compatible w/ HAWK)
- Extension  $\sim 60 \text{ pc}$  at  $100 \text{ GeV}$
- $\gamma_e = 1.8-2$

Di Mauro, Manconi, FD PRD 2020



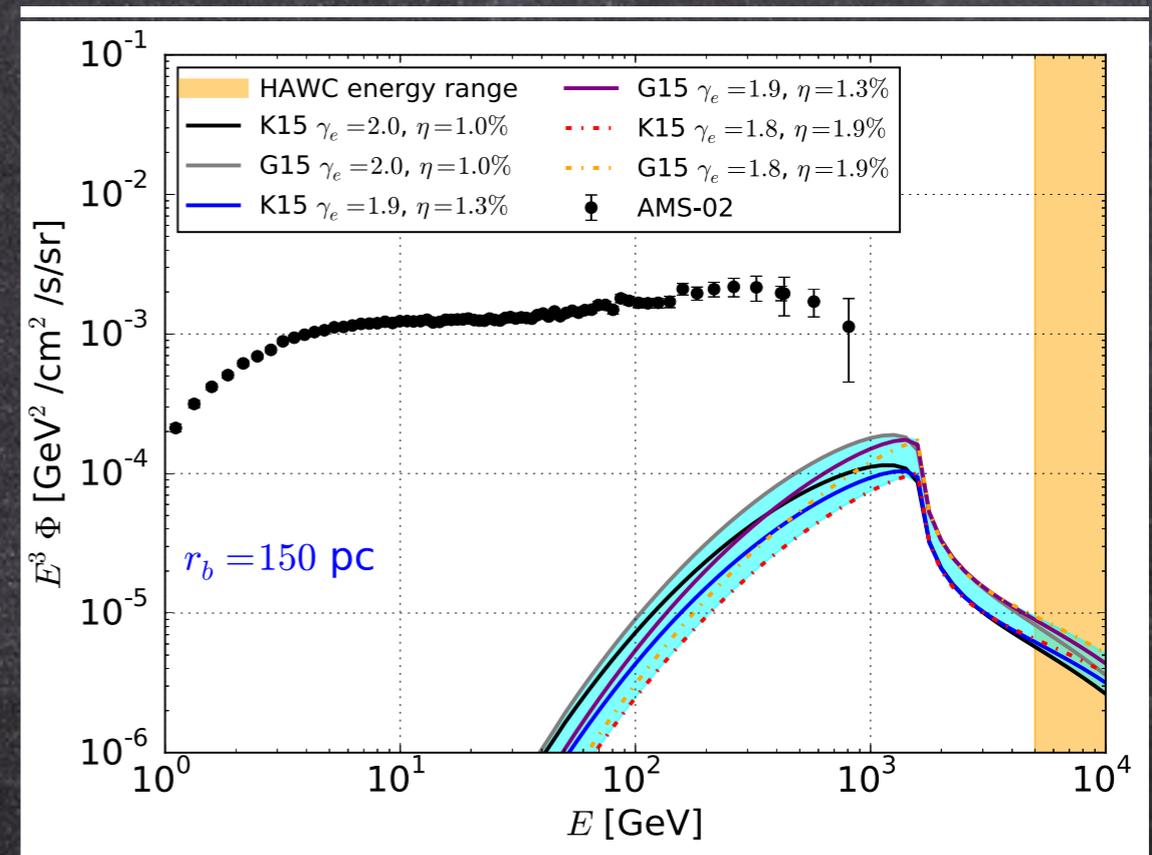
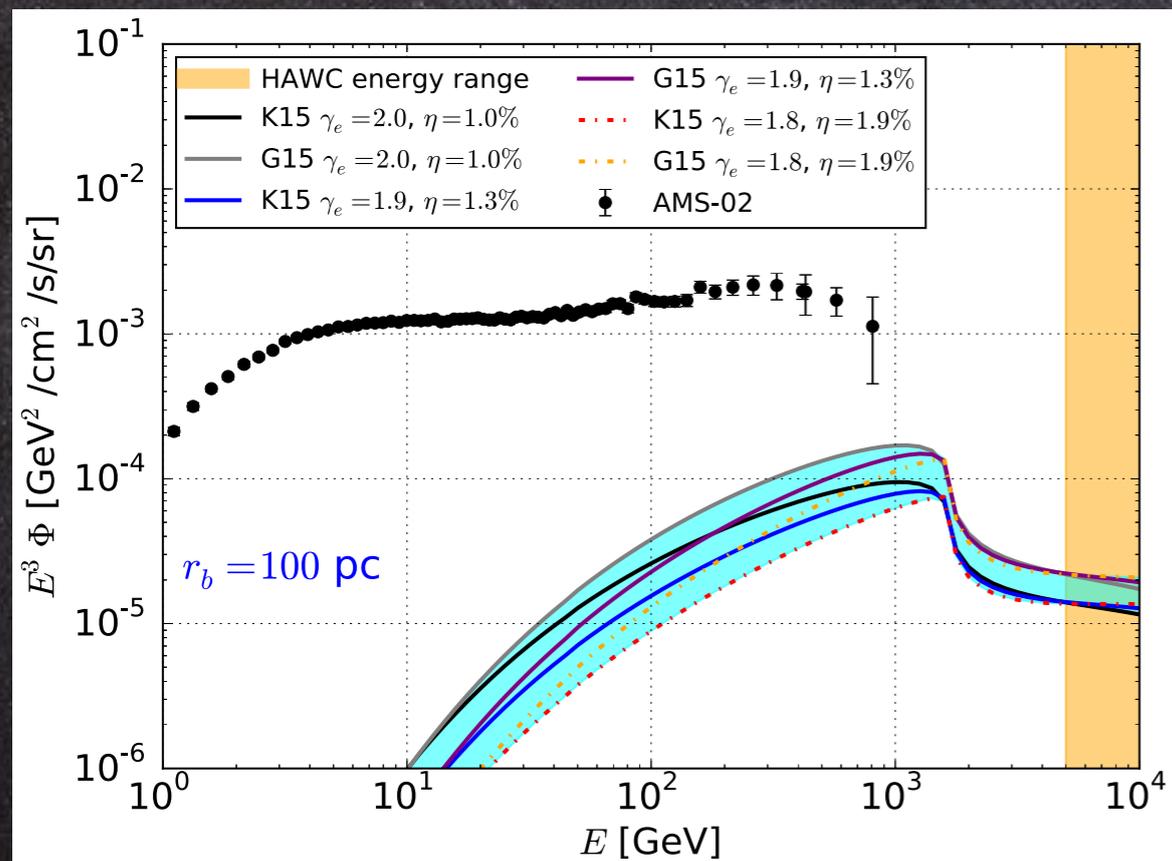
# $\gamma$ -ray haloes influence $e^+$ at the Earth

2-zones diffusion model:

$$D(r) = \begin{cases} D_0(E/1 \text{ GeV})^\delta & \text{for } 0 < r < r_b, \\ D_2(E/1 \text{ GeV})^\delta & \text{for } r \geq r_b, \end{cases}$$

$r_b$  is the boundary between low and high diffusion zones

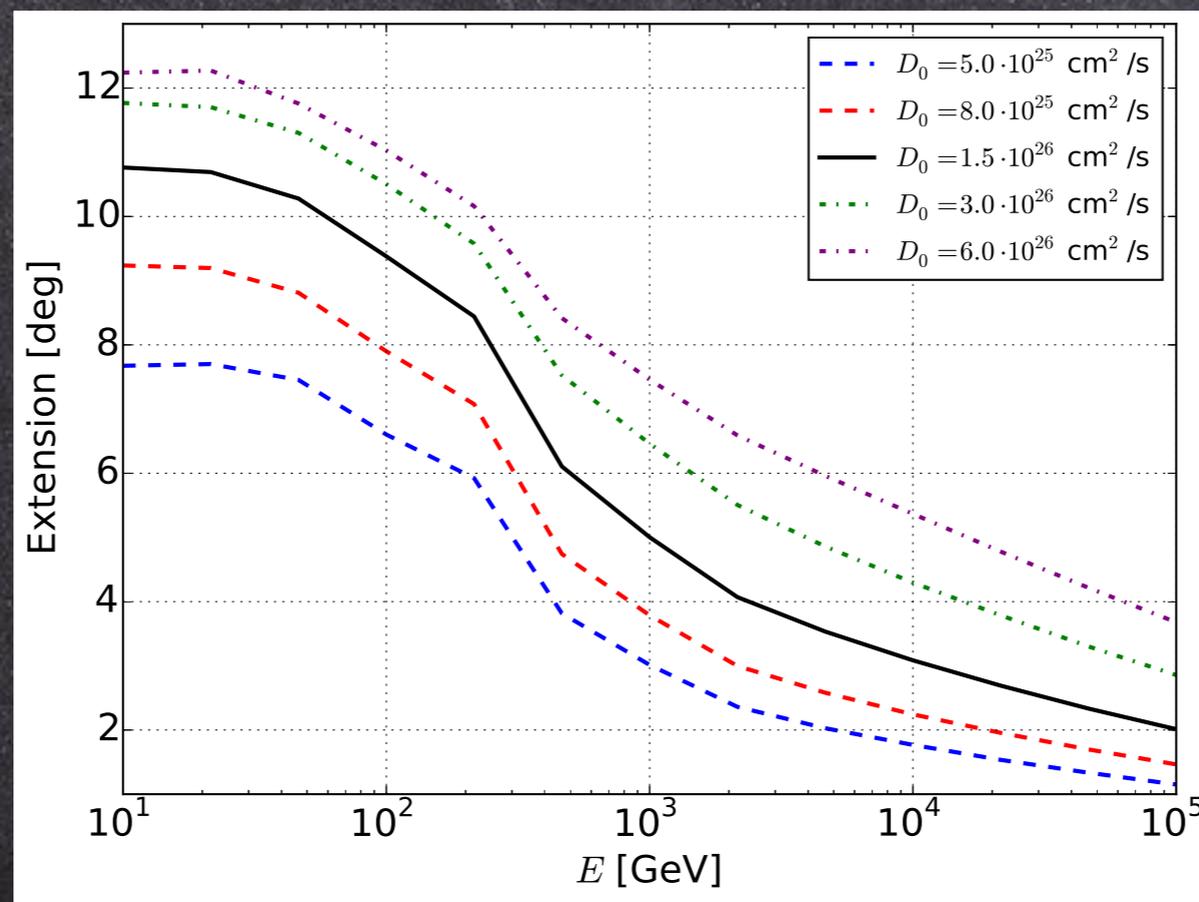
Di Mauro, Manconi, FD PRD 2020



Geminga contributes few % to the  $e^+$  flux at the Earth.  
The ICS halo is about 80 pc at Fermi-LAT energies

# Halo extension depends on energy and diffusion physics

Smaller haloes expected at higher energies



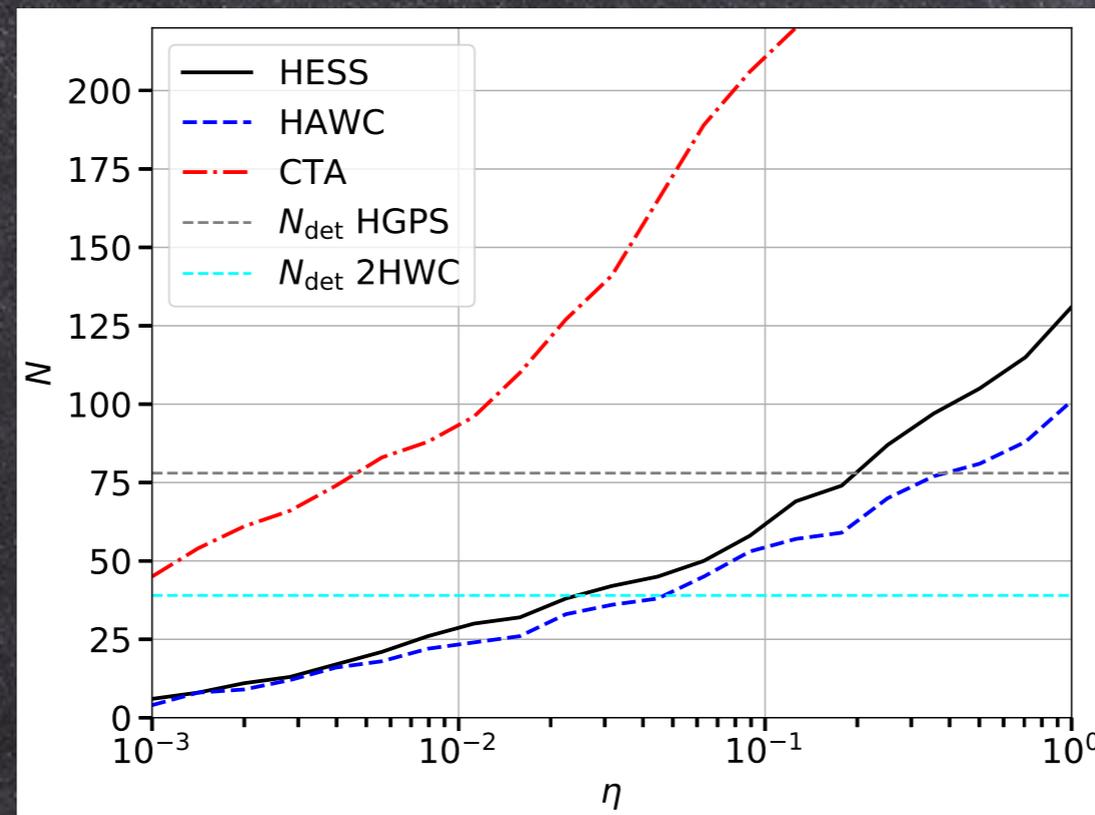
Higher diffusion coefficients get the halo spread out.  
ISM  $D(E)$  values ( $\sim 10^{28} \text{ cm}^2/\text{s}$ ) would get low energy  $\gamma$ -rays around Geminga spread widely in the ISM (no longer a halo...)

# $\gamma$ -ray haloes: a general property of pulsars?

M. Di MAuro, S. Manconi, FD arxiv:1908.03216

We select sources from ATNF catalog with highest ICS halo above 1 TeV

Compute the number of sources above HESS, HAWC and CTA sensitivity as a function of  $\eta$ , efficiency conversion into  $e^\pm$ .

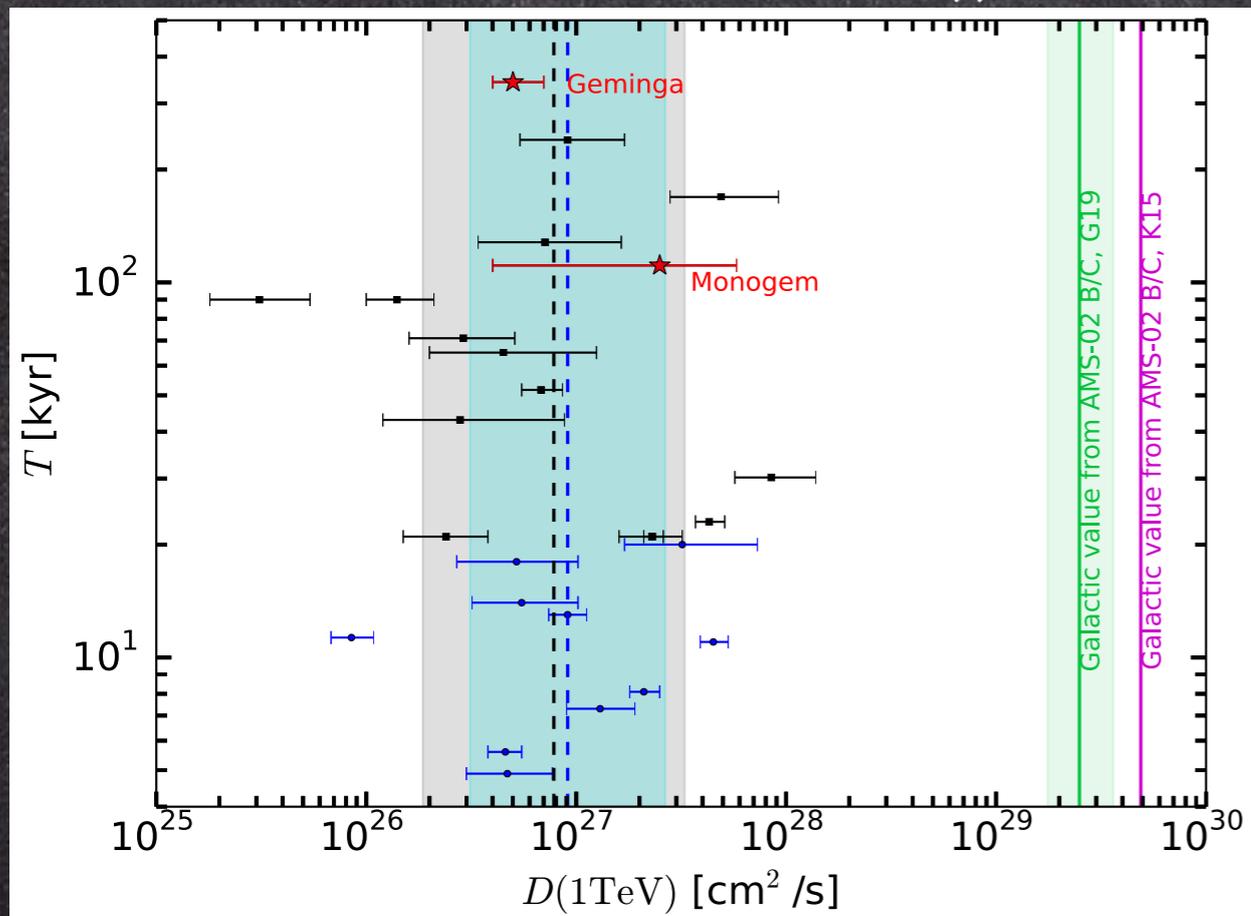


Tens of haloes could/will be detected even with 1% efficiency conversion into  $e^\pm$

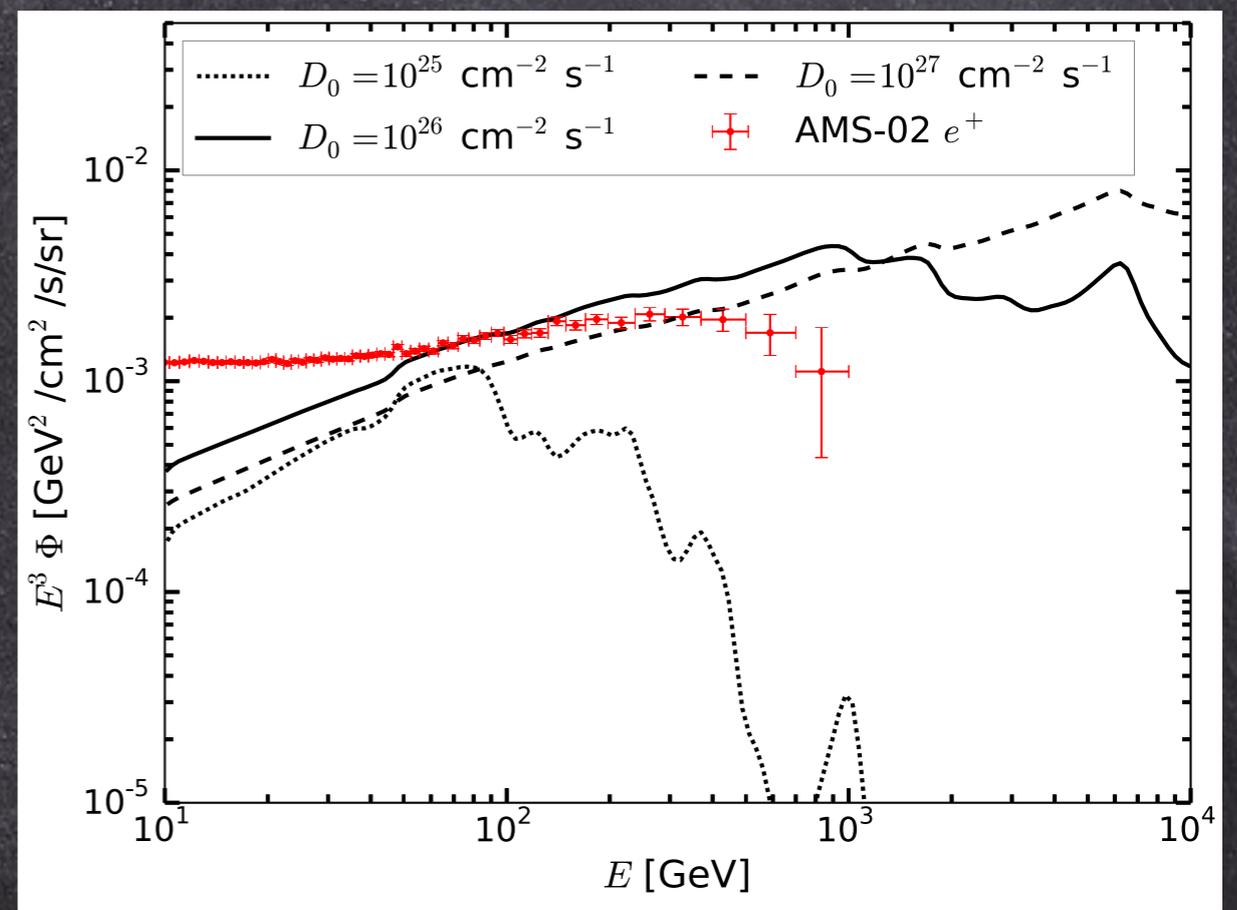
# A low diffusion zone around PWNe

We select sources detected mainly by HESS (they provide flux maps)  
Interpret the data in terms of ICS halo and fit  $D(1\text{TeV})$  and size

Di Mauro, Manconi, FD 1908.03216 to app. PRD



Manconi, Di Mauro, FD 2001.09985, subm. PRD



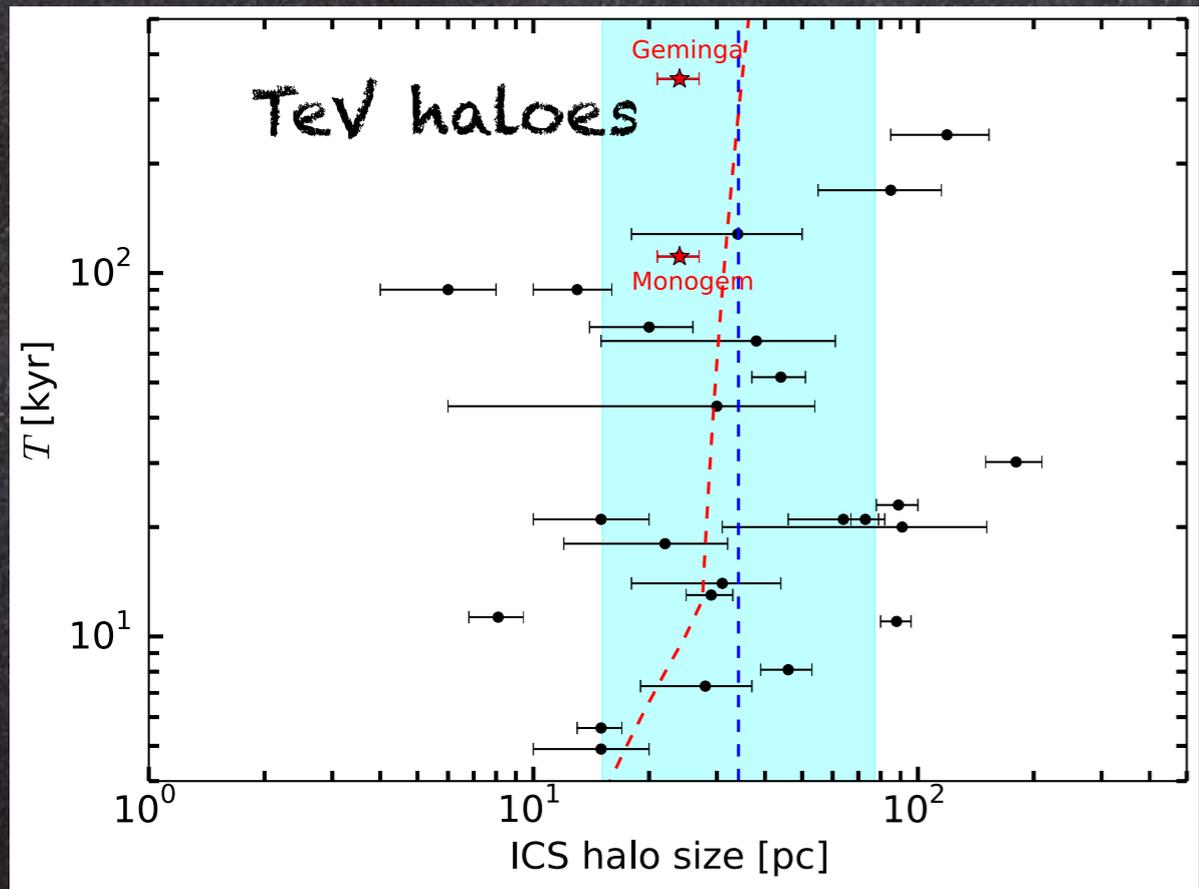
The diffusion coefficient around PWNe is systematically lower by 2 orders of magnitude w.r.t. the ISM diffusion coefficient found from CR data (B/C).

# $\gamma$ -ray haloes have tens pc size

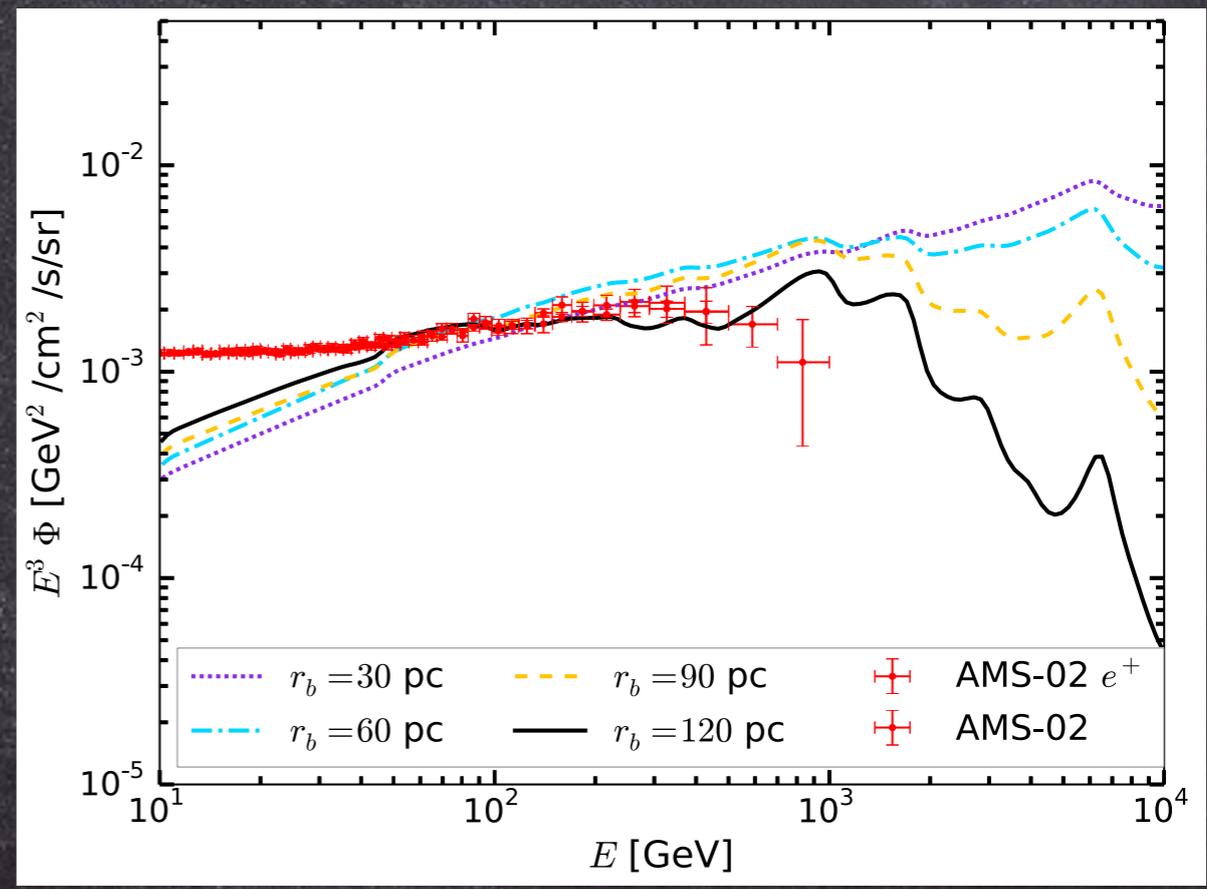
We fit also the ICS halo size

The trend with the age is compatible with models of PWN evolution  
The low diffusion zone around PWN should be larger than the halo size

Di Mauro, Manconi, FD 1908.03216 to app. PRD



Manconi, Di Mauro, FD 2001.09985, subm. PRD



The halo size can shape the high energy  $e^+$  flux

# ATNF pulsar catalog

Manconi, Di Mauro, FD 2001.09985, subm. PRD

Distance  $d$ , age  $T$  ( $50 \text{ kyr} < T < 10^5 \text{ kyr}$ ), spin-down energy  $dE/dt$

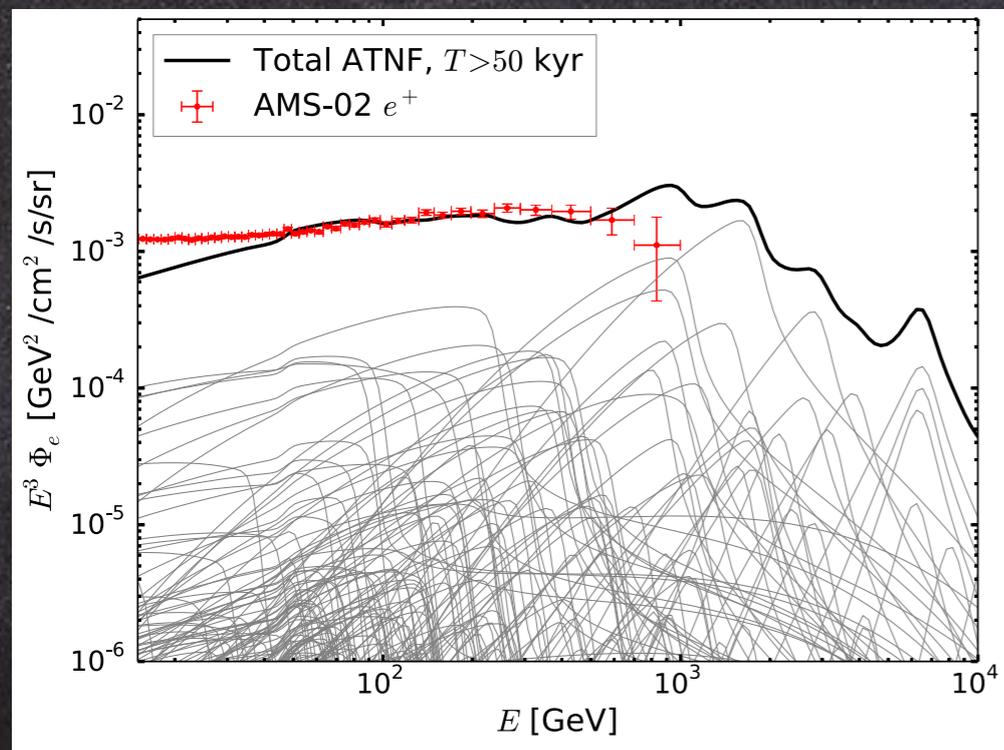
Continuous injection:

$$Q(E, t) = L(t) \left( \frac{E}{E_0} \right)^{-\gamma_e} \exp\left(-\frac{E}{E_c}\right) \quad L(t) = \frac{L_0}{\left(1 + \frac{t}{\tau_0}\right)^{\frac{k+1}{k-1}}}$$

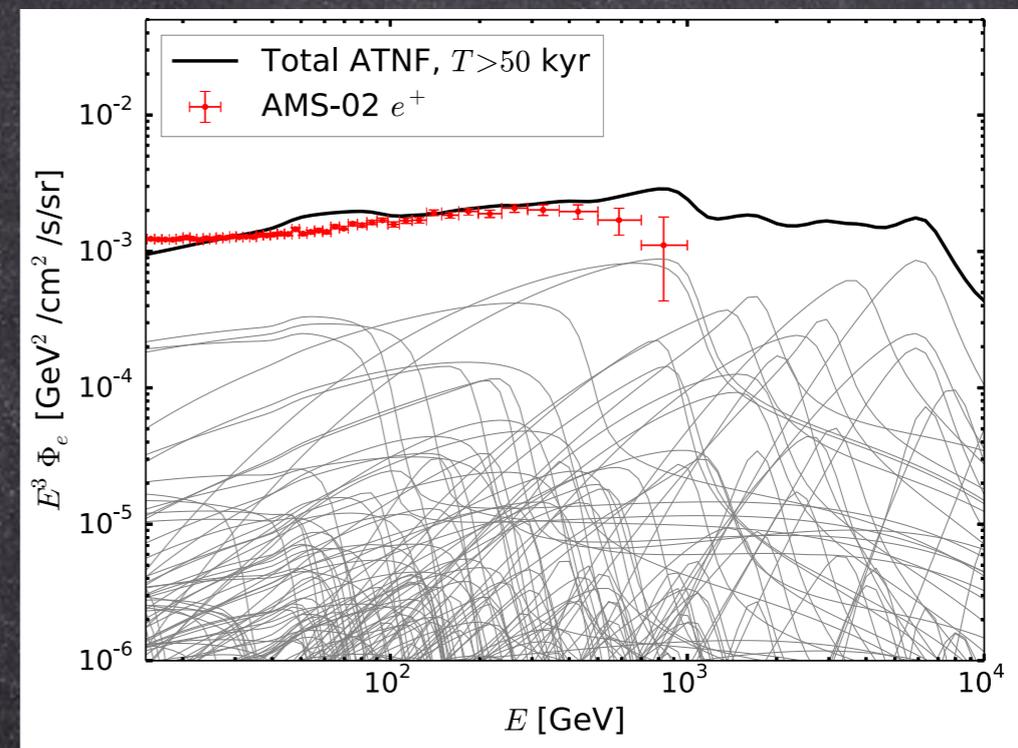
From catalog data:

$$W_0 = \tau_0 \dot{E} \left(1 + \frac{T}{\tau_0}\right)^{\frac{k+1}{k-1}} \quad E_{\text{tot}} = \eta W_0 = \int_0^T dt \int_{E_1}^{\infty} dE E Q(E, t)$$

$r_b = 120 \text{ pc}$ ,  $D_0 = 7.8 \cdot 10^{25} \text{ cm}^2/\text{s}$ ,  $\eta = 0.12$ ,  $\gamma_e = 1.9$



$r_b = 90 \text{ pc}$ ,  $D_0 = 7.8 \cdot 10^{25} \text{ cm}^2/\text{s}$ ,  $\eta = [0.02, 0.30]$ ,  $\gamma_e = [1.4, 2.0]$



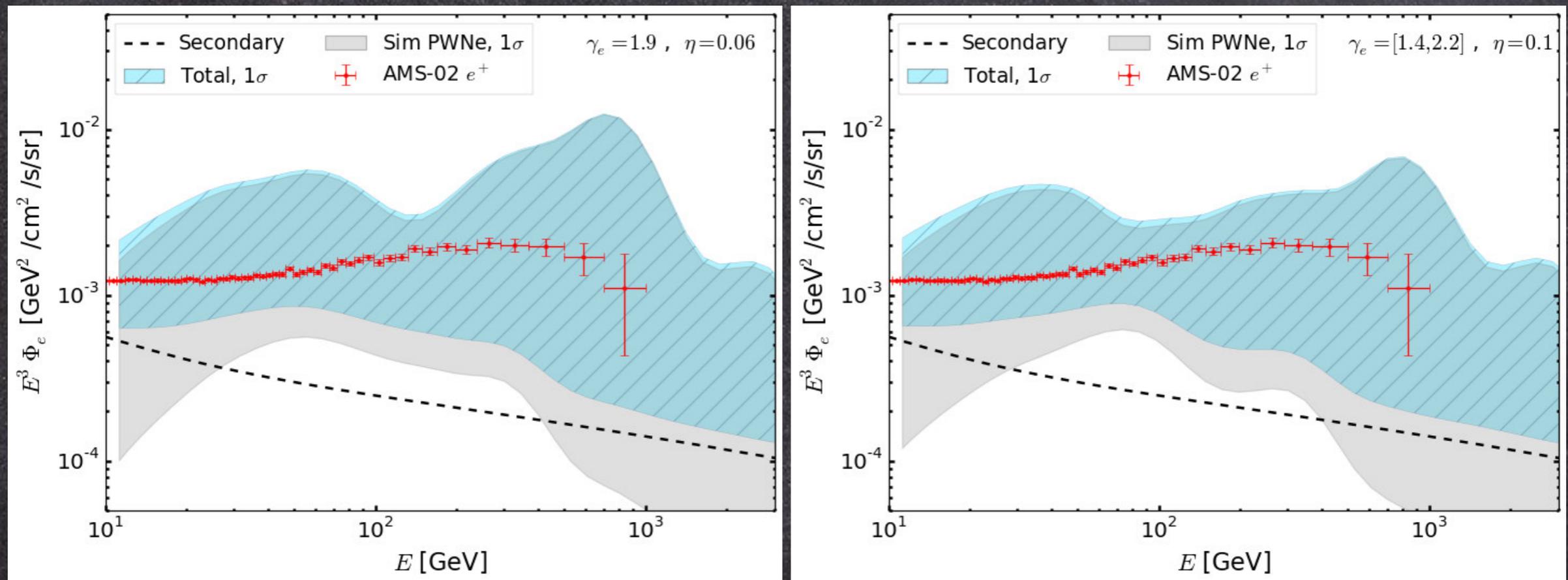
Few sources above 10% data - Cumulative flux  $\sim$  AMS-02 data  
The h.e. trend dictated by low diffusion ( $D_0$ ) within  $r_b$

# Simulations of galactic pulsar population

Manconi, Di Mauro, FD 2001.09985, subm. PRD

The ATNF catalog may be incomplete

We simulate a galactic Pulsar population with physics inputs:



The uncertainty band is  $\sim 10$

The flux is NEVER negligible up to TeV

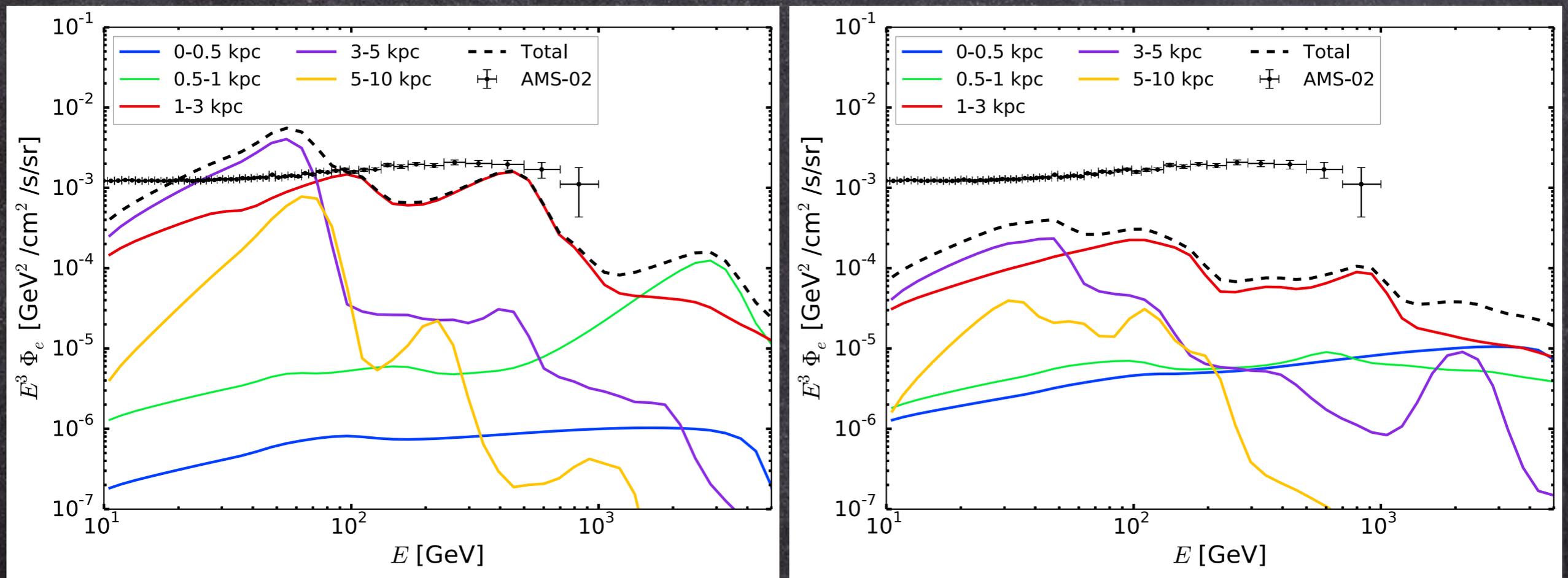
Predictions fall around AMS-02 data

# Conclusions

- Leptons at Earth have a composite origin:  $e^-$  from far smooth and near catalog SNR,  $e^+e^-$  from PWN,  $e^+e^-$  as secondaries in the ISM
- $e^+$  (with  $p^-$ ) are a major antimatter CR component in the Galaxy
- The violent radiative cooling provides invaluable and different messengers
- We discovered an ICS  $\gamma$ -ray halo around Geminga pulsar (U.L. for Monogem) in the Fermi-LAT data at 8-100 GeV
- CRs at high energies, dipole anisotropy, radiation from radio to  $\gamma$  rays are giving us an innovative view of the Galaxy and an explanation of the CR  $e^+$  data on the whole energy range

# Contribution of pulsars to $e^+$ flux as a function of distance from Earth

Manconi, Di Mauro, FD 2001.09985, subm. PRD



# Burst-like or continuous injection

Di Mauro, Manconi, FD PRD 2020

