

# Neutron Stars in the laboratory: how we can test the hypothesis of Hyperon Stars

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E62 - Dense and Strange Matter

<https://www.denseandstrange.ph.tum.de>

- (strange) Hadron interactions
- Hyperon Puzzle in Neutron Stars
- The measurement of Hadron Hadron Correlations
- Experimental Results: RUN1 and RUN2
  - pp Collisions at 7 TeV, 5 TeV and 13 TeV, p-Pb at 5.02 TeV measured by ALICE
  - pp, p $\Lambda$ ,  $\Lambda\Lambda$ , p $\Xi^-$  Correlations
- Outlook

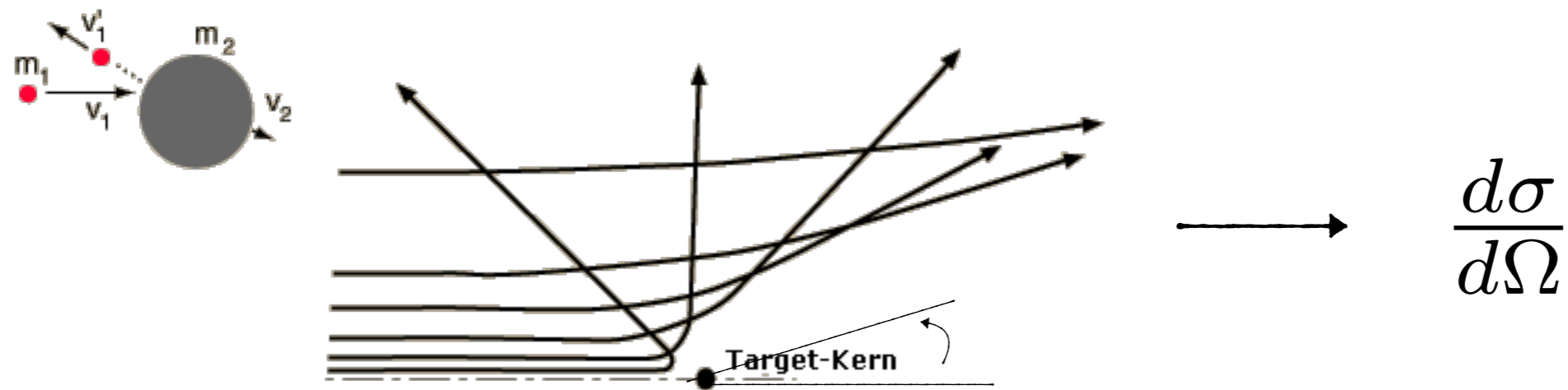
# (Strange) Hadron Interactions

$$\Lambda = uds \quad I(J^P) = 0\left(\frac{1}{2}^+\right)$$

$$\Sigma^0 = uds \quad I(J^P) = 1\left(\frac{1}{2}^+\right)$$

$$\Xi^- = dss \quad I(J^P) = \frac{1}{2}\left(\frac{1}{2}^+\right)$$

Scattering experiments -> Extraction of the differential cross section

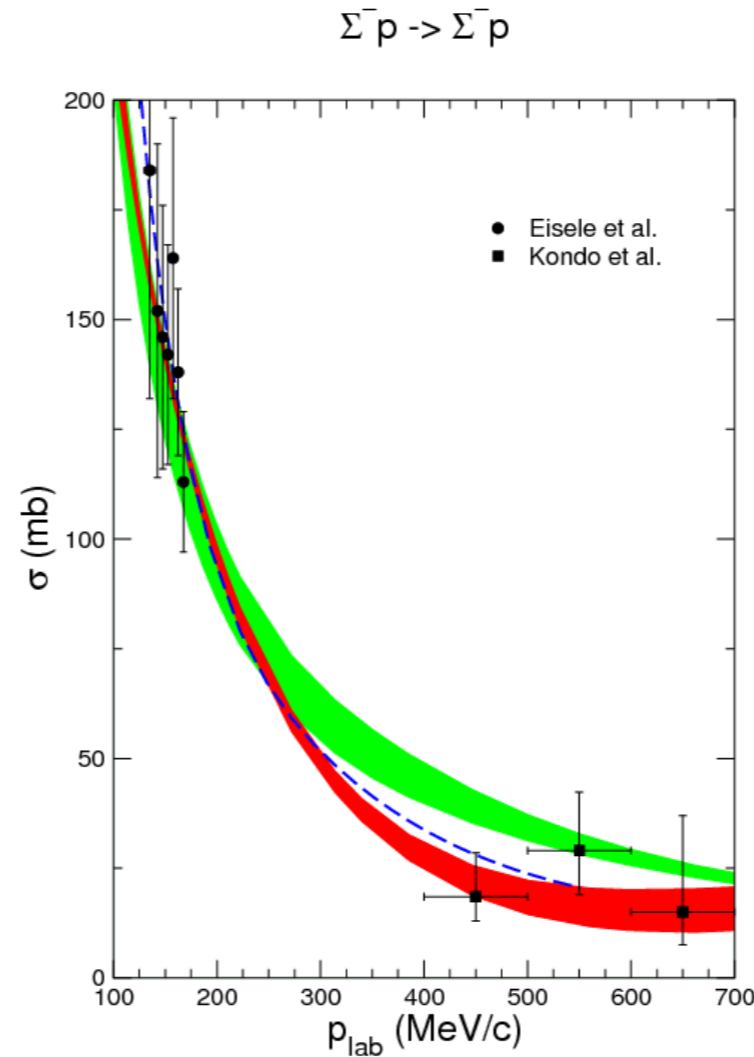
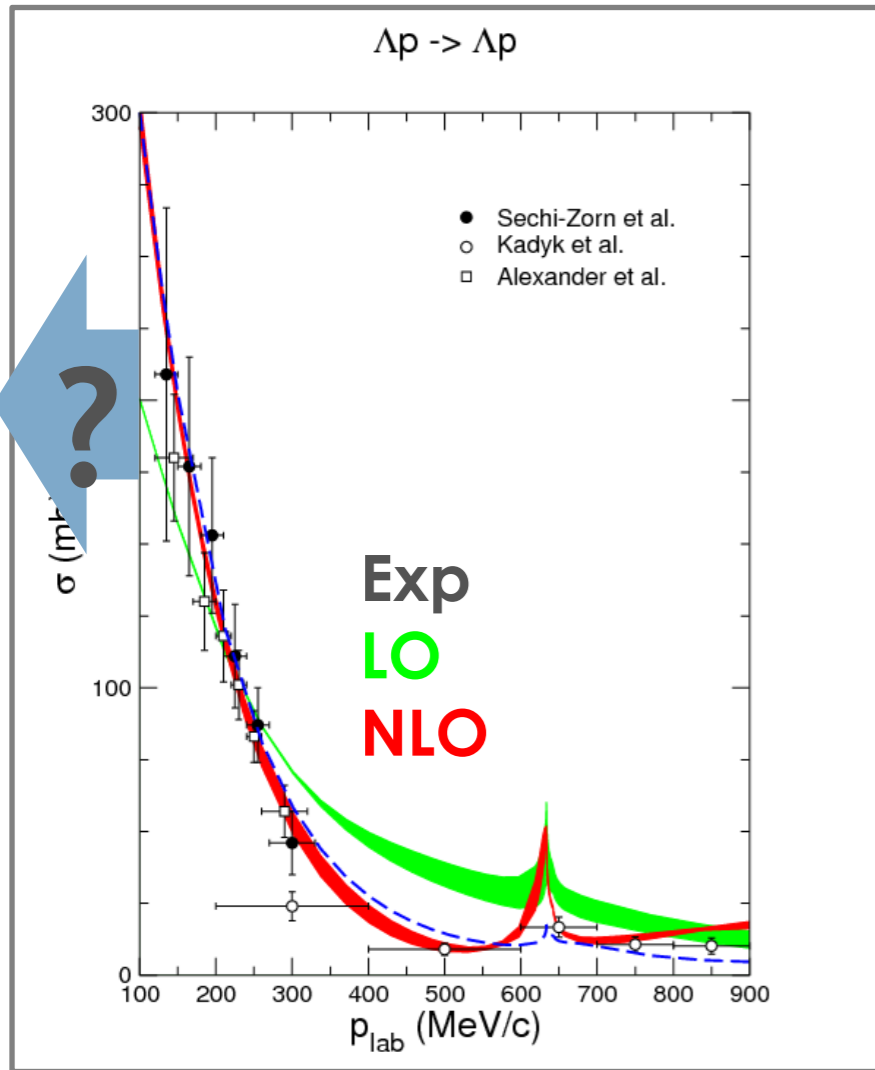


**Partial Wave Expansion:**

$$\sigma = \frac{4\pi}{k^2} \sum_l (2l + 1) \sin^2(\delta_l). \quad \delta_l = \text{phase shifts}$$

**Scattering Length**

$$f_0 = - \lim_{k \rightarrow 0} \frac{1}{k} \tan \delta_0(k) \quad l=0, \text{ s-wave Only!}$$



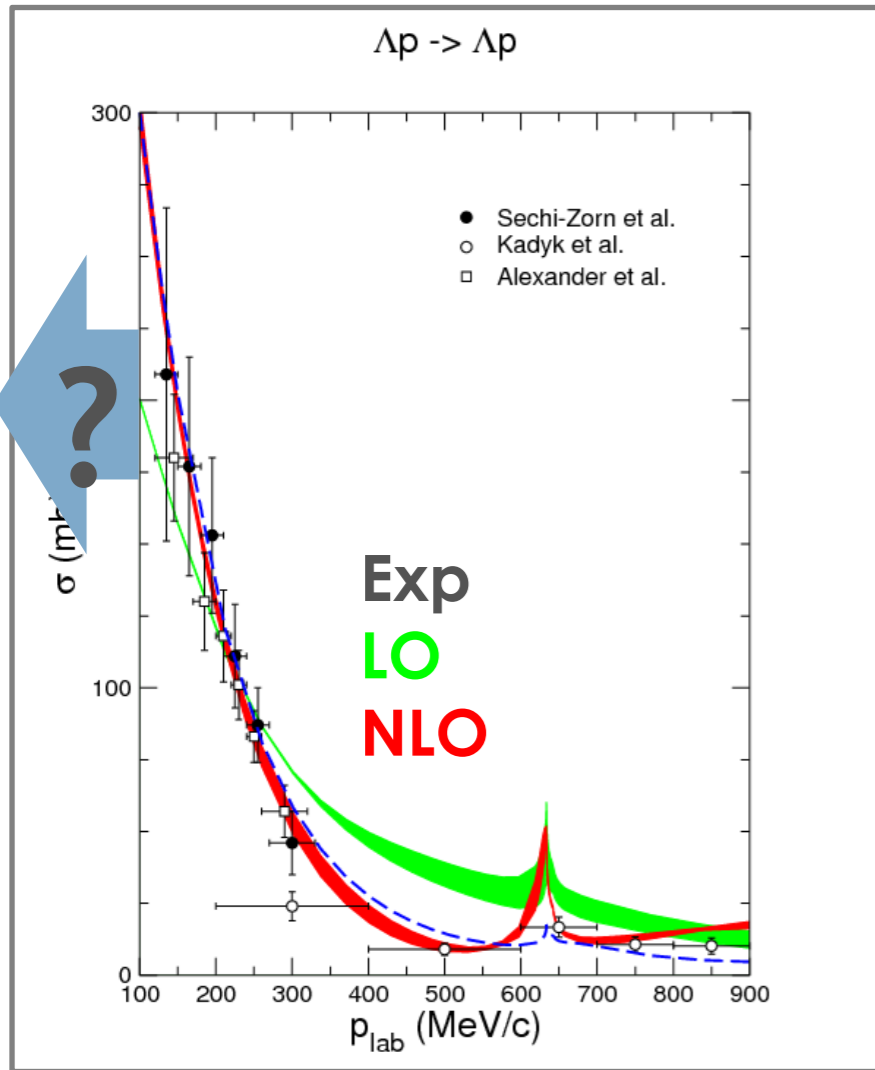
LO: H. Polinder, J.H., U. Meißner, NPA 779 (2006) 244  
 NLO: J.Haidenbauer., N.Kaiser, et al., NPA 915 (2013) 24

Data from scattering experiments and bubble chambers detectors from 1968 and 1971

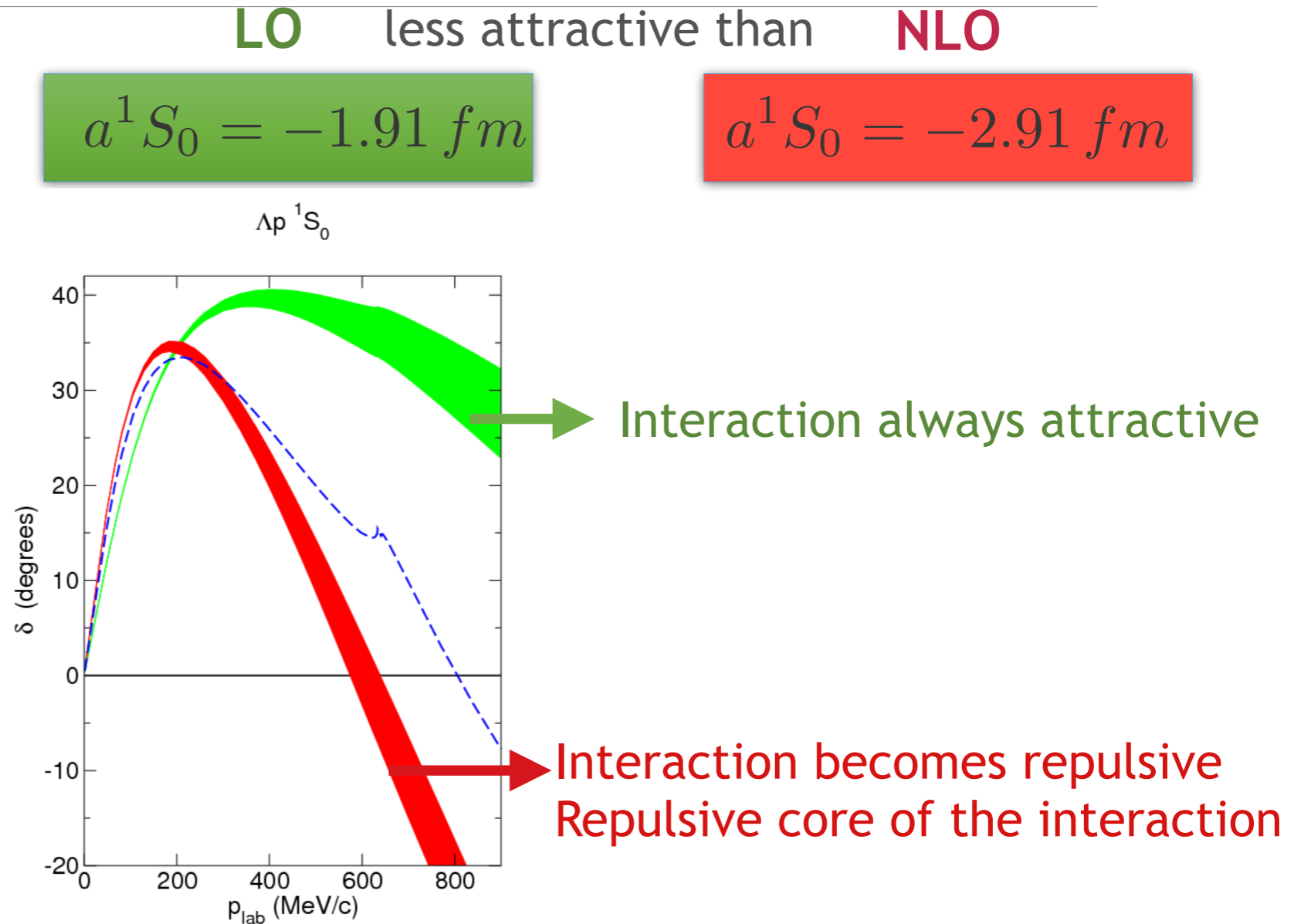
$$K^- + p \rightarrow \Sigma^0 + \pi^0, \Sigma^0 \rightarrow \Lambda + \gamma$$

$$K^- + p \rightarrow \Sigma^- + \pi^+ \dots$$

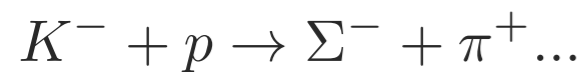
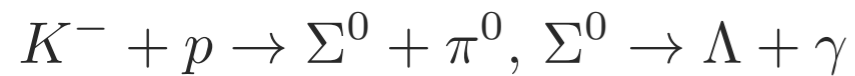
Production Threshold for  $\Lambda'$ s :  $p \geq 100 \text{ MeV}$



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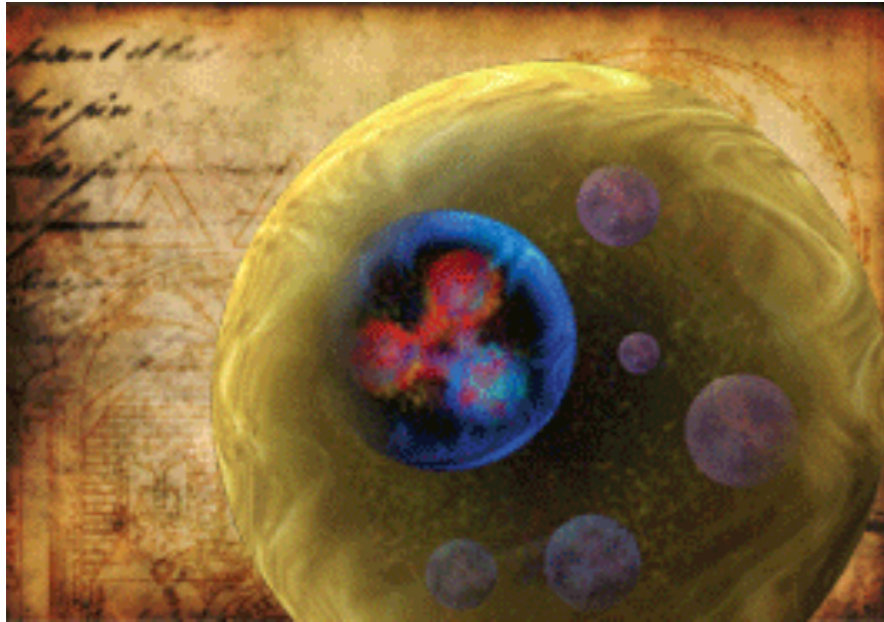


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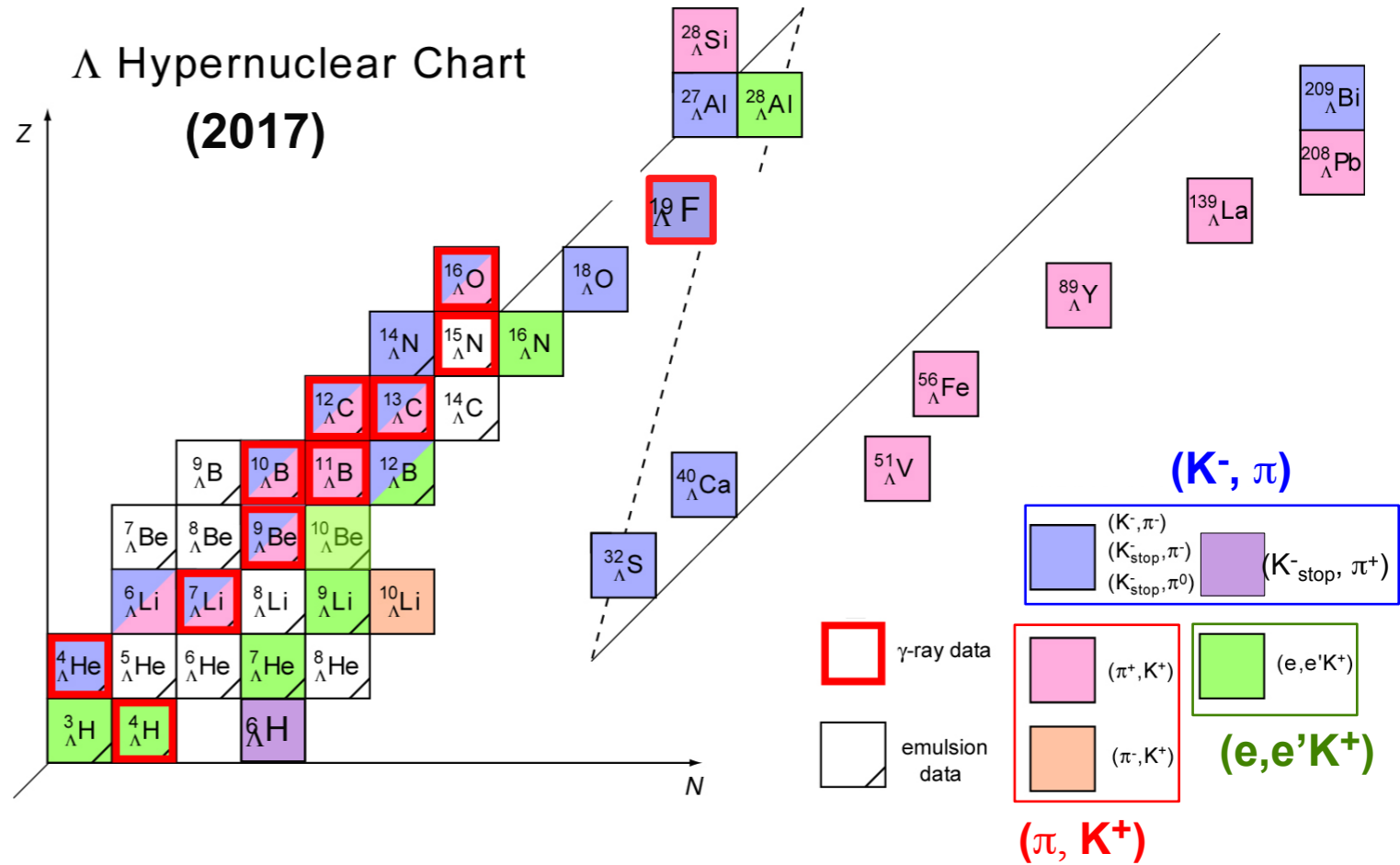


Production Threshold for  $\Lambda'$ s :  $p \geq 100 \text{ MeV}$

<http://eaae-astronomy.org/blog/?cat=254>



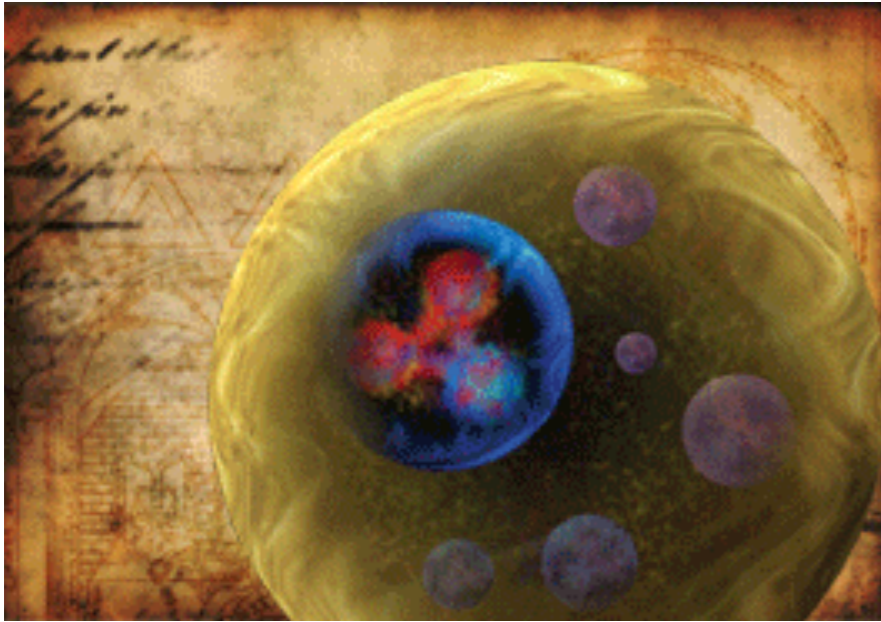
Hypernuclei can be produced  
 Binding Energy of  $\Lambda$  to nucleus = 30 MeV



O. Hashimoto and H. Tamura, Prog. Part. Nucl. Phys. 57 (2006) 564.

Wirth and Roth Phys.Rev.Lett. 117 (2016) 182501

<http://eaae-astronomy.org/blog/?cat=254>

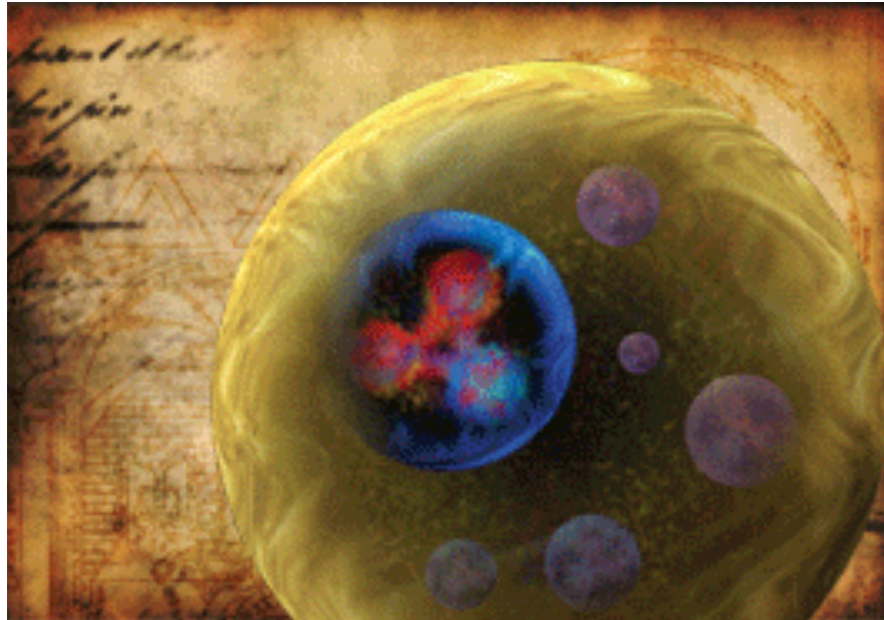


Hypernuclei can be produced  
 Binding Energy of  $\Lambda$  to nucleus = 30 MeV

Nothing is known about  $\Sigma$  - hypernuclei



<http://eaae-astronomy.org/blog/?cat=254>

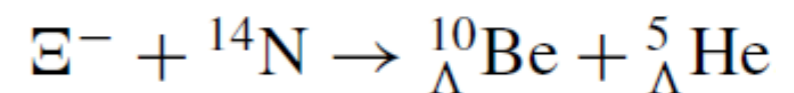
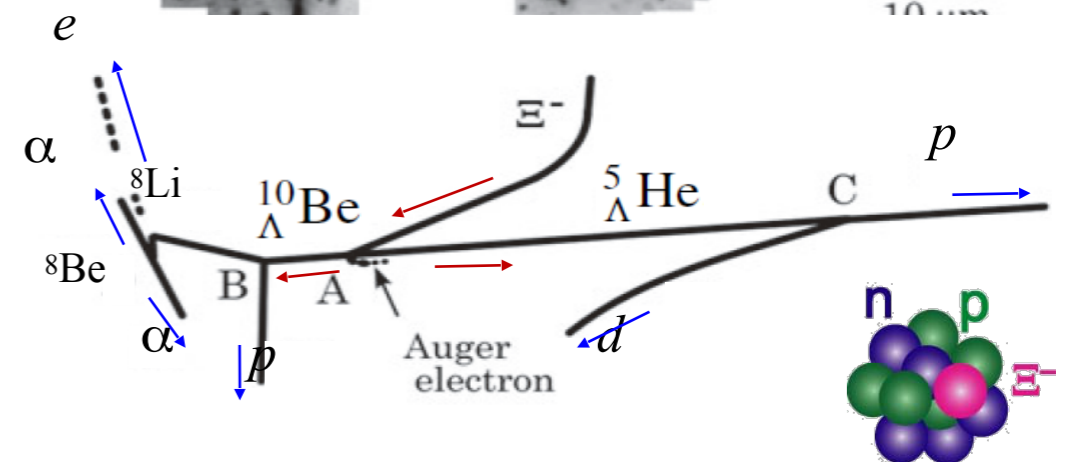
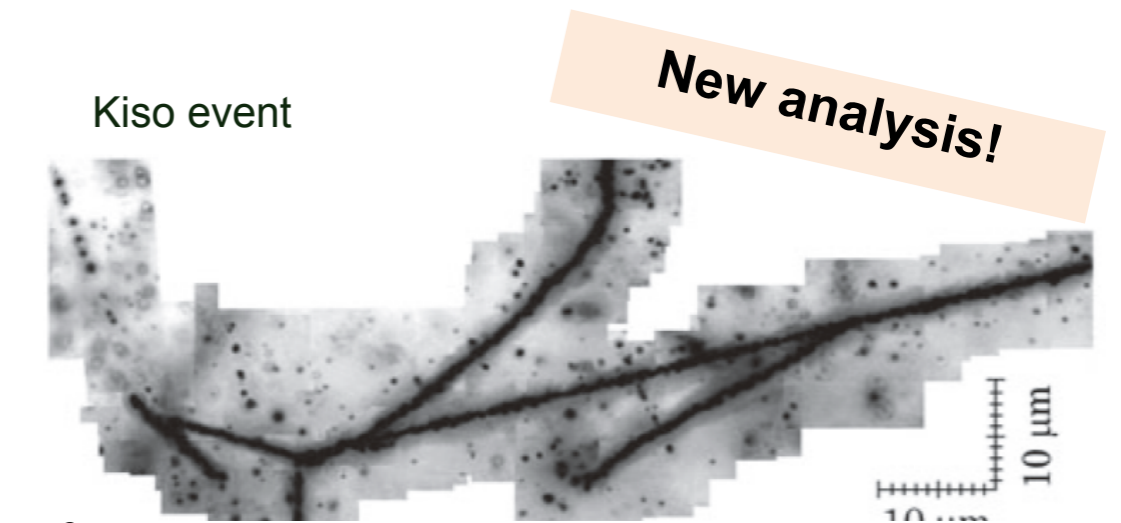


Hypernuclei can be produced  
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Nothing is known about  $\Sigma$  - hypernuclei

$\Xi$  - Hypernucleus shows a shallow attractive interaction

Courtesy H. Tamura, Bormio Winter Meeting 2018



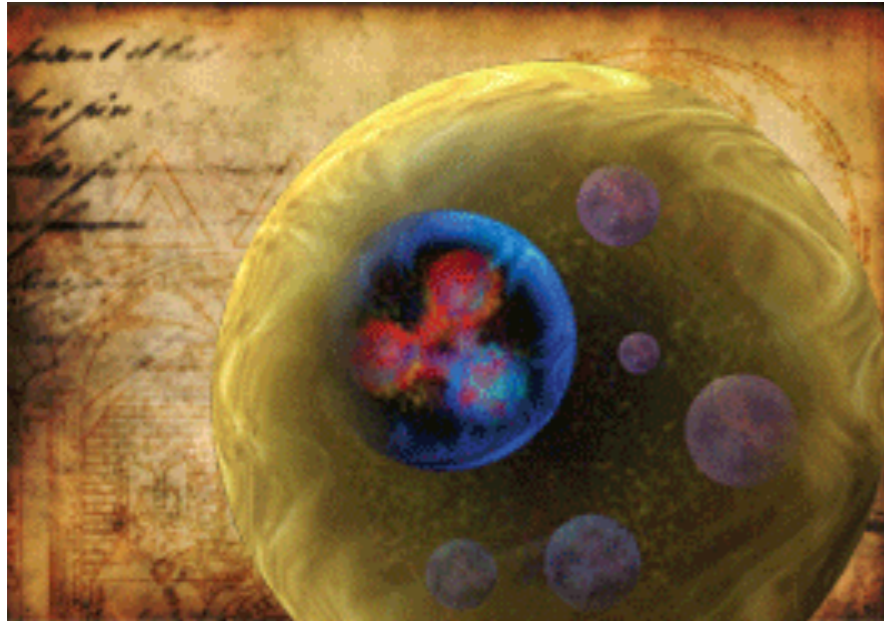
The first clear  $\Xi$  hypernucleus

$$B_{\Xi^-} = 4.38 \pm 0.25 \text{ MeV},$$

$$B_{\Lambda} = 1.11 \pm 0.25 \text{ MeV}$$

*K. Nakazawa et al. PTEP 2015, 033D02*

<http://eaae-astronomy.org/blog/?cat=254>



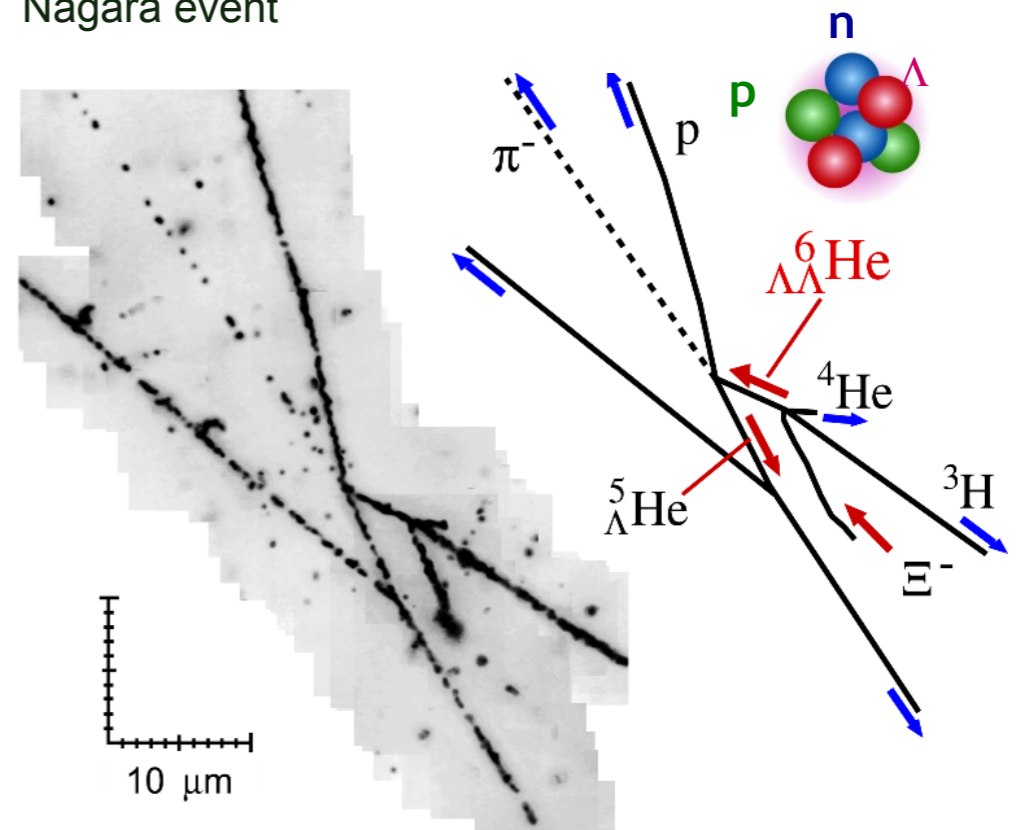
Hypernuclei can be produced  
 Binding Energy of  $\Lambda$  to nucleus = 30 MeV

Nothing is known about  $\Sigma$  - hypernuclei

$\Xi$  - Hypernucleus shows a shallow attractive interaction

Even  $\Lambda\Lambda$ -hypernuclei exist

Nagara event



$$\Delta B_{\Lambda\Lambda} = 0.67 \pm 0.17 \text{ MeV}$$

*H. Takahashi et al., PRL 87 (2001) 212502*

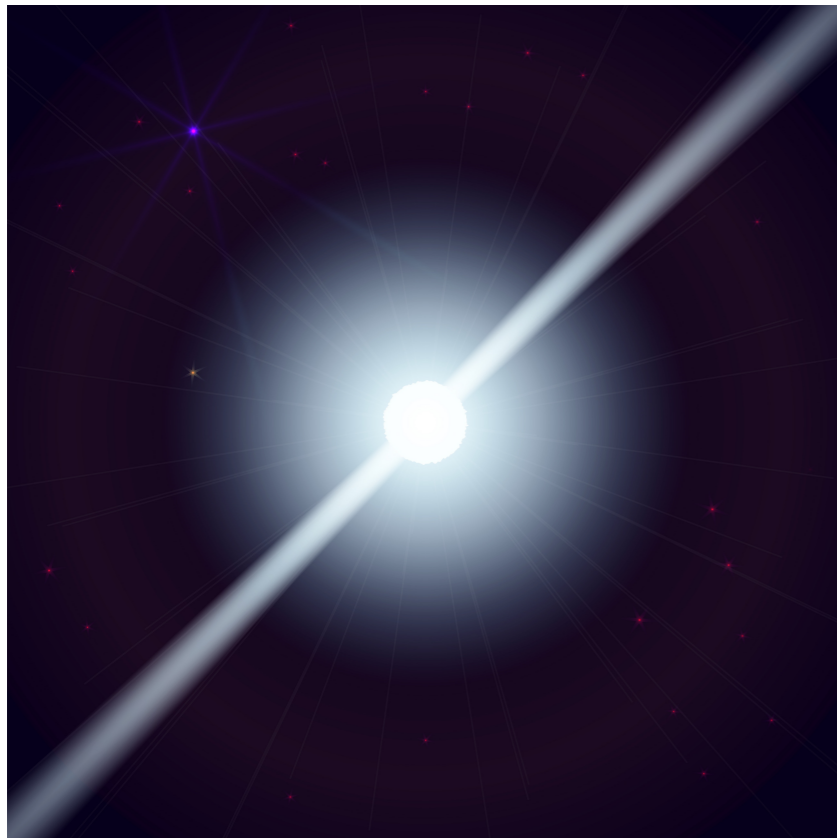
**$\Lambda$ - $\Lambda$  is weakly attractive**

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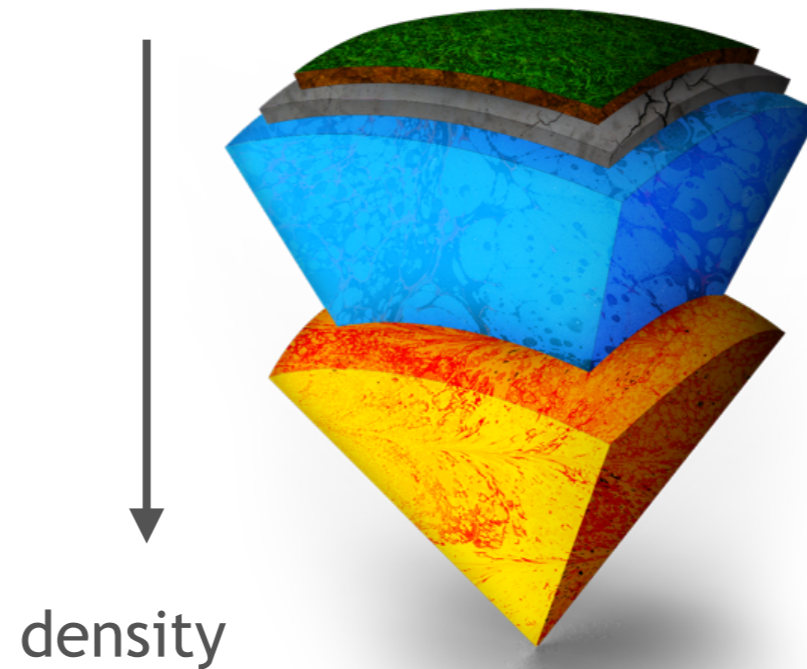
# Hyperon Puzzle in Neutron Stars

$$R \approx 10 - 15 \text{ Km}$$

$$M \approx 1.5 - 2 M_{\odot}$$



Courtesy of Shutterstock



Outer Crust:  
Ions, electrons Gas,  
Neutrons

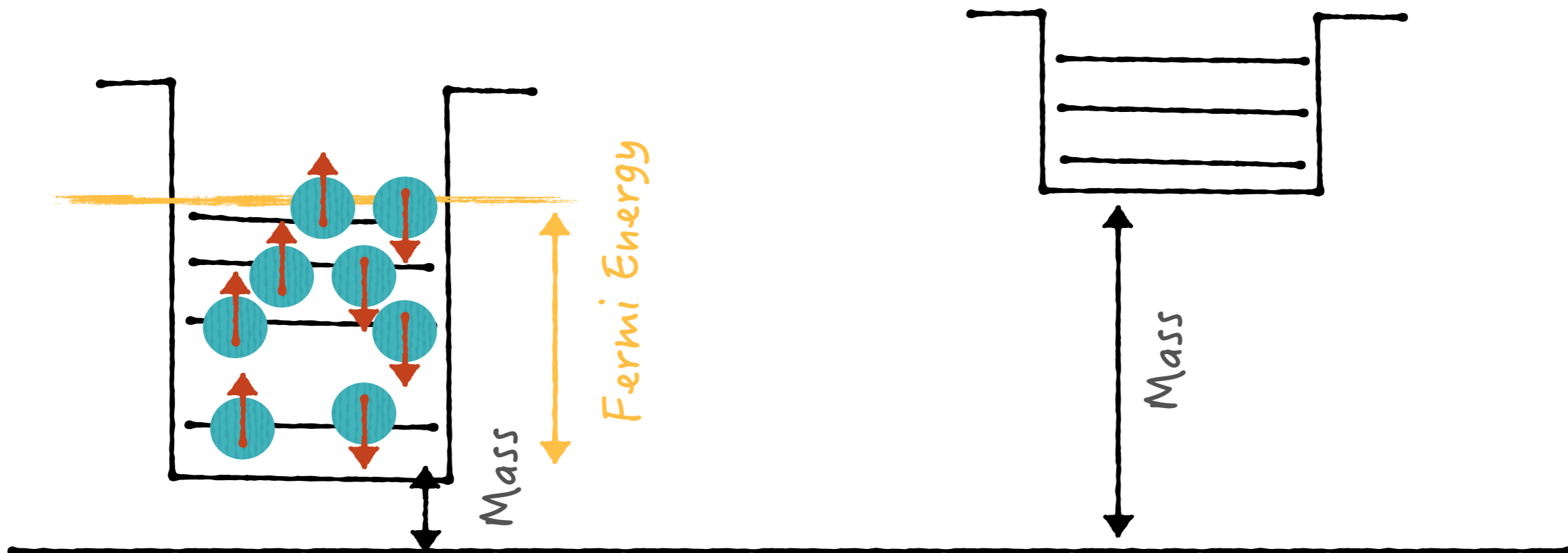
Inner Core: ??  
Neutrons ?? Protons ??  
Hyperons ??  
Quark Matter ??

- Very high density in the interior
- Rotating object emitting Synchrotron radiation in Radio-Frequency (Pulsar character)
- Mass measured in binary systems with White Dwarfs (Shapiro Delay, WD Spectroscopy)
- Radius Measurement very difficult

## What is inside Neutron Stars??

Neutrons (uud,  $m = 938 \text{ MeV}$ )

$\Lambda$  Hyperons (uds,  $m = 1115 \text{ MeV}$ )

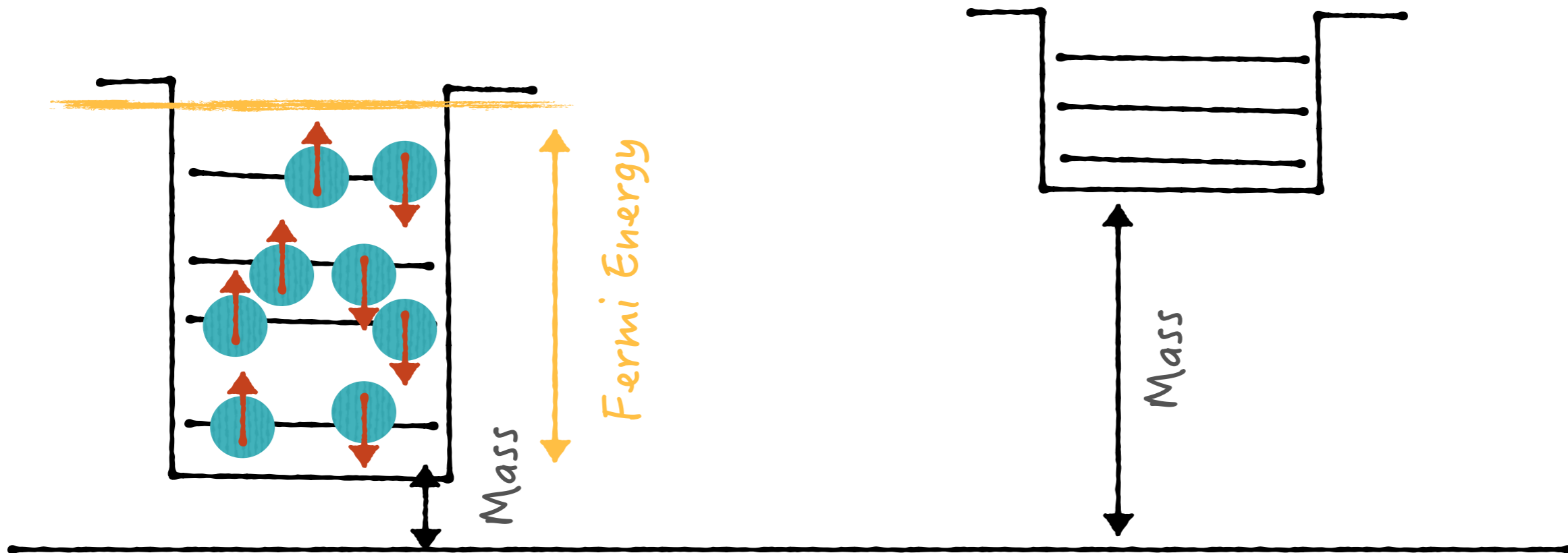


Chemical Potential  $\mu = E_F + mass$

If the density increases also the Fermi Energy increases and hence the chemical potential

Neutrons (uud,  $m = 938 \text{ MeV}$ )

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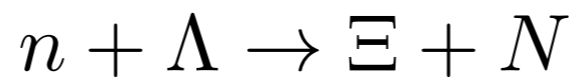
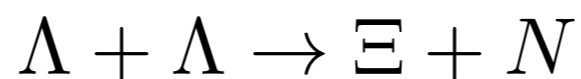
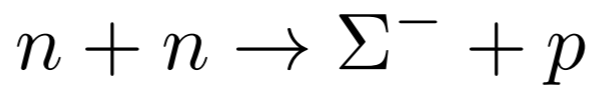
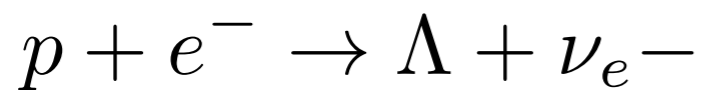
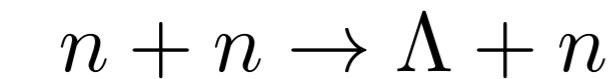
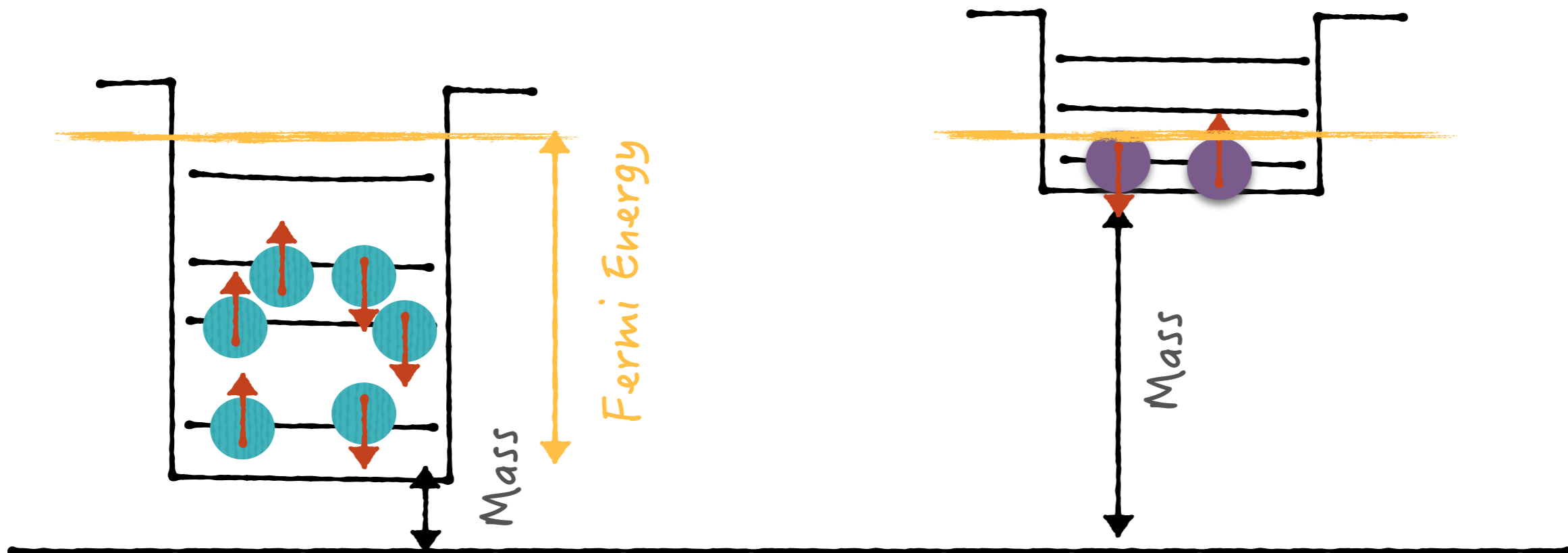


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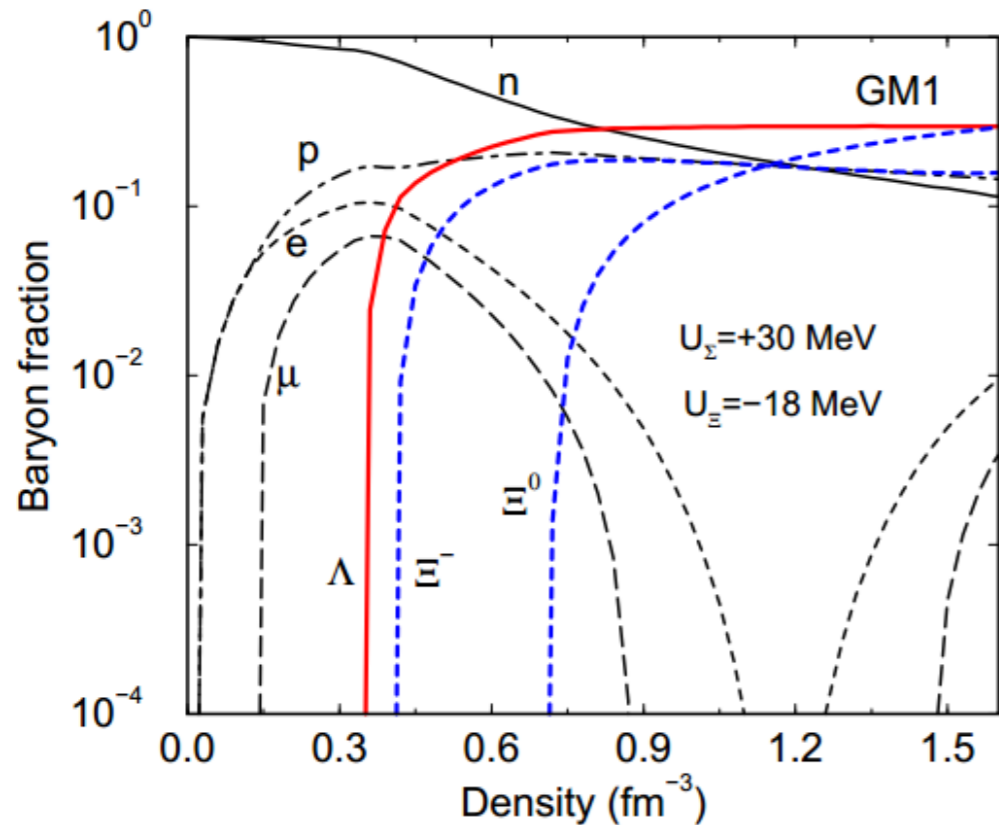
In order to have chemical equilibrium  $\mu_{neutron} = \mu_{\Lambda}$

If the  $\Lambda$ -nucleon interaction is attractive the processes is even more likely

Hyperons should appear in dense neutron-rich matter starting from moderate large densities

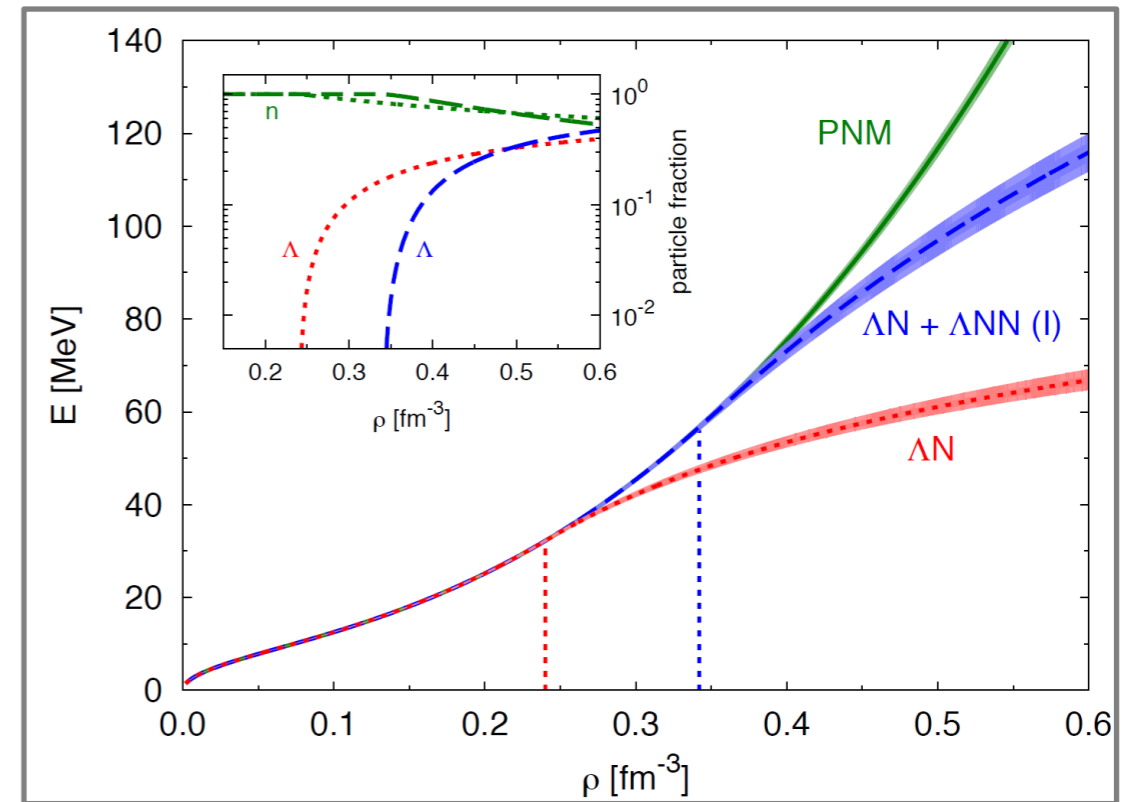
Threshold depends on the Y-N interaction

The appearance of Hyperons softens the EoS



$\rho_0$   $2\rho_0$   $3\rho_0$   $4\rho_0$  .....

J. Schaffner-Bielich, NPA 804 (2008)



D. Lonardoni, A. Lovato, S. Gandolfi, F. Pederiva Phys. Rev. Lett. 114, 092301 (2015)

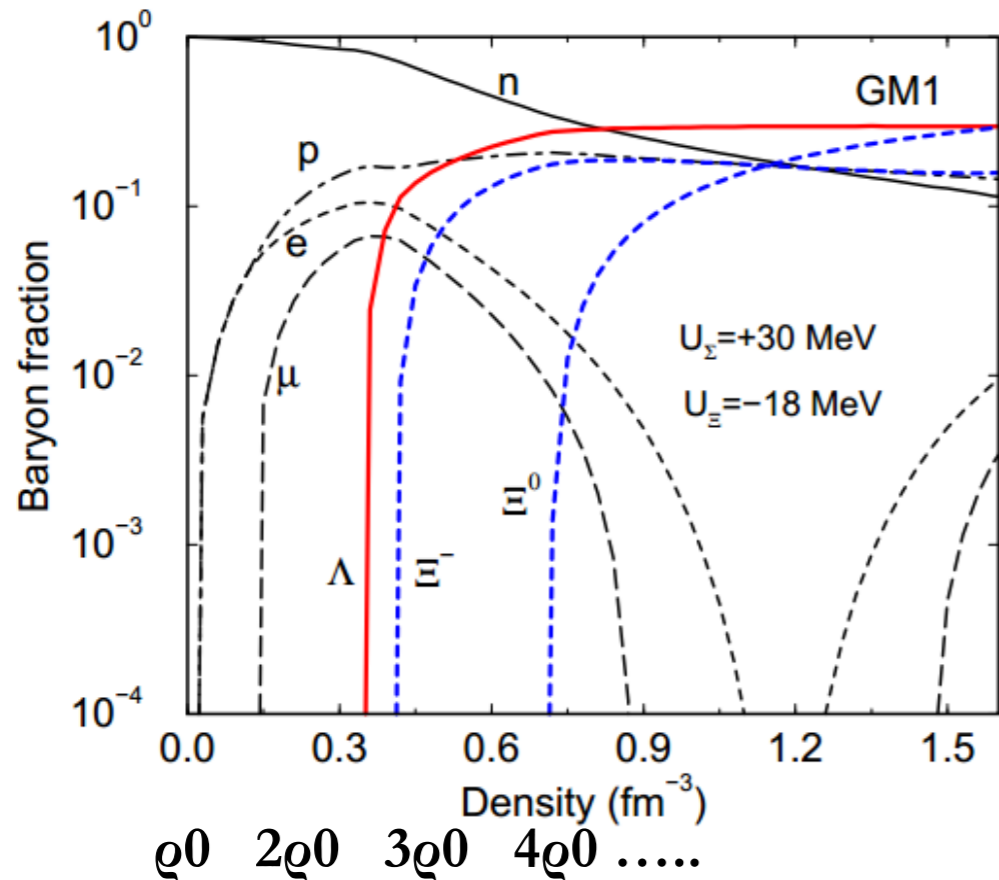


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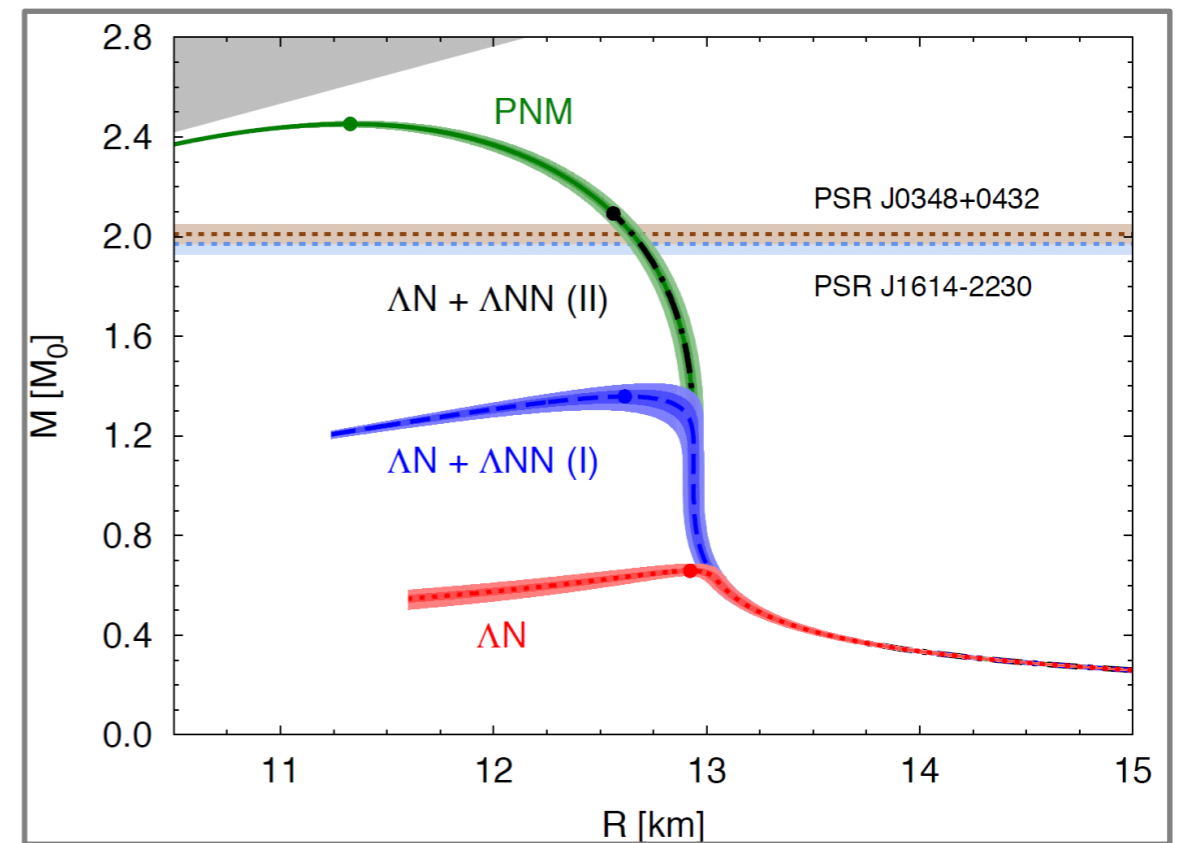
Threshold depends on the Y-N interaction

The appearance of Hyperons softens the EoS

➡ Maximum NS masses get smaller



J. Schaffner-Bielich, NPA 804 (2008)



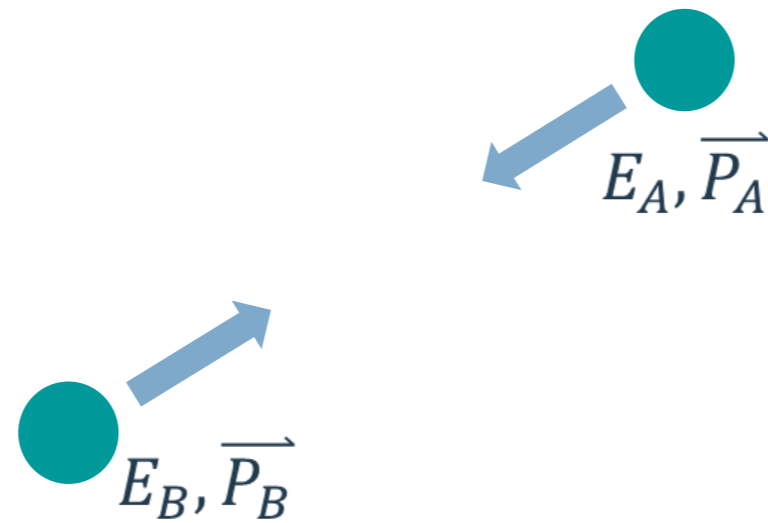
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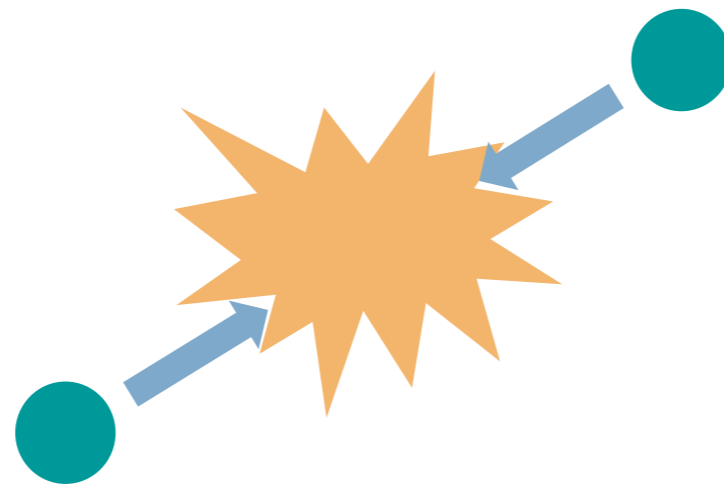
These predictions are only qualitative to this end

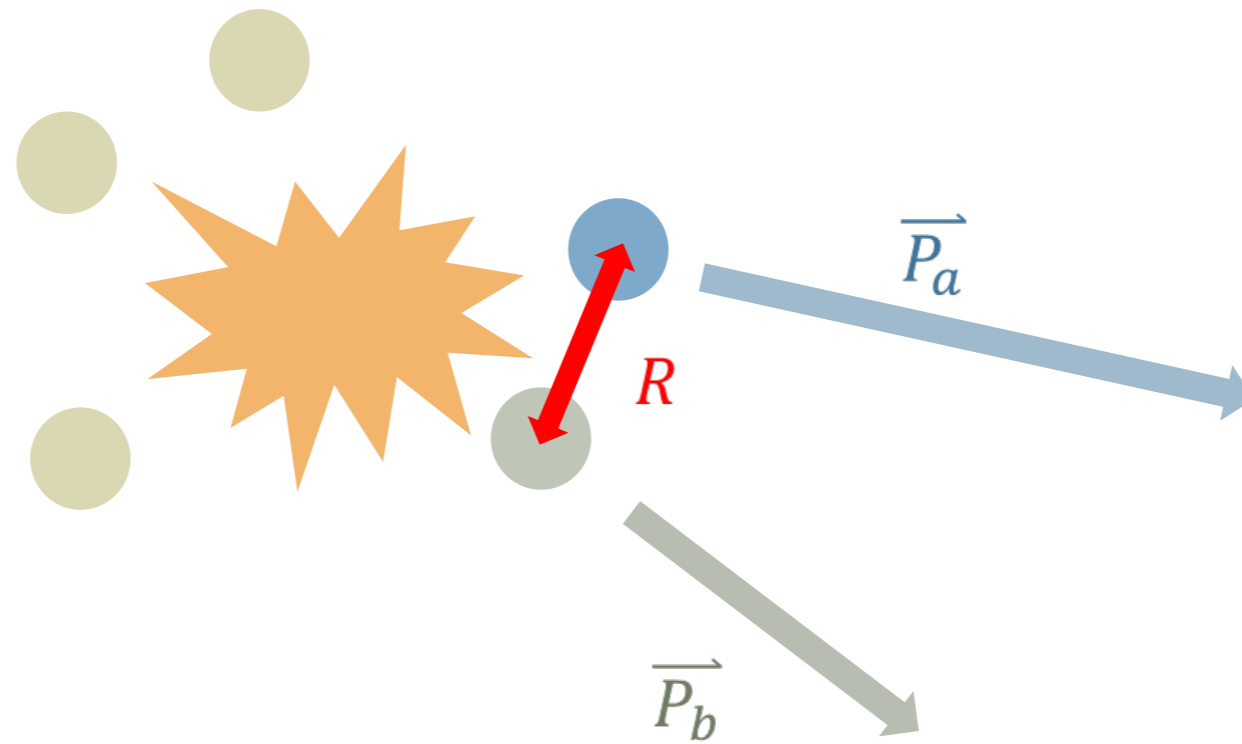
- \* Measure 2 body and 3 bodies interactions that allow to constrain a theory as  $\chi EFT$
- \* Extend this search to all hyperons
- \* Compute more reliable single particle potentials
- \* Compute more solid EoS for dense nuclear matter with hyperons and kaons content

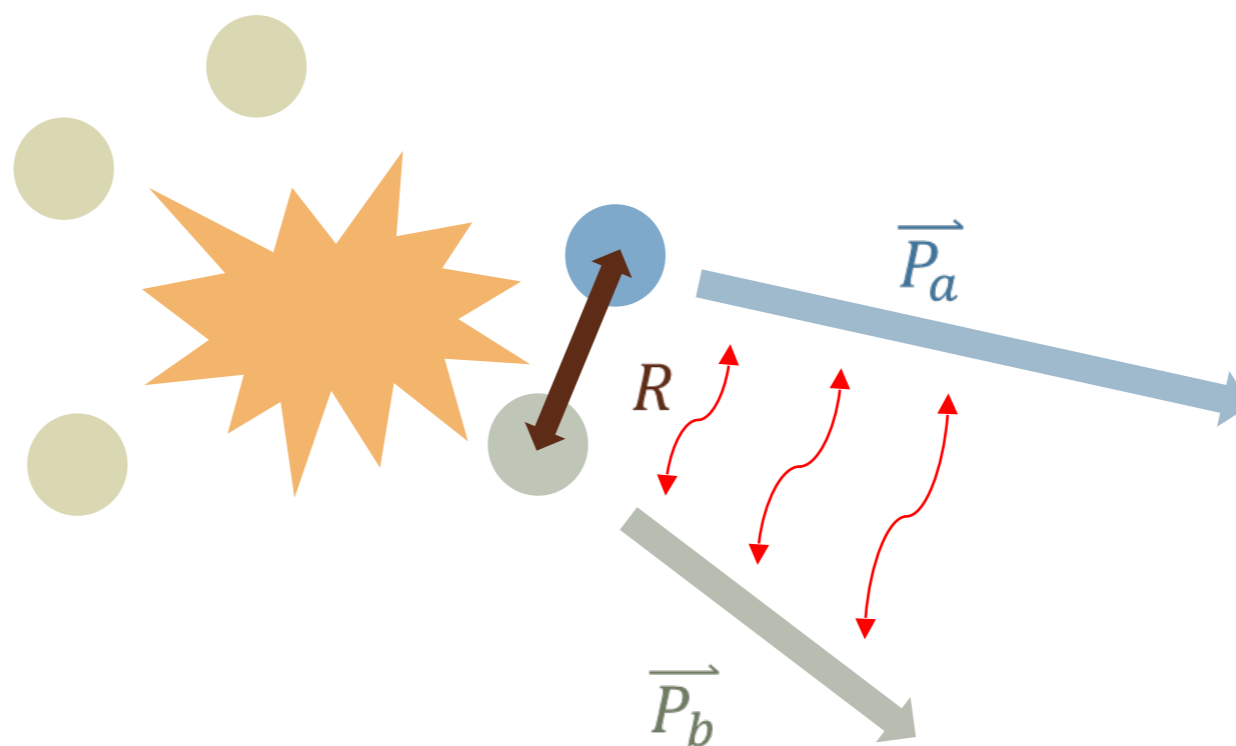
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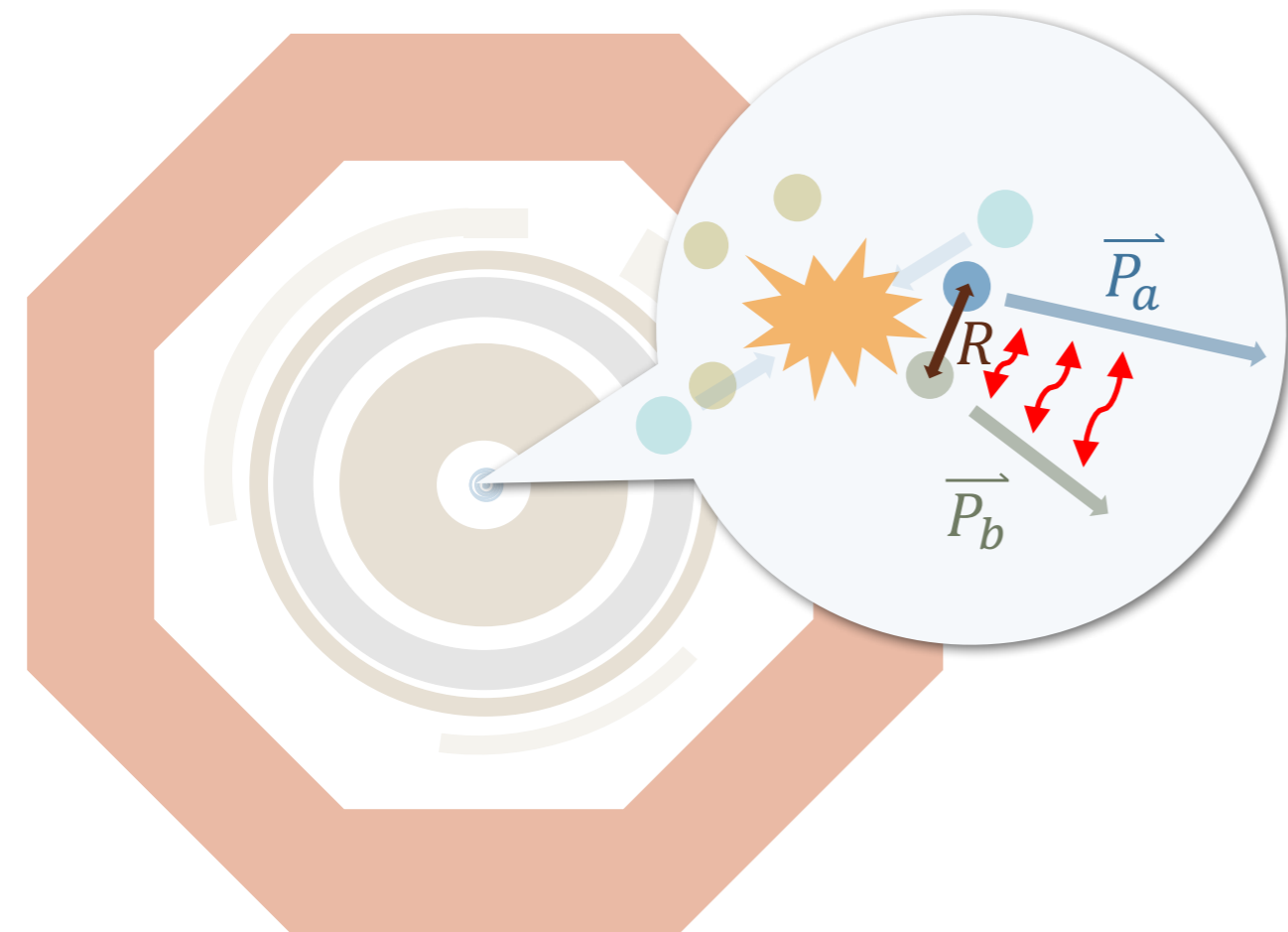
# The measurement of Hadron Hadron Correlations







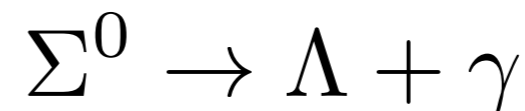
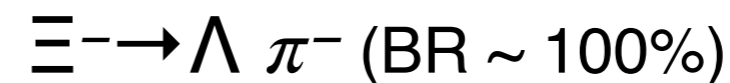
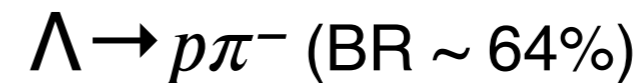




We measure **pp, p $\Lambda$ ,  $\Lambda\Lambda$ , p $\Xi$ , pK, p $\Sigma^0$ , p $\Omega$ -**

Proton and Pion identification with TPC and TOF

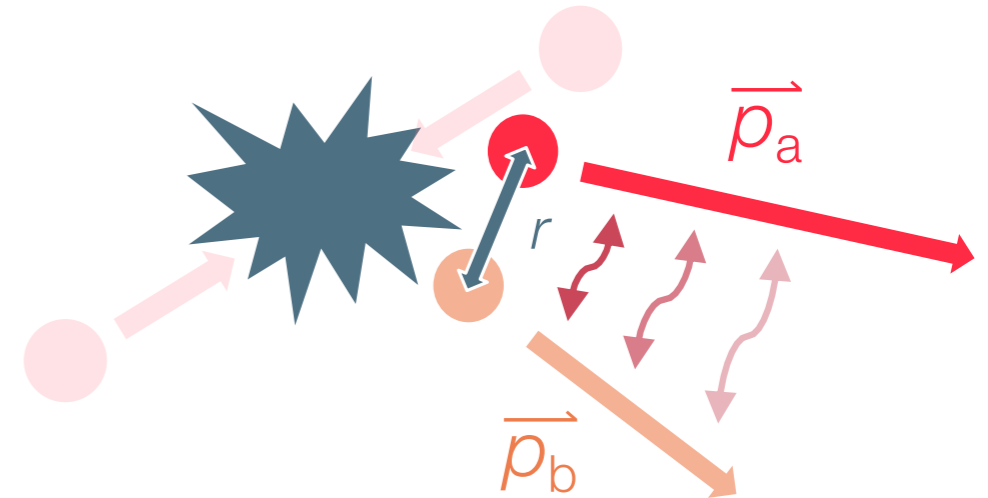
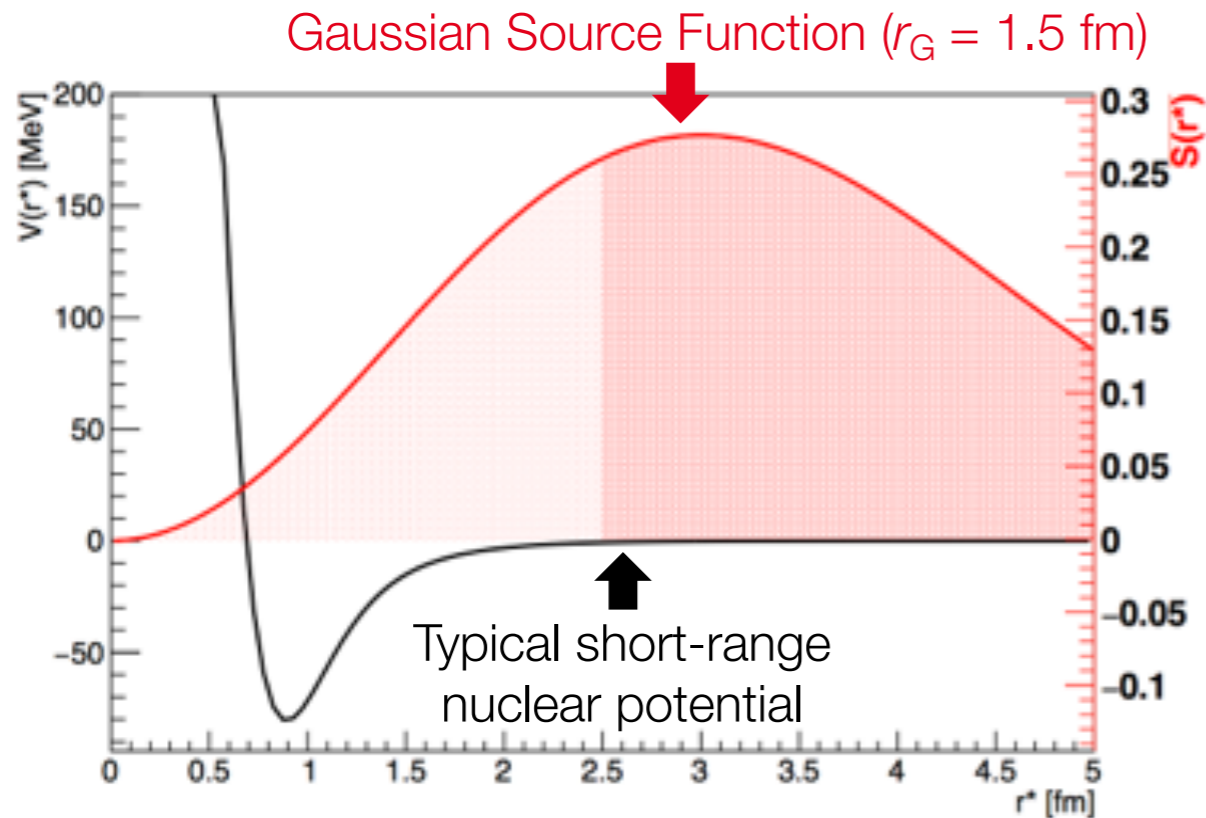
Reconstruction of hyperons



Datasets:

- pp 7 TeV:  $3.4 \cdot 10^8$  MB Events
- pp 5 TeV:  $10 \cdot 10^8$  MB Events
- pp 13 TeV:  $15 \cdot 10^8$  MB Events
- pp 13 TeV:  $1 \cdot 10^8$  HM Events
- p-Pb 5.02 TeV:  $6.0 \cdot 10^8$  MB Events





Small particle-emitting source created in pp and p–Pb collisions at the LHC

- Essential ingredient for detailed studies of the strong interaction
- Assuming the same particle source for all pairs

» **p–p correlation** is used to constrain the source, since Coulomb and Strong interactions are well known

The correlation function:

$$C(k^*) = \frac{P(\mathbf{p}_a, \mathbf{p}_b)}{P(\mathbf{p}_a)P(\mathbf{p}_b)},$$

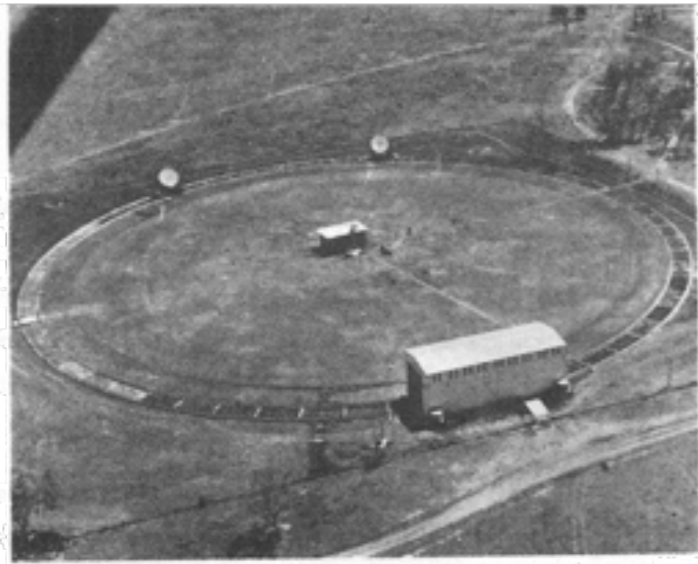
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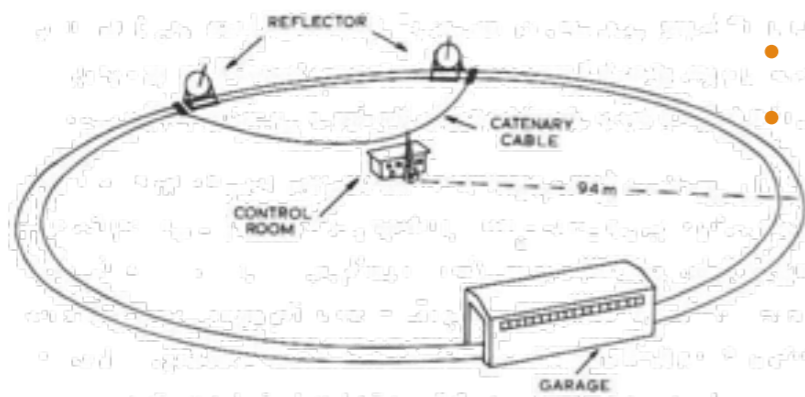
Experimentally obtained as:

$$C(k^*) = \mathcal{N} \frac{N_{Same}(k^*)}{N_{Mixed}(k^*)}$$

- **1950 Hanbury Brown-Twiss (HBT) interferometry** (Nature volume 178 (1956)1046–1048)



- **Two mirror connected to PMT**
- **“correlator”** ⇒ “.. ... collected light as rain in a bucket ..” (more or less what we do when collecting pairs of particles)
- **REMARK:** it was already known that such intensity interferometry was working for radio-wave (classical object). The application to photons has indeed been received initially with skepticism ⇒ wave-particle duality and B/F Quantum Statistics
- **HB and T applied the technique to measure the angular size of the star Sirius (8.6 ly) by studying optical intensity correlations between two telescopes**



Data taken for 18 hours over 5 month period  
 Data yielded an angular diameter of (0.005936")

$$0.0068'' \pm 0.0005'' = 3.1 \times 10^{-8} \text{ radians}$$

G.Baym Acta Phys.Polon. B29 (1998) 1839-1884

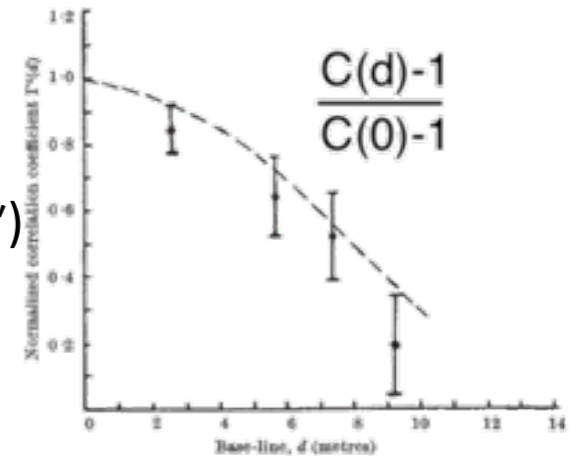
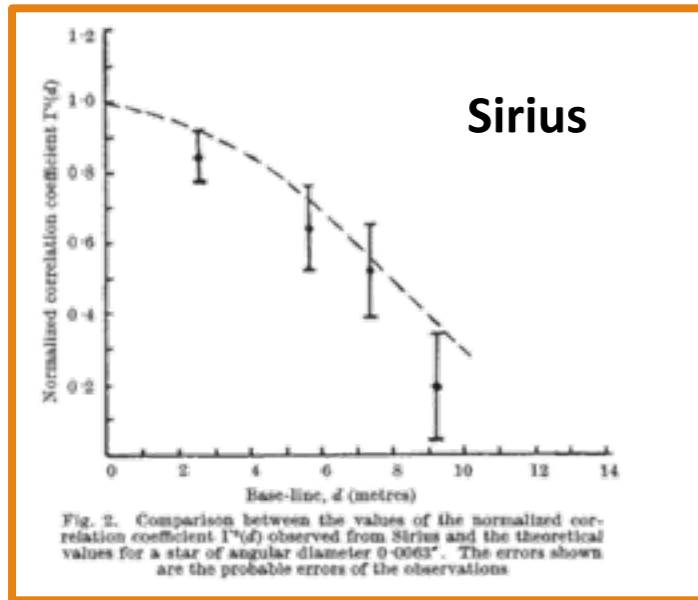
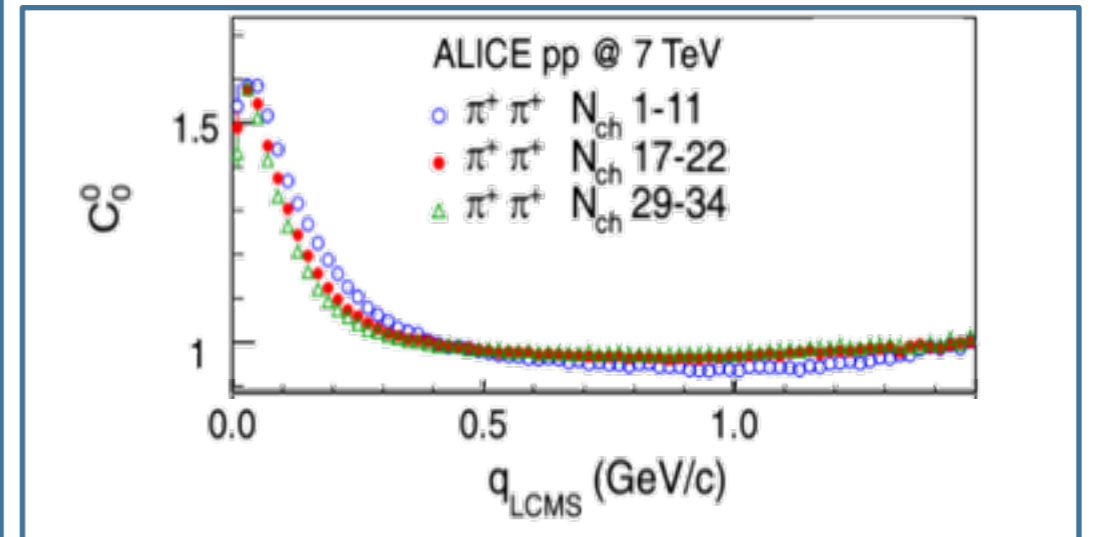
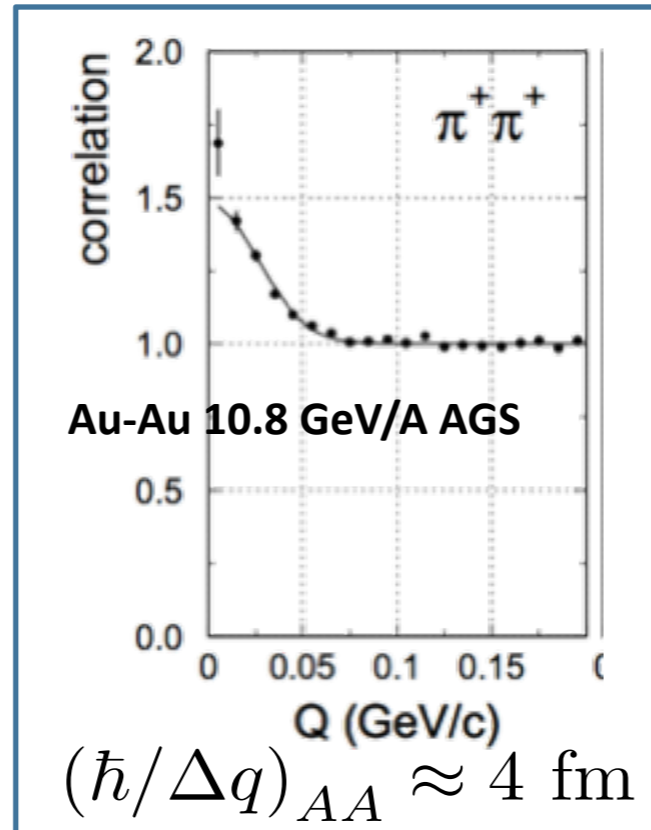


Fig. 2. Comparison between the values of the normalized correlation coefficient  $I^2(d)$  observed from Sirius and the theoretical values for a star of angular diameter 0.0063". The errors shown are the probable errors of the observations



1950 Hanbury Brown-Twiss (HBT)  
Sirius



**p-p 7 TeV ALICE**  
(ALICE Coll. PRD84, 112004  
(2011))  
 $(\hbar/\Delta q)_{pp} \approx 1 - 1.5 \text{ fm}$

- 1959 Goldhaber, Goldhaber, Lee and Pais, 1977 Koonin
- From end of 80s  $\Rightarrow$  used to study the **space-time geometry evolution in heavy-ion collisions and elementary collisions**

The correlation function:

$$C(k^*) = \frac{P(\mathbf{p}_a, \mathbf{p}_b)}{P(\mathbf{p}_a)P(\mathbf{p}_b)},$$

Experimentally obtained as:

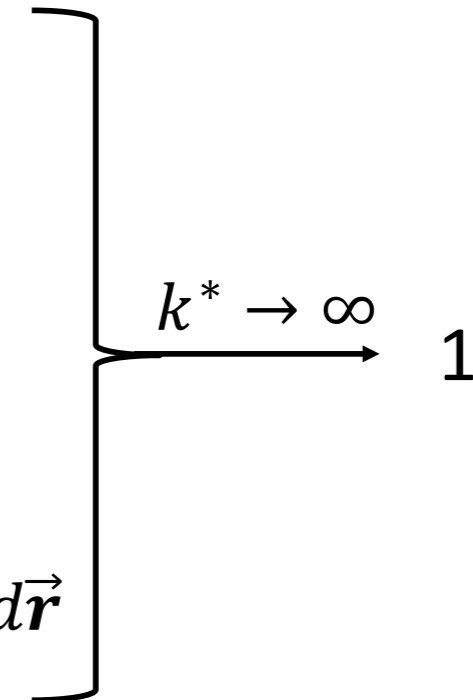
$$C(k^*) = \mathcal{N} \frac{N_{Same}(k^*)}{N_{Mixed}(k^*)}$$

Given by:

$$C(k^*) = \int S(\mathbf{r}, k^*) |\psi(\mathbf{r}, k^*)|^2 d\vec{r}$$



$$k^* = \frac{|\mathbf{p}_a^* - \mathbf{p}_b^*|}{2} \text{ and } \mathbf{p}_a^* + \mathbf{p}_b^* = 0$$



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Assumption of a **common source** with **Gaussian shape** for the **pp, pΛ, pΞ, ΛΛ** and **pK** Correlation Function

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Source

Relative Wave Function

$$k^* = \frac{|\mathbf{p}_a^* - \mathbf{p}_b^*|}{2} \text{ and } \mathbf{p}_a^* + \mathbf{p}_b^* = 0$$

$k^* \rightarrow \infty$  1

Assumption of a **common source** with **Gaussian shape** for the **pp**, **pΛ**, **pΞ**, **ΛΛ** and **pK** Correlation Function

**Strong constraint**



The correlation function:

$$C(k^*) = \frac{P(\mathbf{p}_a, \mathbf{p}_b)}{P(\mathbf{p}_a)P(\mathbf{p}_b)},$$

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Assumption of a **common source** with **Gaussian shape** for the **pp, pΛ, pΞ, ΛΛ and pK** Correlation Function

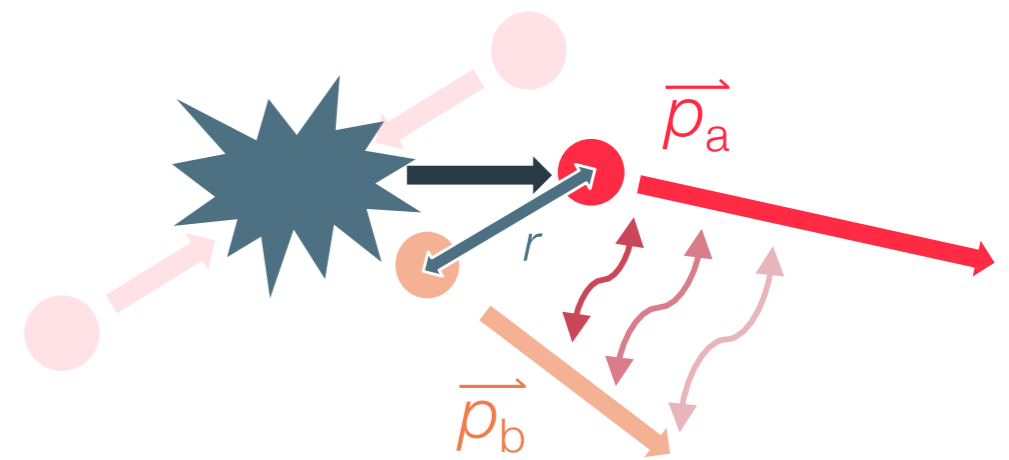
**Strong constraint**

correlations functions allow to study the interactions

## (An)isotropic flow



## + Strongly decaying resonances



- (An)isotropic pressure gradients affect the emission
  - Initial geometric anisotropies introduce a transverse modulation
  - Expanding source with *common velocity field*
- Affects particles depending on their mass

- Resonances with  $c\tau \sim r_0 \sim 1 \text{ fm}$  ( $\Delta, N^*$ ) introduce an *exponential tail* to the source
- Different for each particle species!

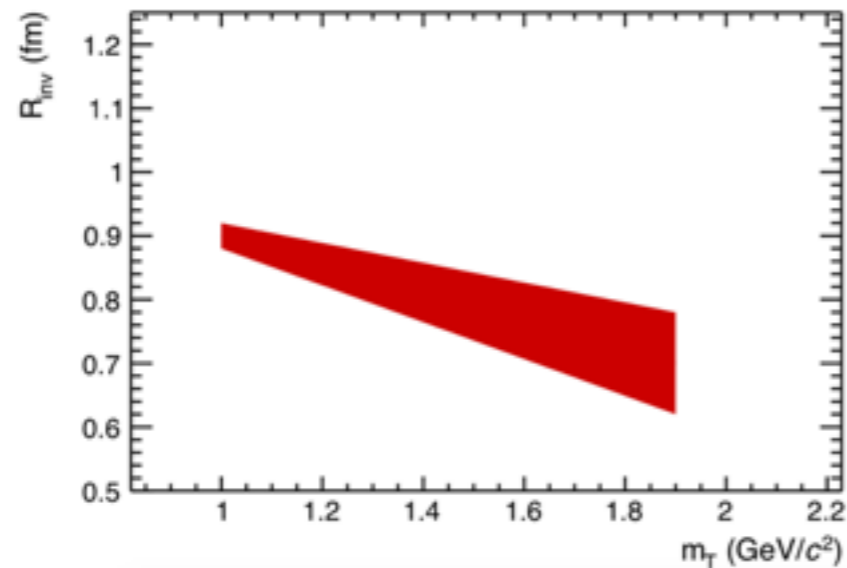
(An)isotropic flow

+ Strongly decaying resonances

Gaussian core

⊗

Exponential tail



| Particle | Primordial fraction | Resonances                 |                        |
|----------|---------------------|----------------------------|------------------------|
|          |                     | $1 < c\tau < 2 \text{ fm}$ | $c\tau > 2 \text{ fm}$ |
| Proton   | 33 %                | 56 %                       | 2 %                    |
| Lambda   | 35 %                | 8 %                        | 58 %                   |

- Yield of resonances determined from Canonical Statistical Hadronization Model
- Priv. Comm. with Prof. F. Becattini  
J.Phys. G38 (2011) 025002.

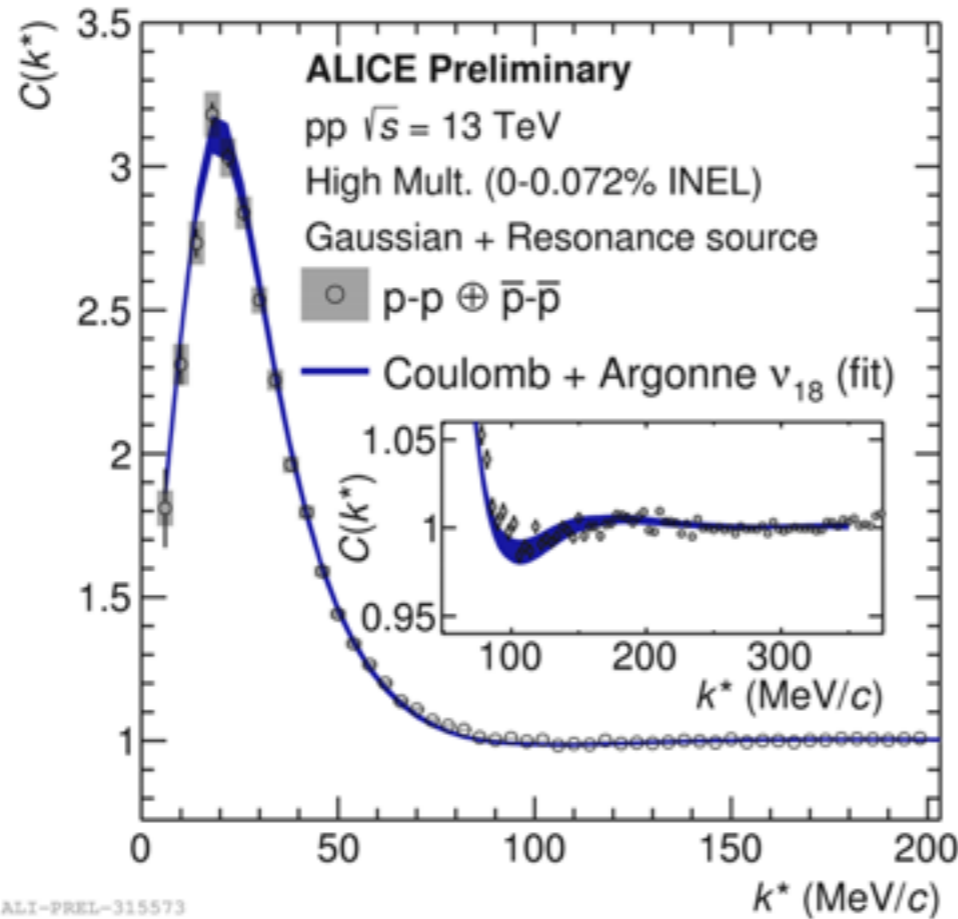
(An)isotropic flow

Strongly decaying resonances

Gaussian core



Exponential tail



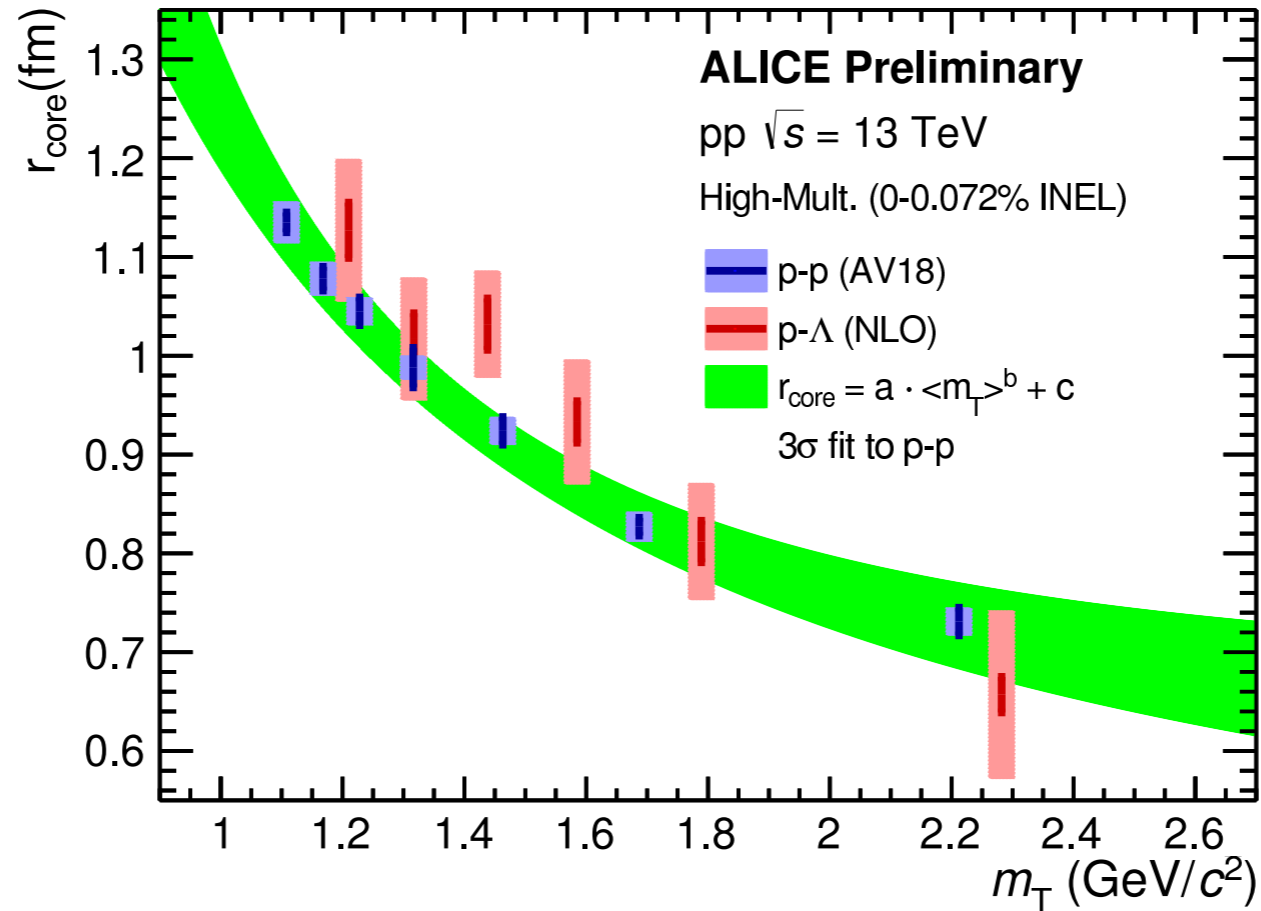
$$r_{\text{Core}} = 0.995 \pm 0.005 \begin{matrix} +0.024 \\ -0.022 \end{matrix} \text{ fm}$$

$$r_{\text{Eff}} = 1.249 \pm 0.008 \begin{matrix} +0.024 \\ -0.021 \end{matrix} \text{ fm}$$

$$C(k) = \int dr^3 \phi_{rel}^2(r, k) \exp\left(-\frac{r^2}{4R_G^2}\right)$$

Coulomb + AV18 Potential  
 + Modelled Source

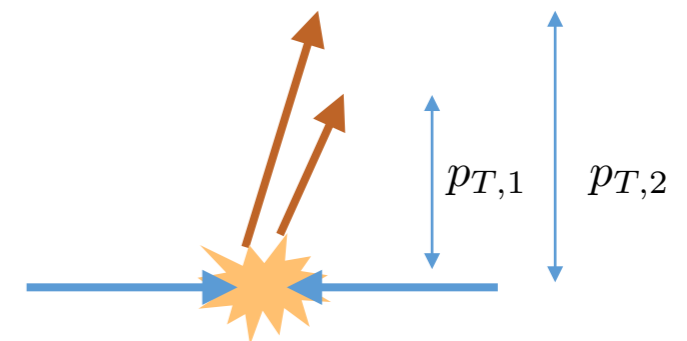
arXiv:2004.08018



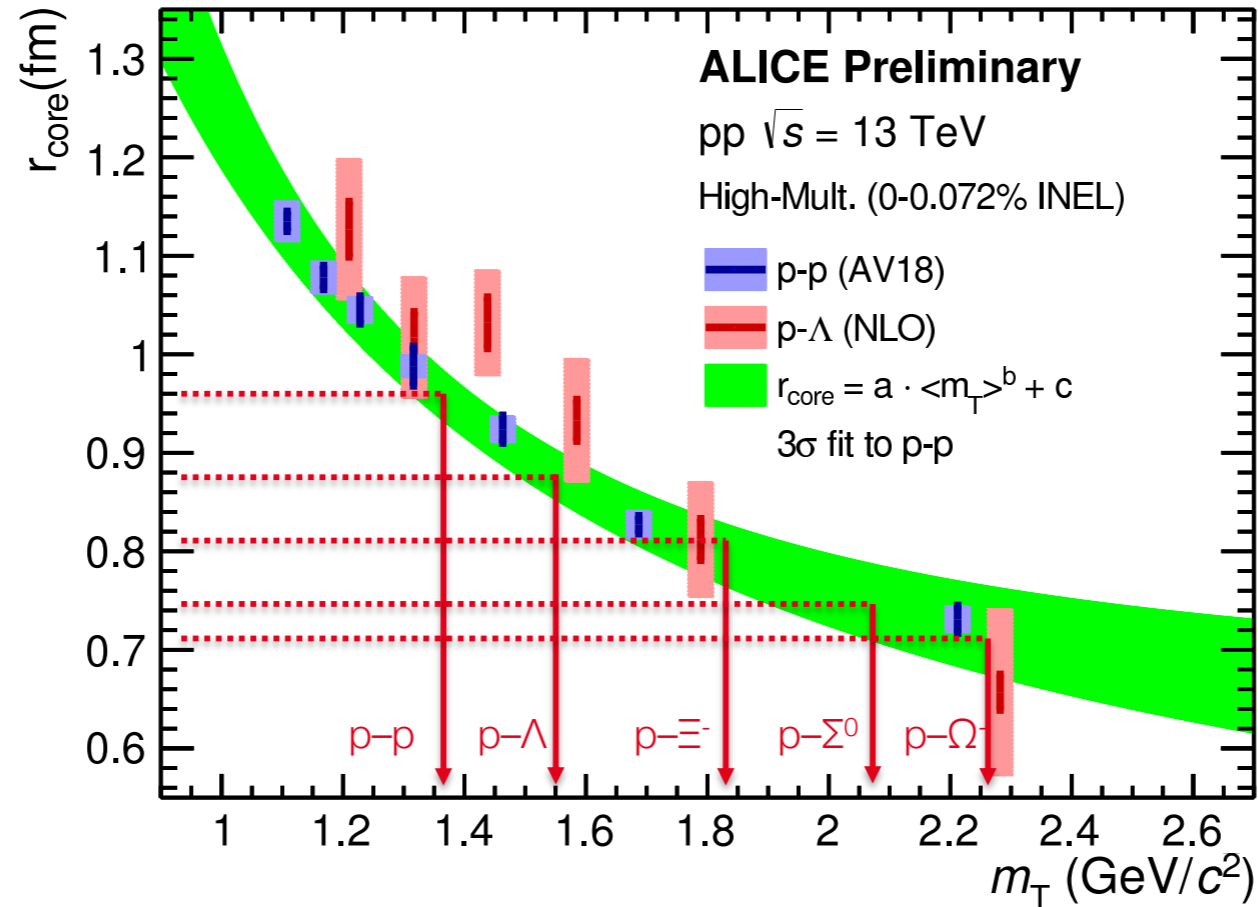
ALI-PREL-315640

- Core radius for p-p and p- $\Lambda$  in good agreement

$$m_T = \sqrt{m_{\text{pair}}^2 + k_T^2} \quad k_T = |\vec{p}_{T,1} + \vec{p}_{T,2}|/2$$



arXiv:2004.08018



| Pair          | $r_{\text{Core}}$ (fm) | $r_{\text{Eff}}$ (fm) |
|---------------|------------------------|-----------------------|
| p-p           | 1.00                   | 1.25                  |
| p- $\Lambda$  | 0.88                   | 1.30                  |
| p- $\Sigma^0$ | 0.75                   | 1.14                  |
| p- $\Xi^-$    | 0.80                   | 0.92                  |
| p- $\Omega^-$ | 0.73                   | 0.85                  |

ALI-PREL-315640

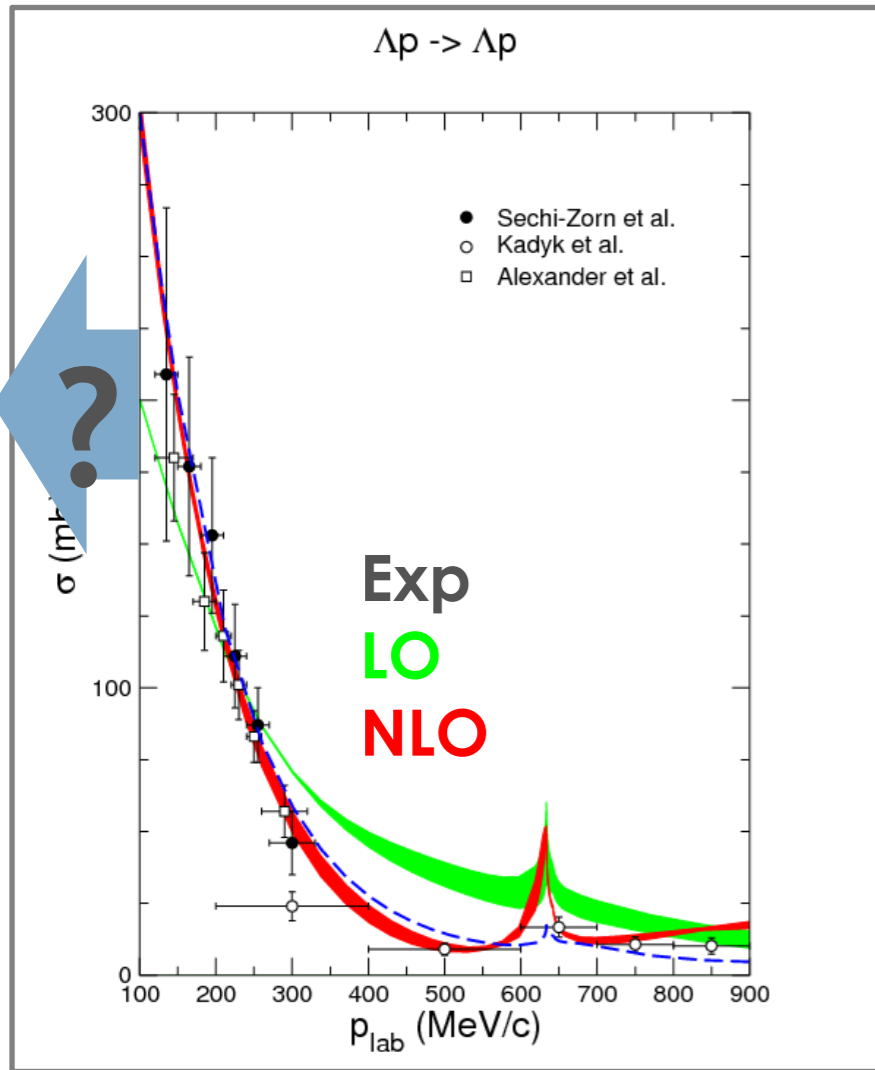
- Fix the value of  $r_{\text{Core}}$  of each particle species based on their  $\langle m_T \rangle$ 
  - Add specific resonance contribution to obtain the corresponding pair source

# Femtoscscopy with $\Lambda$ and $\Sigma$ baryons

Precision and pioneering feasibility studies in the  $|S| = 1$  sector

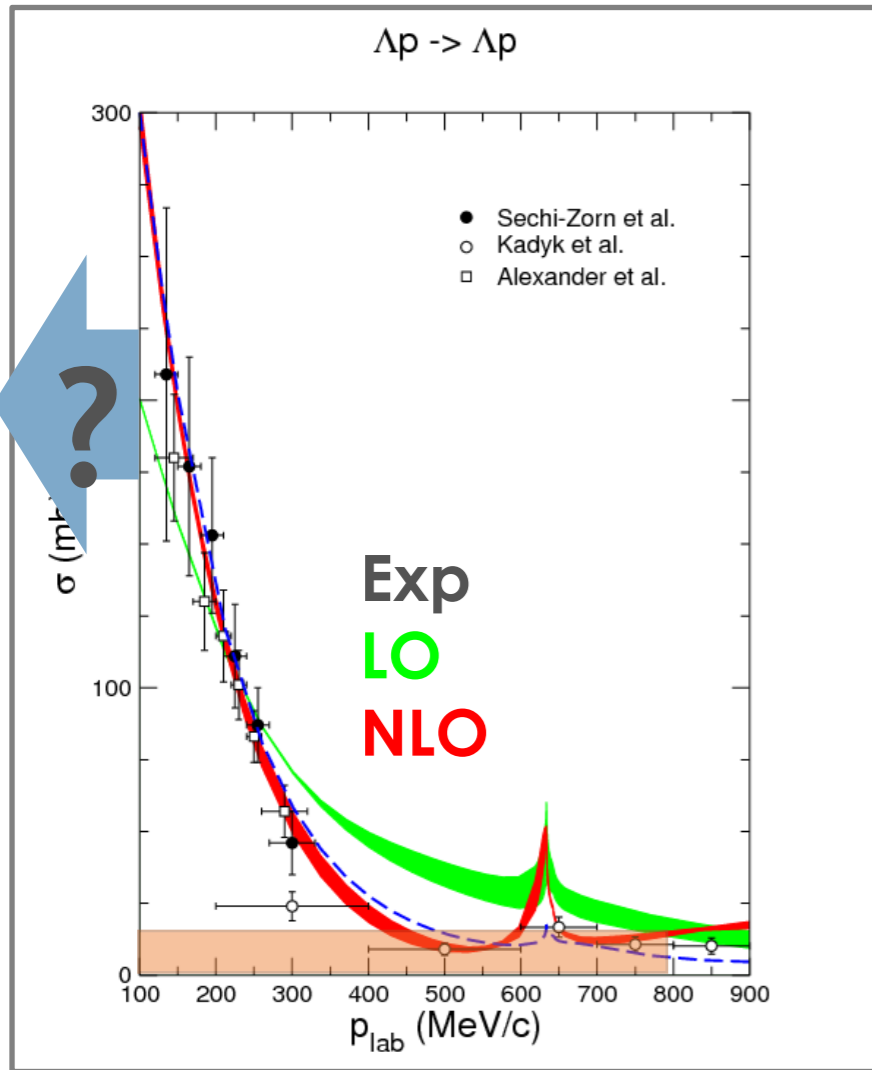
ALICE Collaboration, *Phys.Lett.B* 805 (2020) 135419

pp  $\sqrt{s} = 13$  TeV (high mult.)



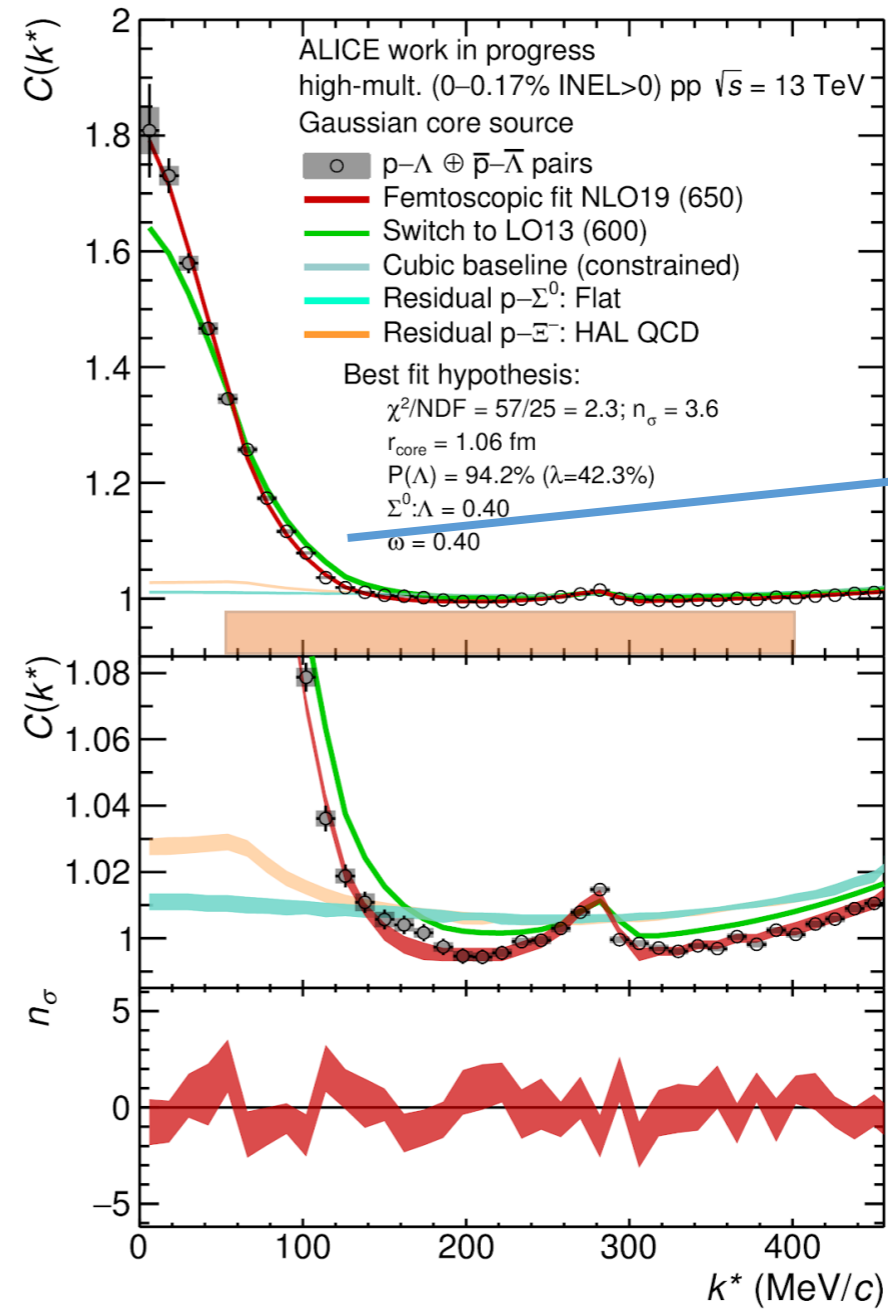
LO: H. Polinder, J.H., U. Meißner, NPA 779 (2006) 244  
 NLO: J.Haidenbauer., N.Kaiser, et al., NPA 915 (2013) 24





LO: H. Polinder, J.H., U. Meißner, NPA 779 (2006) 244  
NLO: J.Haidenbauer., N.Kaiser, et al., NPA 915 (2013) 24

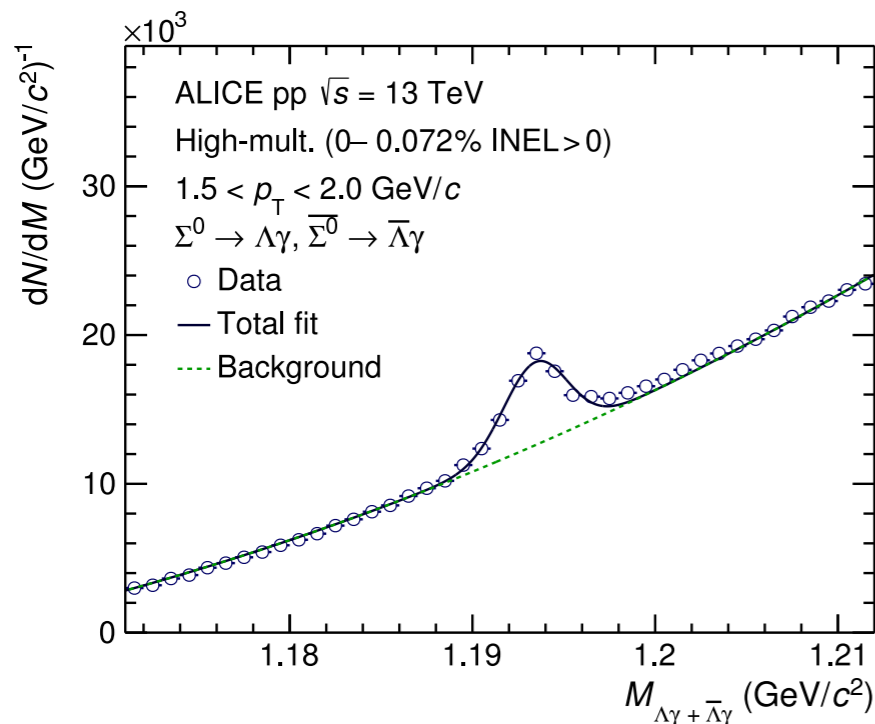
- Combination of spin singlet and triplet
- \* Extension to the low momentum regime
- \* <1% errors !!



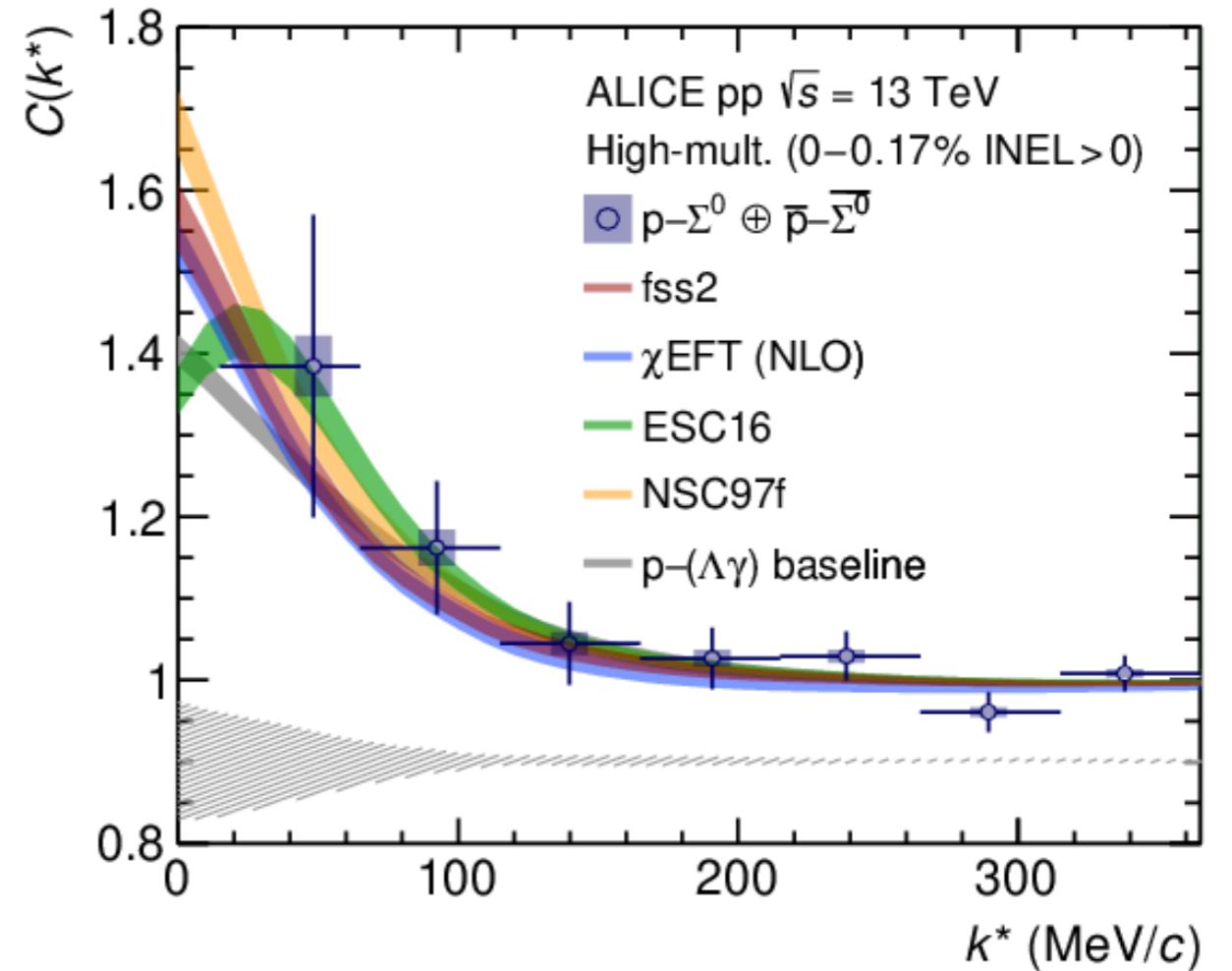
$C(k^*) > 1$ :  
Attractive  
interaction

ALICE Collaboration, *Phys.Lett.B* 805 (2020) 135419

- $\Sigma^0 \rightarrow \Lambda \gamma$  (BR: almost 100 %)
  - Identification of the photon via conversions
  - Significant contribution from correlated  $p-(\Lambda\gamma)$  background due to low purity
- Significant differences among the models will allow decisive measurements in future



ALI-PUB-337371



$\chi$ EFT: J. Haidenbauer et. al, Nucl. Phys. A915 (2013) 24–58.  
 NSC97f: T. A. Rijken et. al, Phys. Rev. C59 (1999) 21–40  
 ESC16: M. M. Nagels et. al, Phys. Rev. C99 (2019) 044003  
 fss2: Y. Fujiwara et al., Prog. Part. Nucl. Phys. 58, 439–520 (2007)

# $p-\Xi^-$ femtoscopy

## Benchmarking lattice QCD

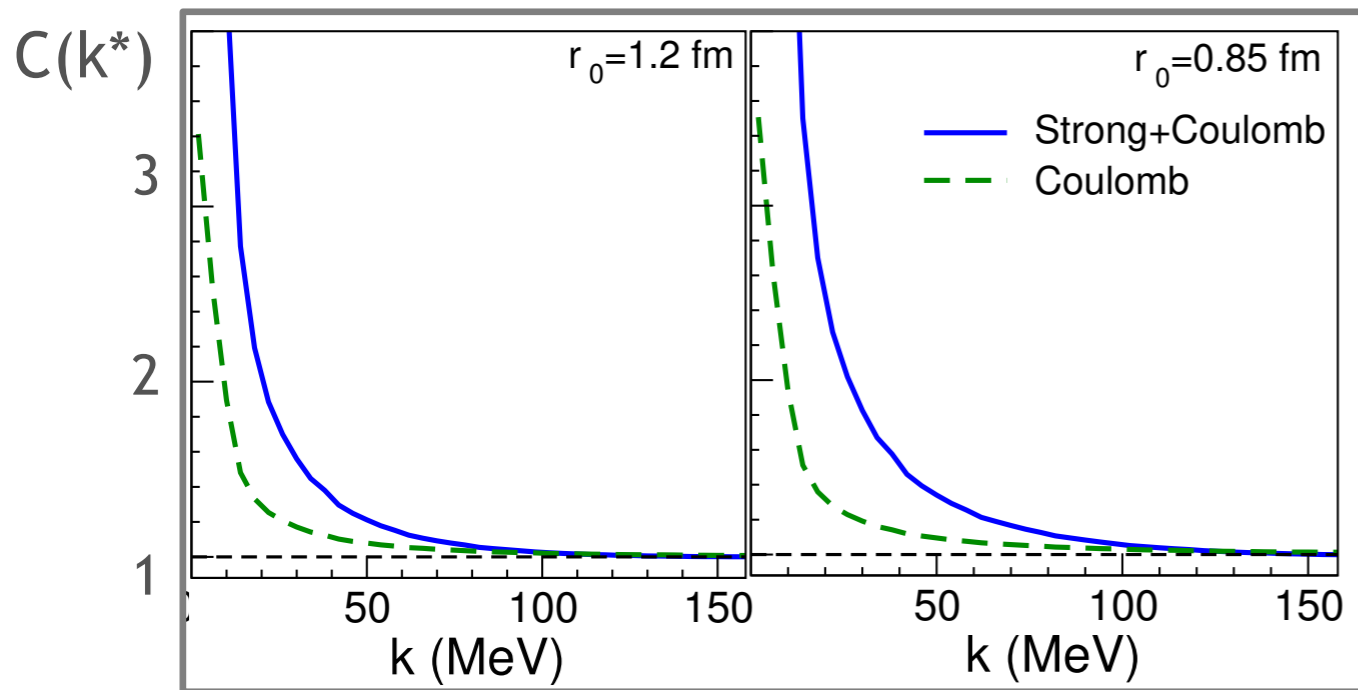
ALICE Collaboration, Phys. Rev. Lett. 123 (2019) 112002

$pp \sqrt{s} = 13 \text{ TeV}$  (high mult.)

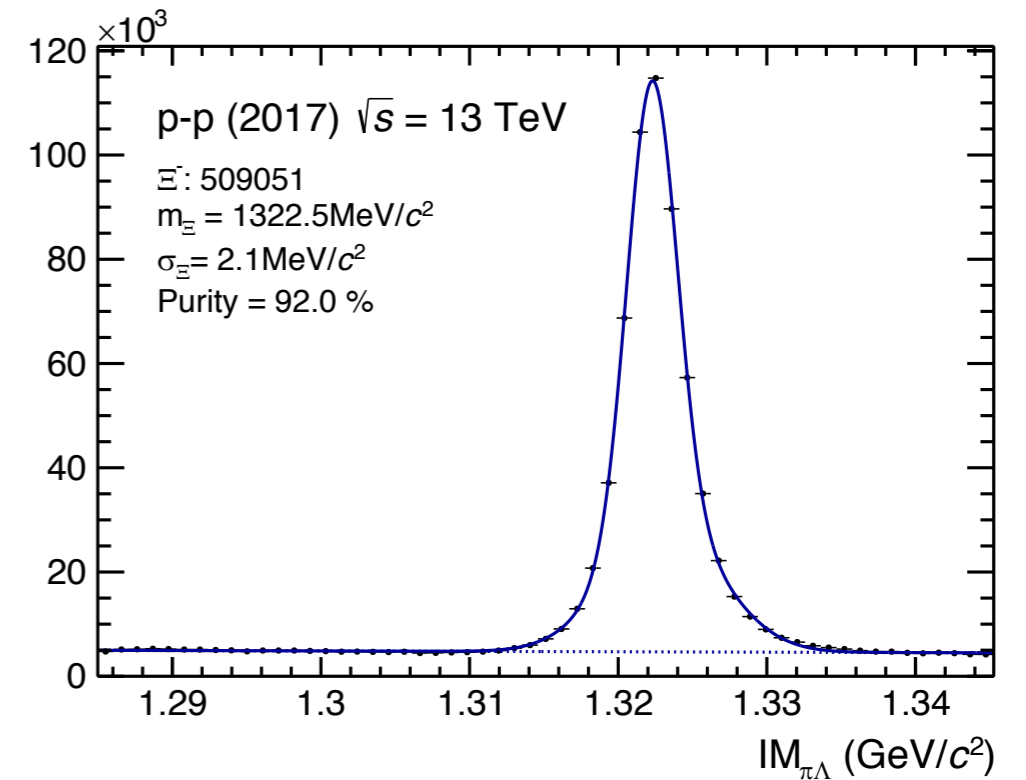
$p\text{-Pb } \sqrt{s_{NN}} = 5.02 \text{ TeV}$

- **Preliminary** calculations by the HAL QCD Collaboration
- Taking the strong interaction into account creates a significantly different Correlation function than Coulomb only

- Decay mode  $\Xi^\pm \rightarrow \Lambda + \pi$   
 $p + \pi$



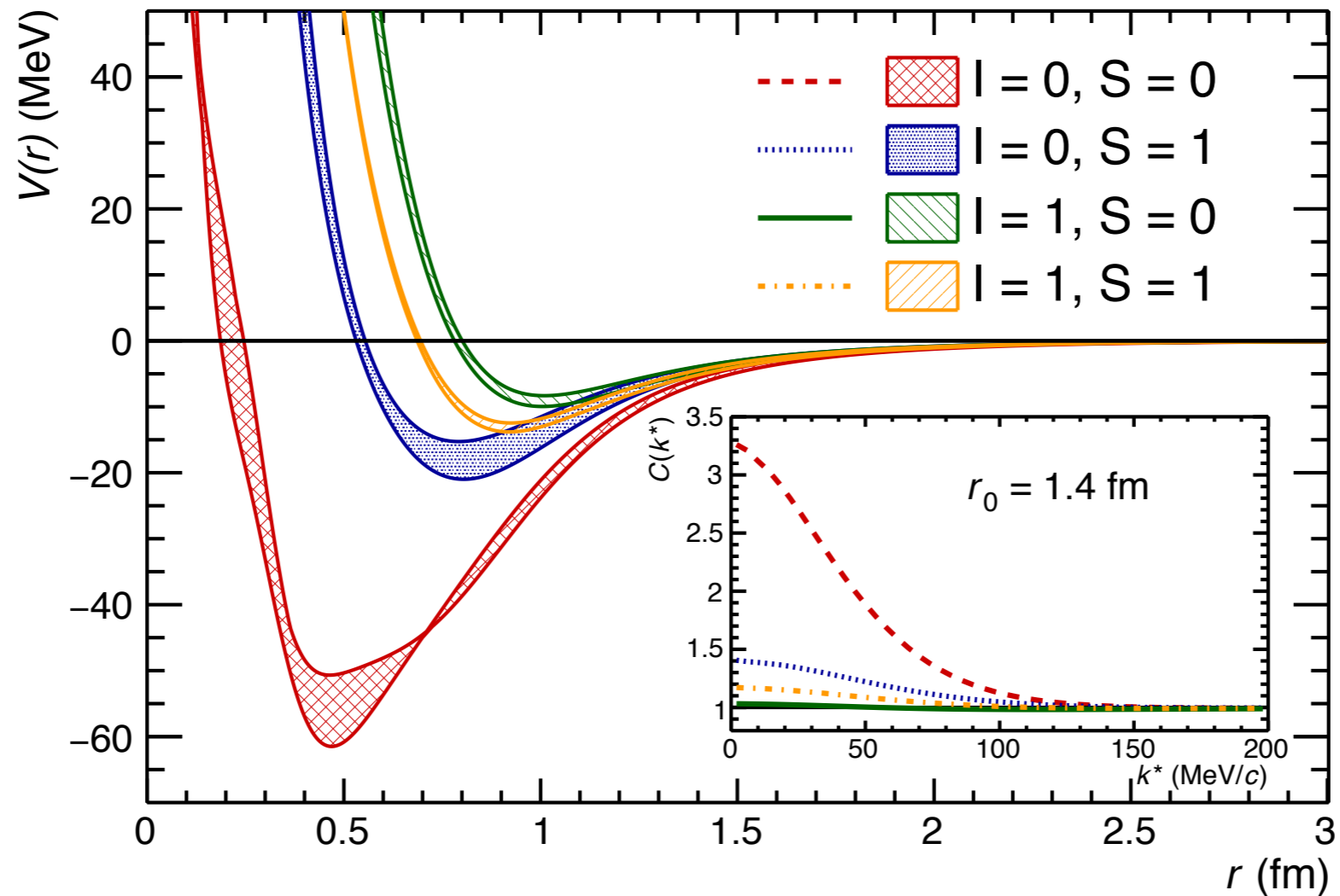
arXiv:1702.06241 , Nuclear Physics A 967 (2017) 856–859



**CATS** (D.L.Mihaylov et al. Eur.Phys.J. C78 (2018) no.5,394)

# Lattice Interaction

(Potential from Hatsuda et al., NPA967 (2017) 856, PoS Lattice2016 (2017) 116)



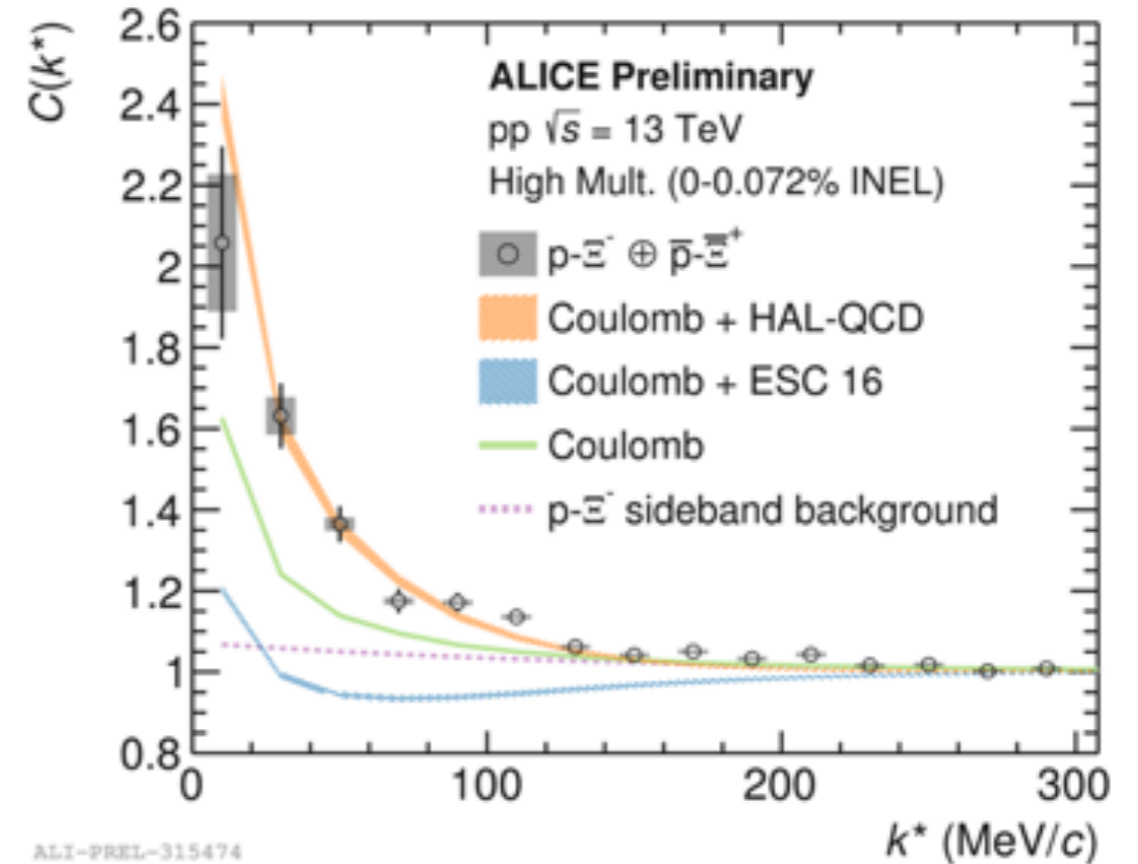
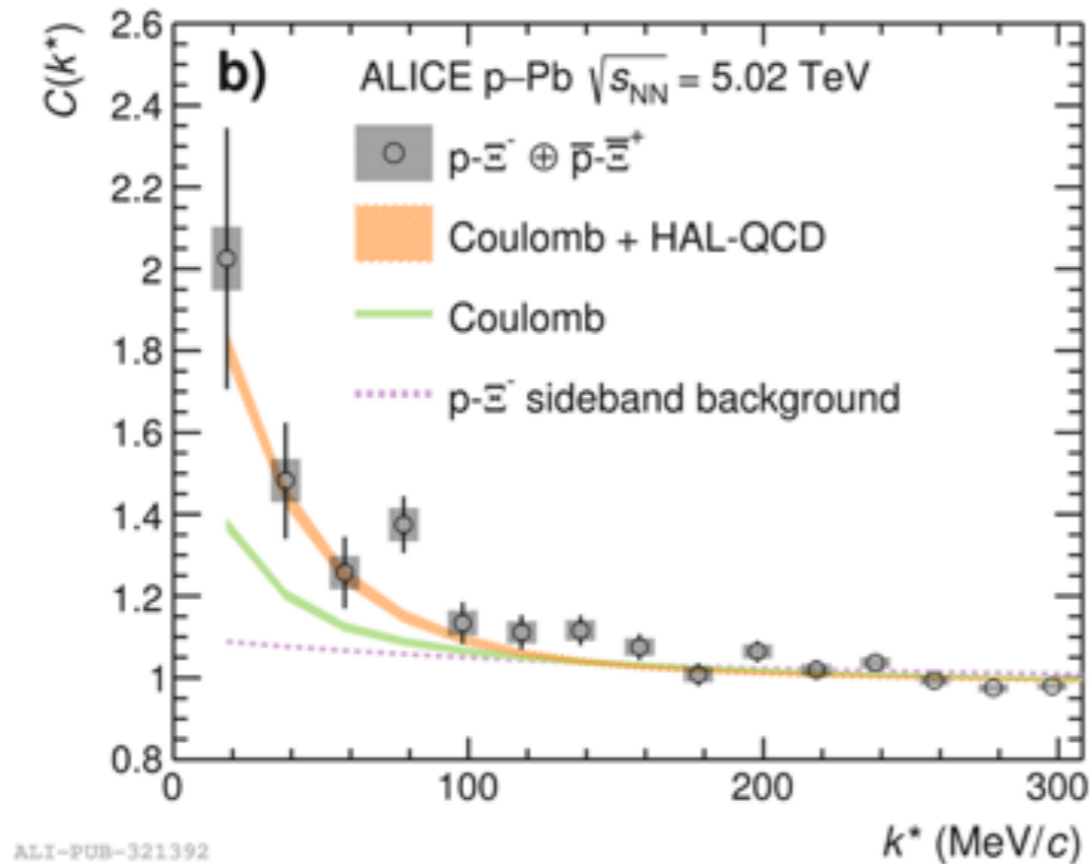
Errors due to different integration times

Each Potential can be converted in a correlation function via CATS

$$C(k^*) = \frac{1}{8} (C_{I=0}^{S=0} + C_{I=1}^{S=0}) + \frac{3}{8} (C_{I=0}^{S=1} + C_{I=1}^{S=1})$$

# proton- $\Xi^-$ Correlation Function

ALICE Collaboration, PRL 123 (2019) 112002



## First observation of the strong interaction in $p-\Xi^-$

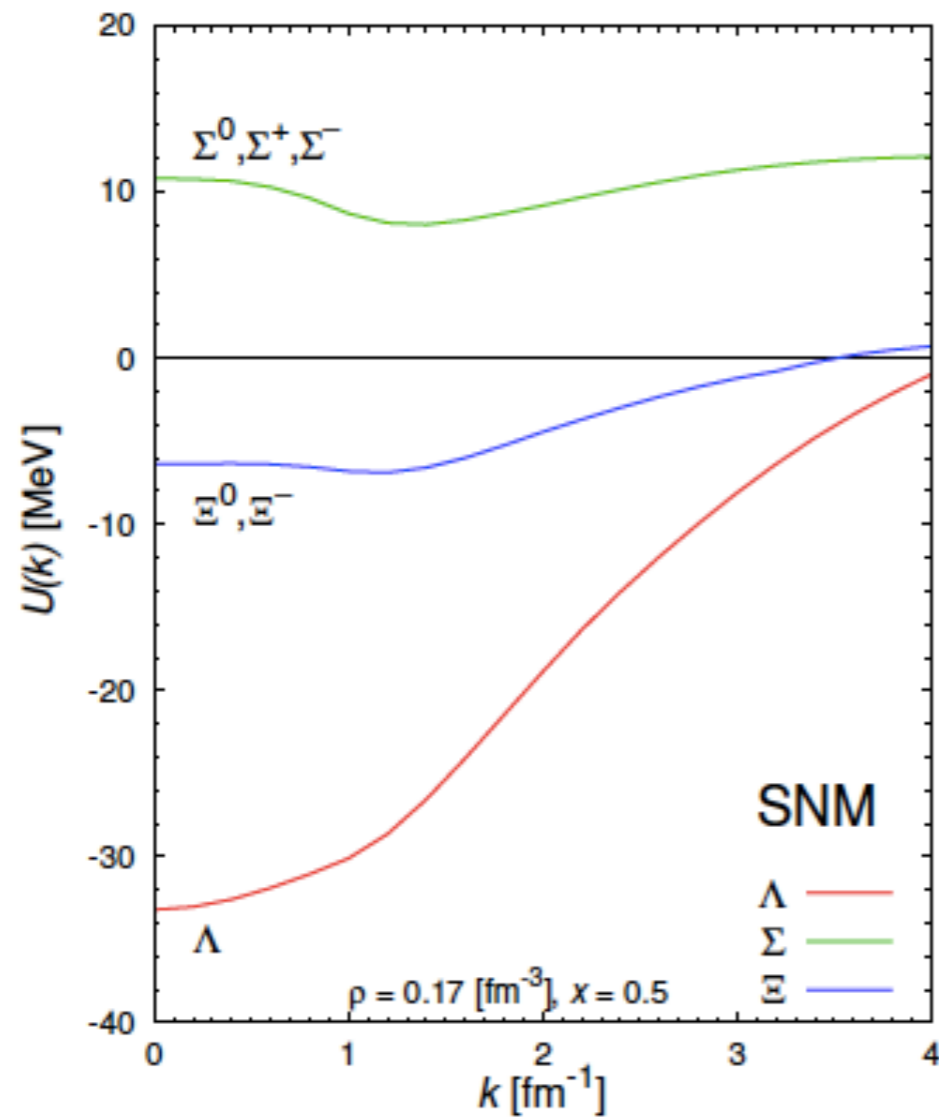
- Coulomb-only excluded ( $> 4 \sigma$ )
- Compatible with Lattice (HAL-QCD) calculations

## $p-\Xi^-$ in pp 13 TeV (high mult.)

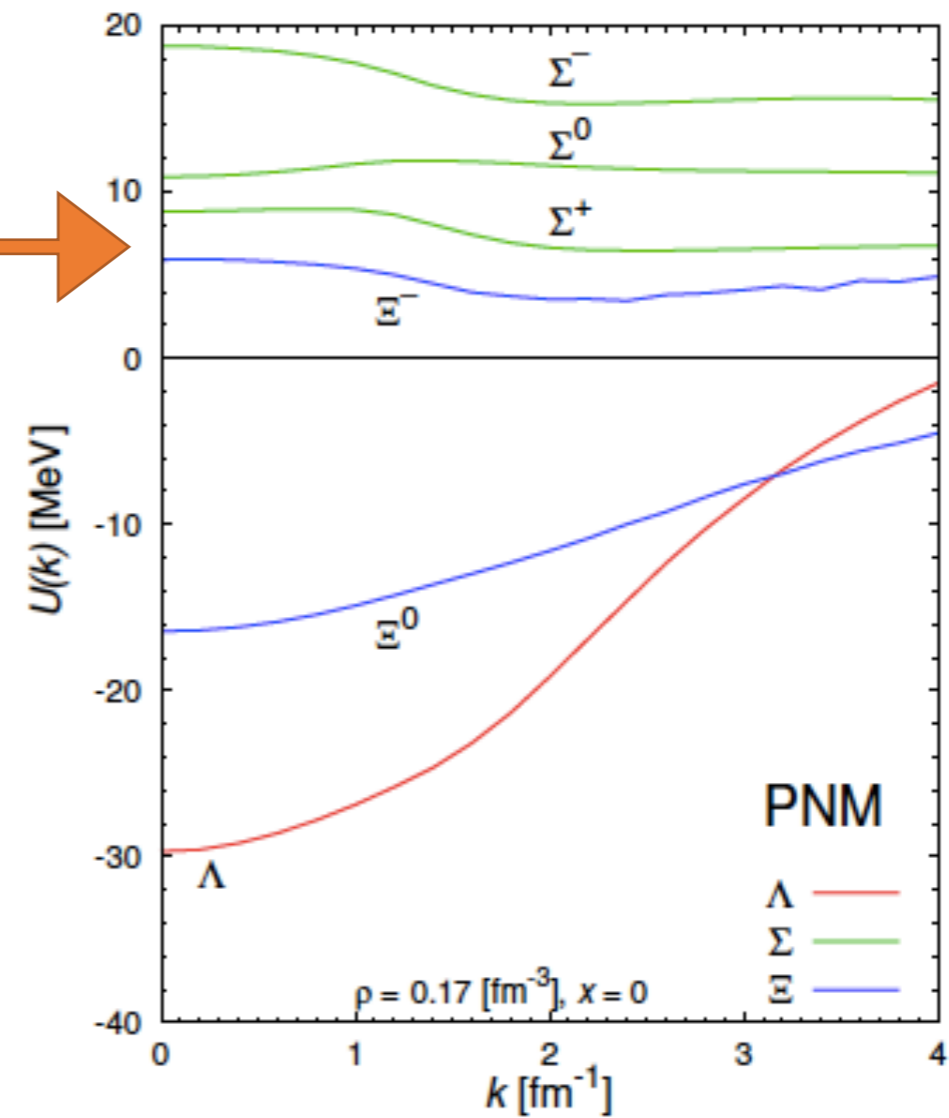
- Coulomb-only:  $> 5.7 \sigma$
- HAL-QCD: (1.3-2.5)  $\sigma$
- ESC16:  $> 18 \sigma$

## Symmetric Nuclear Matter

PoS Lattice2016 (2017) 116)



## Pure Neutron Matter



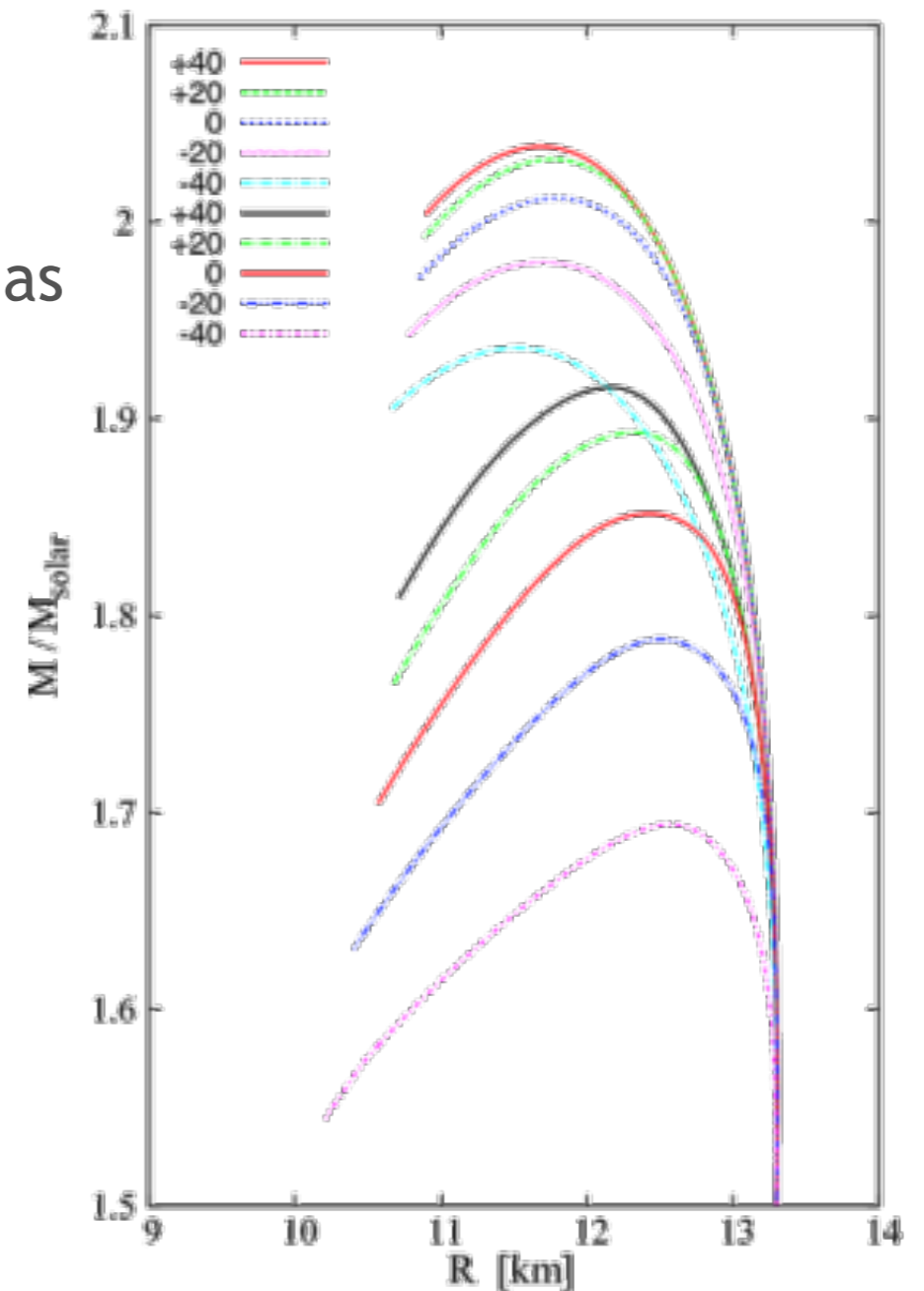
(Weissborn et al., NPA881 (2012) 62-77)

RMF models: EOS of neutron-rich matter with hyperon content

-> uses single particle potential at saturation densities as input

$$U_{NN}(\rho_0), U_{\Lambda N}(\rho_0), U_{\Sigma N}(\rho_0), U_{\Xi N}(\rho_0),$$

$= -30 \text{ MeV}$ 
 $= +30 \text{ MeV}$ 
variable ->





(Weissborn et al., NPA881 (2012) 62-77)

RMF models: EOS of neutron-rich matter with hyperon content

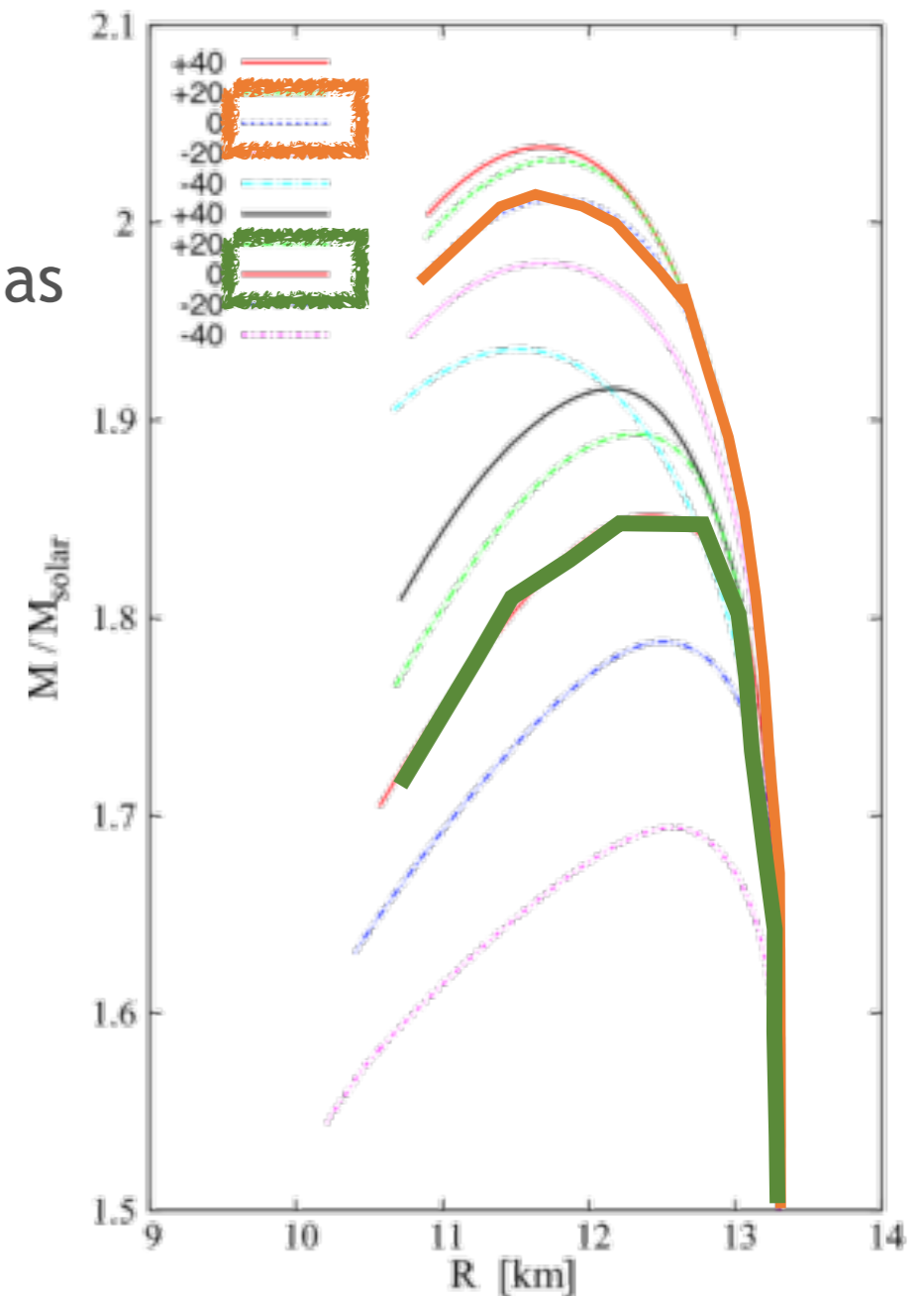
-> uses single particle potential at saturation densities as input

$$\begin{aligned}
 U_{NN}(\rho_0), U_{\Lambda N}(\rho_0), U_{\Sigma N}(\rho_0), U_{\Xi N}(\rho_0), \\
 = -30 \text{ MeV} \quad = +30 \text{ MeV}
 \end{aligned}$$

## Repulsive interaction

⇒ Production of  $\Xi$  pushed to higher densities

⇒ stiffer EoS, higher masses



We managed to become the new reference for the study of two hadrons interaction with strangeness content

- > Most precise data on  $\Lambda p$
- > First Measurement of the  $p\Sigma^0$
- > First evidence of the  $p\Xi^-$  strong attraction

In RUN3 ( from 2021 on) we expect factor 100 in statistics

What do we want more?

- > Three body interactions